

30 for observing dynamic behavior of the CSWs.

31

32 **Keywords:** Continental shelf wave; Arrested topographic wave; Cross-shelf SLA
33 structure; South China Sea

34

35 **1 Introduction**

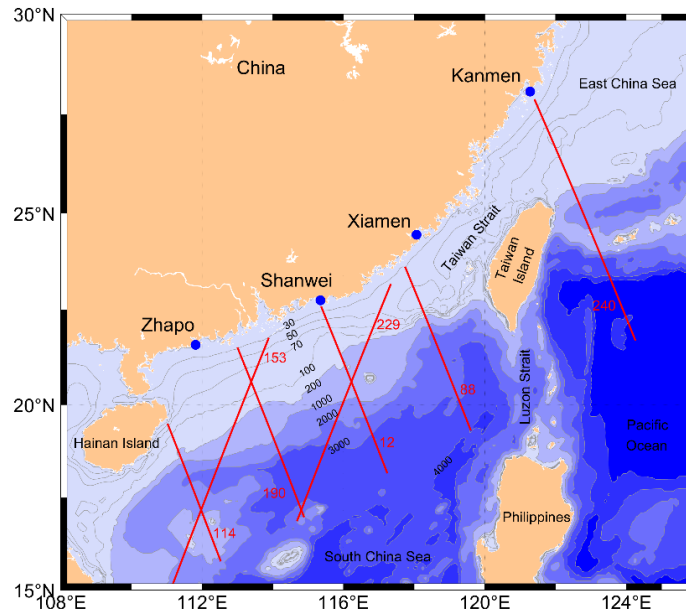
36 Continental shelf wave (CSW) is a type of topographic Rossby wave (TRW)
37 trapped in the continental shelf with amplitudes ranging from several tens' centimeters
38 to more than one meter (Aydın and Beşiktepe, 2022; Clarke and Brink, 1985; Heaps et
39 al., 1988; Morey et al., 2006; Mysak, 1980; Robinson, 1964; Zheng et al., 2015). CSW
40 is a sub-inertial motion with a wavelength much greater than the depth (Li et al., 2015;
41 Schulz et al., 2011). It propagates along the shelf with the coast on its right (left) in the
42 northern (southern) hemisphere (Clarke, 1977). During the impact of typhoon, an
43 excessive flooding in the coastal zone could be induced by a propagating CSW that
44 added to the locally wind-generated surge (Dukhovskoy and Morey, 2011; Han et al.,
45 2012). Therefore, CSW is particularly important for coastal sea-level variations.

46 CSW is generally generated by large-scale weather systems moving across or
47 along the shelf (Thiebaut and Vennell, 2010). CSW events have been reported by
48 previous investigators lasted from 2 days to 2 weeks (Chen and Su, 1987; Li et al., 2015;
49 Li et al., 2021; Zheng et al., 2015). The phase speed of CSWs depends on the bottom
50 topography, ranging from 5 to 20 m s⁻¹ (Li et al., 2015; Li et al., 2016; Shen et al., 2021).
51 CSWs could be taken as barotropic motions in a homogeneous coastal area. While in a
52 stratified ocean, it could be classified into coastal trapped waves. Overall, they are
53 resulting from conserving potential vorticity over the shelf (Chen et al., 2022; Quan et
54 al., 2021; Wang and Mooers, 1976).

55 The sea level variations in the South China Sea (SCS) are well depicted by these
56 previous studies as the continental shelf occupies about its half area (Ding et al., 2012;
57 Li et al., 2023b; Shen et al., 2021; Zhao et al., 2017; Zhou et al., 2023). However, two
58 issues should be improved. The first is that the primary data are usually obtained from
59 the tide-gauge stations along the coastline. The data have high accuracies, but represent

60 the sea level at the coasts only. Thus, satellite-altimeter-observed the sea level variations
 61 are often used to fill the data gaps in between the tide-gauge stations and on the
 62 continental shelf. The second issue is that the repeat period of the satellite altimeter is
 63 9.9 d, which is challenging to investigate the sea level variations with periods shorter
 64 than 10 d. Previous studies used the along-track sea level anomaly (SLA) from satellite
 65 altimeters to describe CSW (Chen et al., 2014; Li et al., 2016). However, the satellite
 66 altimeter with the sparse tracks (as shown in Fig. 1) could only capture the cross-
 67 structure of CSW with one or two snapshots of one CSW.

68



69

70 **Fig. 1.** Study area. Blue dots represent locations of tide-gauge stations Kanmen,
 71 Xiamen, Shanwei, and Zhapo. Red lines represent segments of ground tracks 12, 88,
 72 114, 153, 190, 229, and 240 for altimeter satellites over the continental shelf. Tracks 12,
 73 88, 114, and 190 are almost perpendicular to the coastline. Isobaths are in m.

74

75 Under the background of global warming, the sea level variation with very low
 76 frequency has been investigated by previous investigators. Ho et al. (2000) found
 77 seasonal sea level variability in the SCS using data from a satellite altimeter. Kajikawa
 78 and Yasunari (2005) investigated the interannual variability of the intra-seasonal
 79 variation over the SCS. Fang et al. (2006) analyzed low sea level along the eastern
 80 boundary of the SCS. Zhuang et al. (2010) found strong intra-seasonal variability in the

81 northern SCS. In addition, the sea level variations are influenced by thermodynamic
82 processes, e.g., eddies and thermal change of upper layer of the SCS (Cheng and Qi,
83 2007; Xie et al., 2018; Zheng et al., 2014).

84 Meanwhile, the upper layer thermal changes significantly influence the sea level
85 variations (Cheng and Qi, 2007). Using sea surface height (SSH) data from satellite
86 altimeters, Xu et al. (2016) found that sea level variations in the coastal area of the SCS
87 are still strongly influenced by the coastal current system in summer and winter.
88 Seasonal circulation is mainly driven by the monsoon wind stress (Gan et al., 2006).
89 Lin et al. (2021) applied the arrested topographic waves (ATWs) model to the coastal
90 mean dynamic topography along the East China Sea (ECS) and SCS. The mean
91 circulation in a coastal zone of variable depth may be modeled by linear equations (Wu,
92 2021). The result suggests that the mean dynamic topography is a counterbalance of
93 contributions from the along-shelf wind and bottom friction and well predicted by the
94 ATW model. Therefore, the monsoon winds are a control factor for the sea level
95 variation.

96 This study aims to investigate the cross-shelf structures of sea level over the
97 continental shelf. As the repeat period of the satellite altimeters, 9.9 d, is comparable
98 with that of CSWs in the northern SCS, the statistical characteristics of the along-track
99 SLA are applied to show the cross-shelf structure of CSWs using a long-term data set
100 from 1993 to 2020. To figure out the cross-shelf structures of ATWs is another goal.

101 The rest of the paper is organized as follows: Section 2 describes the observed data.
102 Section 3 presents the theory of CSWs and ATWs. Section 4 presents the characteristics
103 of the signals derived from the tide-gauge data and along-track SLA. Section 5
104 discusses the CSWs detected from tide-gauge data and the cross-shelf structure of sea
105 level over the continental shelf. Section 6 gives summaries.

106

107 **2 Data**

108 2.1 Along-track sea level anomaly data

109 Satellite altimeter along-track SLA data are produced and distributed by the
110 Archiving, Validation, and Interpretation of Satellite Oceanographic Data (AVISO),

111 Centre National d'Etudes Spatiales (CNES) of France. The data from 1993 to 2020 are
112 derived from TOPEX/Poseidon, Jason-1, Jason-2, and Jason-3 measurements. The
113 satellite repeat period is 9.9 d, and the temporal resolution of the along-track data is 1
114 Hz. The along-track SLA is calculated by subtracting the twenty-year mean from the
115 SSH measured by the satellite altimeters. The along-track SLA is low pass filtered using
116 a 7-point moving average. The ground tracks in the study area, 12, 88, 114, 153, 229,
117 and 240, are shown in Fig. 1.

118 Satellite altimetry provides a unique sea level dataset for the coastal sea level
119 research. A few recent studies have stressed the importance of small-scale coastal
120 processes on coastal sea-level variance (Cazenave and Moreira, 2022; Vignudelli et al.,
121 2019). The along-track SLA has been successfully validated and applied to the coast
122 zone by Birol et al. (2021). These studies present the availability of along-track SLA in
123 the coastal zones.

124

125 2.2 Sea level anomaly from tide-gauge

126 The tide-gauge data at stations Kanmen, Xiamen, Shanwei, Hongkong and Zhapo
127 (as shown in Fig. 1) are obtained from the Global Sea Level Observing System
128 (GLOSS). The data cover a period from 1993 to 1997, with a temporal resolution of 1
129 h. De-tided sea level anomaly (DSLAs) is calculated by removing tidal signals using a
130 Matlab toolbox (Pawlowicz et al., 2002).

131 The monthly sea level means at stations Xiamen, Shanwei, and Zhapo are obtained
132 from the Permanent Service for Mean Sea Level (PSMSL). Monthly mean data cover
133 periods of 1993-2003, 1993-1994, and 1993-2020 for stations Xiamen, Shanwei, and
134 Zhapo, respectively.

135

136 2.3 Sea surface wind stress

137 Monthly sea surface wind stress is derived from the Copernicus Marine
138 Environment Monitoring Service (CMEMS). The dataset covers a period from 1993 to
139 2020 with a spatial resolution of $0.25^{\circ} \times 0.25^{\circ}$. Sea surface wind stress data on the
140 satellite altimeter ground tracks are decomposed into the cross-shelf and along-shelf

141 components. The cross-shelf component is positive seaward and parallel to the satellite
 142 altimeter ground tracks (12, 88, 114 and 190). The along-shelf component is positive
 143 northward and perpendicular to these satellite altimeter ground tracks.

144

145 2.4 Topographic profile

146 The behavior of CSWs is determined by the topography of the continental shelf
 147 and slope, which has been well documented. The topographic profiles along the satellite
 148 altimeter ground tracks are extracted from a dataset of ETOPO-2. This study uses one-
 149 dimensional linear piecewise functions to fit the topographic profiles along the satellite
 150 altimeter ground tracks. The width of the continental shelf, depths of shelf break and
 151 deep basin along the tracks are listed in Table 1. The continental shelf break is extracted
 152 as the location of maximum change in the gradient of the continental slope.

153

154 Table 1. The width of the continental shelf, depths of shelf break and deep basin along
 155 the satellite ground tracks.

Track number	$l(\text{km})$	$H_1(\text{m})$	$H_2(\text{m})$
12	200	-300	-3800
88	178	-200	-2500
114	123	-200	-1800
153	259	-110	-1600
190	247	-255	-3500
229	244	-225	-3500
240	273	-250	-5000

156

157 3 Theory

158 3.1 Momentum equation for CSWs

159 The linearized shallow-water equations governing a barotropic ocean on a rotating
 160 earth are

$$161 \quad \frac{\partial u}{\partial t} - fv = -g \frac{\partial \eta}{\partial x} + \frac{\tau_s^x - \tau_b^x}{\rho H} \quad (1a)$$

$$162 \quad \frac{\partial v}{\partial t} + fu = -g \frac{\partial \eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H} \quad (1b)$$

$$163 \quad \frac{\partial \eta}{\partial t} + \frac{\partial(uH)}{\partial x} + \frac{\partial(vH)}{\partial y} = 0 \quad (1c)$$

164 where cross-shelf and along-shelf velocities (u, v) are depth-averaged in cross-shelf and
 165 along-shelf coordinates (x, y). η is the sea surface height. The Coriolis parameter is f .
 166 The bathymetry $H=H(x)$ is assumed to be a function of the cross-shelf variable, x only.
 167 τ_s and τ_b the surface and bottom stresses. g, ρ are the gravitational acceleration and
 168 the water density.

169 The scales of the along-shelf length of CSW ($L = 2\pi/k \approx 2 \times 10^3$ km), and
 170 cross-shelf length ($l \approx 200$ km) are subject to the long-wave assumption ($l/L \ll 1$).
 171 Under the long-wave assumption, a nondimensionalized could be applied to Eqs. (1a-
 172 b). Then, $\partial u/\partial t \ll fv$. We neglect the $\partial u/\partial t$ term, and Eqs. (1a-b) become (Li et al.,
 173 2016; Schulz et al., 2011):

$$174 \quad v = \frac{g}{f} \frac{\partial \eta}{\partial x} - \frac{\tau_s^x - \tau_b^x}{f \rho H} \quad (2a)$$

$$175 \quad u = -\frac{g}{f} \frac{\partial \eta}{\partial y} - \frac{\partial}{\partial t} \left(\frac{g}{f^2} \frac{\partial \eta}{\partial x} - \frac{\tau_s^x - \tau_b^x}{f^2 \rho H} \right) + \frac{\tau_s^y - \tau_b^y}{f \rho H} \quad (2b)$$

176 Substitute Eqs. (2a-b) into Eq. (1c), we obtain the equation governing SSH of
 177 CSWs

$$178 \quad \frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y - \tau_b^y}{f \rho} \right) + \frac{\partial^2}{\partial t \partial x} \left(\frac{\tau_s^x - \tau_b^x}{f^2 \rho} \right) - H \frac{\partial}{\partial y} \left(\frac{\tau_s^x - \tau_b^x}{f \rho H} \right) = 0 \quad (3)$$

180 Assume that CSWs are forced by along-shelf wind stress, i.e., $(\tau_s^x, \tau_s^y) = (0, \tau_s^y)$.
 181 The bottom friction is neglected to simplify the calculation in Eq. (3). Therefore, the
 182 equation governing the SSH of CSWs becomes

$$183 \quad \frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y}{f \rho} \right) = 0 \quad (4)$$

184 The change in the SSH of CSWs is balanced by the variation of along-shelf wind
 185 stress in cross-shelf direction. Assume a periodic along-shelf wind stress, $\tau_s^y =$
 186 $\tau_0 \exp[i(\alpha y + \omega t)]$ (where α is wavenumber of wind stress, and $\tau_0 = \text{constant}$). Eq.
 187 (4) becomes

$$188 \quad \frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} = 0 \quad (5)$$

189 which means that the change in SSH is independent of along-shelf wind stress.

190 The assumption of a wave solution

$$191 \quad \eta = \phi(x)\exp[i(ky + \omega t)] \quad (6)$$

192 yields an equation for $\phi(x)$,

$$193 \quad H \frac{d^2\phi}{dx^2} + \frac{dH}{dx} \frac{d\phi}{dx} + \left(\frac{fk}{\omega} \frac{dH}{dx} - \frac{f^2}{g} \right) \phi = 0 \quad (7)$$

194

195 3.1.1 Over the shelf

196 For $0 \leq x \leq l$, $H = H_1x/l$, the equation for $\phi(x)$ is

$$197 \quad x \frac{d^2\phi}{dx^2} + \frac{d\phi}{dx} + \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1} \right) \phi = 0 \quad (8)$$

198 which is subject to the following boundary conditions: $\phi(0) = a$, and $\phi(l) = A$. a and
199 A should be arbitrary. We could take a as the amplitude of fluctuation in sea level from
200 the tidal gauge station.

201 The solution to Eq. (8) is expressed as the sum of the first and second kinds of
202 Bessel functions (Robinson, 1964; Schulz et al., 2011)

$$203 \quad \phi(x) = aJ_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) + bY_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) \quad (9)$$

204 where a is arbitrary constant. As the solution for $\phi(x)$ is finite, therefore

$$205 \quad \phi(x) = aJ_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) \quad (10)$$

206

207 3.1.2 In the deep basin

208 For $l < x$, $H = H_2$, the equation for $\phi(x)$ is

$$209 \quad \frac{d^2\phi}{dx^2} - \frac{f^2}{gH_2} \phi = 0 \quad (11)$$

210 which is subject to the following boundary conditions: $\phi(x \rightarrow l) = A$, and $\phi(\infty) = 0$.

211 The solution is

$$212 \quad \phi(x) = A \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) + B \exp \left(\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) \quad (12)$$

213 As the solution for $\phi(x)$ is also finite in the deep basin, i.e.,

$$214 \quad \phi(x) = A \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) \quad (13)$$

215 where A is arbitrary constant.

216

217 3.1.3 Dispersion relation

218 As the fluid is continuous at the edge of continental shelf, therefore

$$219 \quad H_1 \cdot u|_{x \rightarrow l_-} = H_2 \cdot u|_{x \rightarrow l_+} \quad (14a)$$

$$220 \quad aJ_0 \left(2 \left(\frac{fkl}{\omega} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) = A \quad (14b)$$

221 From Eq. (14a), we have

$$222 \quad H_1 \cdot a \left(kJ_0 + \frac{\omega}{f} \left(\frac{fkl}{\omega} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} J_0' \right) = H_2 \cdot A \left(k - \frac{\omega}{\sqrt{gH_2}} \right) \quad (15)$$

223 where J_0' is the derivation of the zero-order of the first kind Bessel function,

$$224 \quad J_0'(x) = -\frac{x}{2} (J_0(x) + J_2(x)) \quad (16)$$

225 where J_2 is the second order of J_0 .

226 Substituting Eq. (16) into Eq. (15) yields a dispersion relation for CSWs

$$227 \quad \left(\frac{cfl}{gH_2} - 1 + \frac{c}{(gH_2)^{\frac{1}{2}}} \right) J_0 - \left(\frac{H_1}{H_2} - \frac{cfl}{gH_2} \right) J_2 = 0 \quad (17)$$

228 where $c (= \frac{\omega}{k})$ is phase speed of CSWs. We solve Eq. (17) using the zero-finding

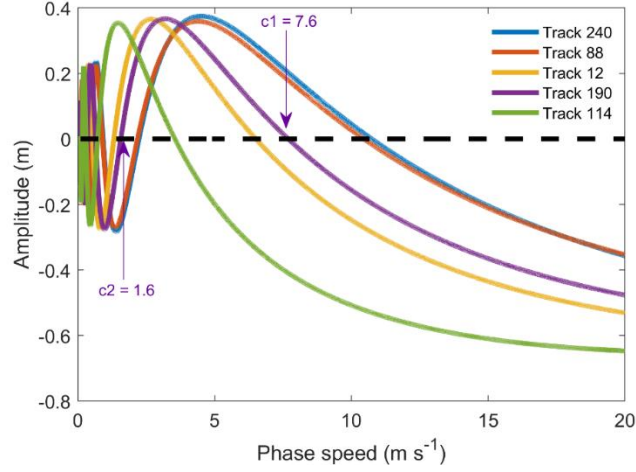
229 function in MATLAB. The solution of phase speed for CSWs is shown as Fig. 2. The

230 zero-crossing points for each curve present the phase speed of CSWs. The first zero-

231 crossing point on the right hand of each curve points out the phase speed of mode-1

232 CSW, e.g., $c = 7.6 \text{ m s}^{-1}$ for track 190.

233



234
 235 **Fig. 2.** Solution of phase speed for CSWs from Eq. (17). The zero-crossing values
 236 represent the phase speed of CSWs for different modes. Colorful curves are amplitude
 237 of zero-order of the first kind Bessel function in different modes for the idealized depth
 238 profile of tracks 12, 88, 114, 190, and 240, respectively. Arrows point out the phase
 239 speed of mode-1 (7.6 m s^{-1}) and mode-2 (1.6 m s^{-1}) CSWs along track 190.

240

241 3.1.4 Cross-shelf structure

242 With Eqs. (10), (13) and (14b), the wave solution of SSH, i.e., Eq. (3), is

$$243 \quad \eta = \begin{cases} \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{fk_i}{\omega_i} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) \exp[i(k_i y + \omega_i t)] & x \leq l \\ \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{fk_i l}{\omega_i} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) \exp[i(k_i y + \omega_i t)] & x > l \end{cases} \quad (18)$$

244

245 where $\exp[i(ky + \omega t)]$ is the waveform propagating along the shelf. $a_i \cdot J_0$ is the
 246 cross-shelf structure of the waveform for mode i over the shelf.

247 In the cross-shelf direction, the SSH becomes

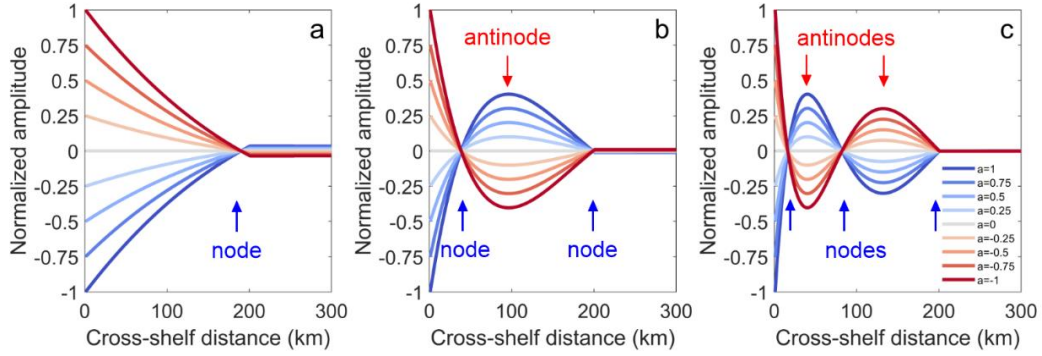
$$248 \quad \eta = \begin{cases} \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{fk_i}{\omega_i} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) & x \leq l \\ \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{fk_i l}{\omega_i} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) & x > l \end{cases} \quad (19)$$

249 where $\eta_{0i}(t) (= a_i \exp[i(k_i y + \omega_i t)])$ is time series of SSH at the coastline for mode
 250 i .

251 Fig. 3 shows the SSH structure over the continental shelf. The SSH mode 1 looks
 252 like a bell mouth. For Modes 2 and 3, nodes and antinodes appear on the shelf. Node
 253 for Mode 2 appears at 50 km off the coast over the shelf. Nodes for Mode 3 appear at

254 30 and 90 km off the coast.

255



256

257 **Fig. 3.** (a-c) Cross-shelf structure of normalized amplitudes of the first three CSW
 258 modes for idealized depth profile of track 12. Gradient color curves represent amplitude
 259 structure over time for CSWs in the cross-shelf direction.

260

261 3.2 Steady situation

262 The governing equations in the steady situation for Eqs. (1a-c) are

$$263 \quad -fv = -g \frac{\partial \eta}{\partial x} \quad (20a)$$

$$264 \quad fu = -g \frac{\partial \eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H} \quad (20b)$$

$$265 \quad \frac{\partial(uH)}{\partial x} + \frac{\partial(vH)}{\partial y} = 0 \quad (20c)$$

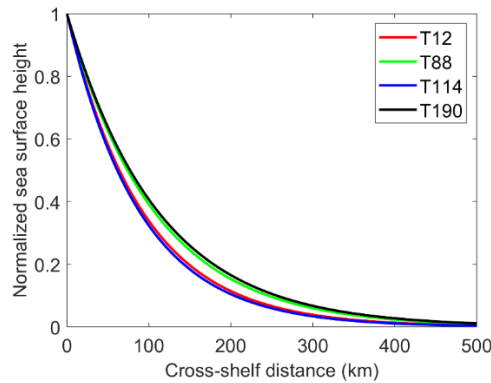
266 A linear drag is used for the bottom friction, i.e., $\overline{\tau_b} = \rho \lambda \vec{u}$ (Hsueh and Pang, 1989; Lin
 267 and Yang, 2011). The solution for SSH based on Eqs. (20a-c) gives (Csanady, 1978)

$$268 \quad \eta = B \tau_0 e^{-\frac{x}{L}} \sin\left(\alpha y + \frac{x}{L} - \frac{\pi}{4}\right) \quad (21)$$

269 where $L = \sqrt{2\lambda l / \alpha f H_1}$ is the scale width of trapped sea level in the cross-shelf
 270 direction, B is an arbitrary constant. A drag coefficient, $\lambda = O(5 \times 10^{-4}) \text{ m s}^{-1}$, is used
 271 for the linear bottom friction (Chapman, 1987; Lin and Yang, 2011). The typical
 272 magnitude for winter wind stress over the northern continental shelf of the SCS is
 273 $O(0.1) \text{ N m}^{-2}$ (Lin and Yang, 2011; Lin et al., 2011). $f = 5 \times 10^{-5} \text{ s}^{-1}$, and $\alpha =$
 274 $2\pi/4000 \text{ km}$ evaluated by Lin et al. (2021).

275 The normalized SSH in the cross-shelf direction of ATW for tracks 12, 88, 114,
 276 and 190 are shown in Fig. 4. One can see that the trapped sea level in the cross-shelf
 277 direction decays quickly from 1 at the coastline to ~ 0.2 at the edge of the continental

278 shelf (~200 km), and 0.1 at a distance of 300 km. The ATW amplitude decays offshore,
279 and $L = \sim 100$ km in the study area, which is much less than the local Rossby radius of
280 deformation (~600 km). Under different wind stresses, the amplitude of ATW evolves
281 similarly to that in Fig. 3a. As Track 240 is not perpendicular to the coastline, it is
282 beyond the scope.
283



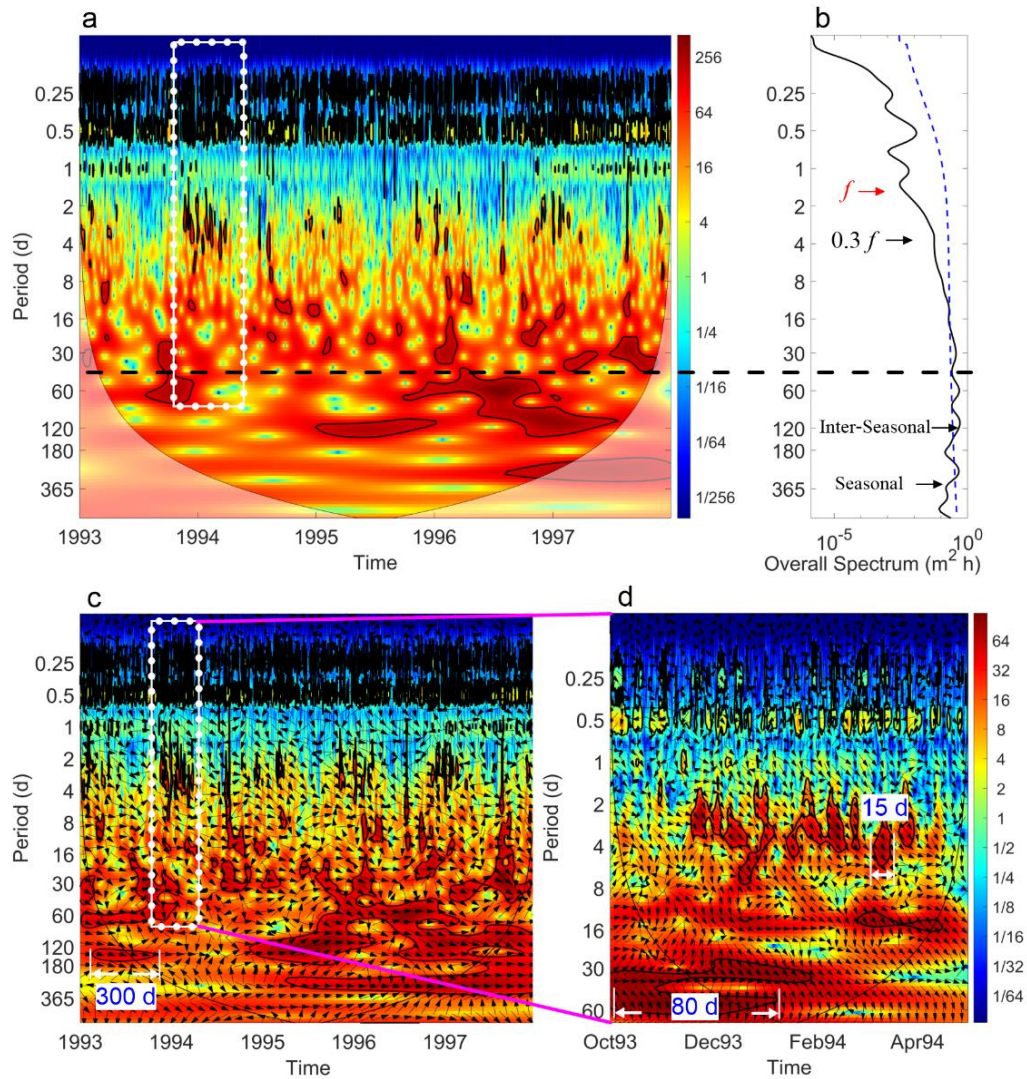
284
285 **Fig. 4.** Normalized SSH in cross-shelf direction for tracks 12, 88, 114, and 190.
286

287 4 Signals in sea level anomaly

288 4.1 Tide-gauge data

289 Fig. 5 shows the wavelet transform (WT) of SLA at station Xiamen from 1993 to
290 1997. One can see abundant signals with periods from several days to one year in Figs.
291 5a-b. A dozen signals with periods from several days to about one month could be seen
292 from the WT of SLA every year. The duration of these signals, with a period of several
293 days, is about ten days (e.g., in Fig. 5d). Moreover, the signals with a period of about
294 one month sustain for about three months (Figs. 5c-d). Even if significance level against
295 red noise is larger than 5%, the power is universal and continuous in the period bands
296 of 10-60 d. The main reason is that these signals are not so significant compared with
297 the signals with large amplitude. The wavelet analysis also exhibits a significant inter-
298 seasonal and seasonal variation of SLA. The characteristics of the signals at stations
299 Kanmen, Shanwei, and Zhapo are almost the same (not shown here). Therefore, the
300 variation of SLA along the northern SCS coast is universal.

301



302

303 **Fig. 5.** Temporal variability of SLA at the coastline. (a) WT of SLA at station Xiamen.
 304 (b) Overall spectrum of SLA at station Xiamen. (c) XWT of SLA between stations
 305 Kanmen and Xiamen. (d) XWT of SLA between stations Kanmen and Xiamen from
 306 October 1993 to April 1994. White line and arrows are auxiliary lines pointed out
 307 duration. The thick line is a 5% significance level against red noise, and the cone of
 308 influence (COI) is shown as the thin line. Color codes of power spectra normalized by
 309 variance are in arbitrary units. White dotted rectangles in (a) and (c) show the temporal
 310 domain of (d). The Blue dashed curve in (b) is a 5% significance level. Arrows in (b)
 311 point out the characteristic frequencies of inertial oscillation and CSWs. The black
 312 dashed lines in (a) and (b) represent signal period boundary (40 d). The arrows in (c)
 313 and (d) show the relative phase relationship between SLA at Kanmen and Xiamen with
 314 in-phase (anti-phase, leading, and lagging) pointing right (left, down, and up). In (d),
 315 arrows point down-right (about $\pi/3$), indicating SLA in Kanmen is leading that in
 316 Xiamen.

317

318 Figs. 5c-d show the cross wavelet transform (XWT) of SLA between stations

319 Kanmen and Xiamen. One can see that the period band and the occurrence time of the
320 significant cross wavelet power are consistent with that in Fig. 5a. The signals with
321 periods shorter and longer than 40 d show remarkably different characteristics. In the
322 period band shorter than 40 d, one can see the arrows generally point down-right. The
323 quasi-uniform phase lag indicates the fixed time delay of the signal between two tidal
324 gauge stations. However, the direction of arrows is in disorder for the period band
325 longer than 40 d, which indicates there is no evidence for propagating.

326 In the period band shorter than 40 d, the signals at station Xiamen lag that at station
327 Kanmen about 15 h (time lag=phase difference/ 2π ×period of signal). The propagation
328 phase speed of the sea level signal could be calculated by the lag time of sea level
329 propagation between stations Kanmen and Xiamen. The result is about 9 m s^{-1} , which
330 is very close to that reported by a number of recent studies (Ding et al., 2012; Li et al.,
331 2015; Li et al., 2016; Zhao et al., 2017).

332 In the period band longer than 40 d, the phase of signals between Kanmen and
333 Xiamen is a little complicated. The signal phase with the period of 360 d from 1995 to
334 1997 is almost 0. However, the seasonal signals at station Xiamen lag (lead) that at
335 station Kanmen in winter (summer) about $\pi/4$ - $\pi/2$, implying that the signal propagates
336 very slowly along the coast. Csanady (1978) concluded that these signals are a kind of
337 ATWs.

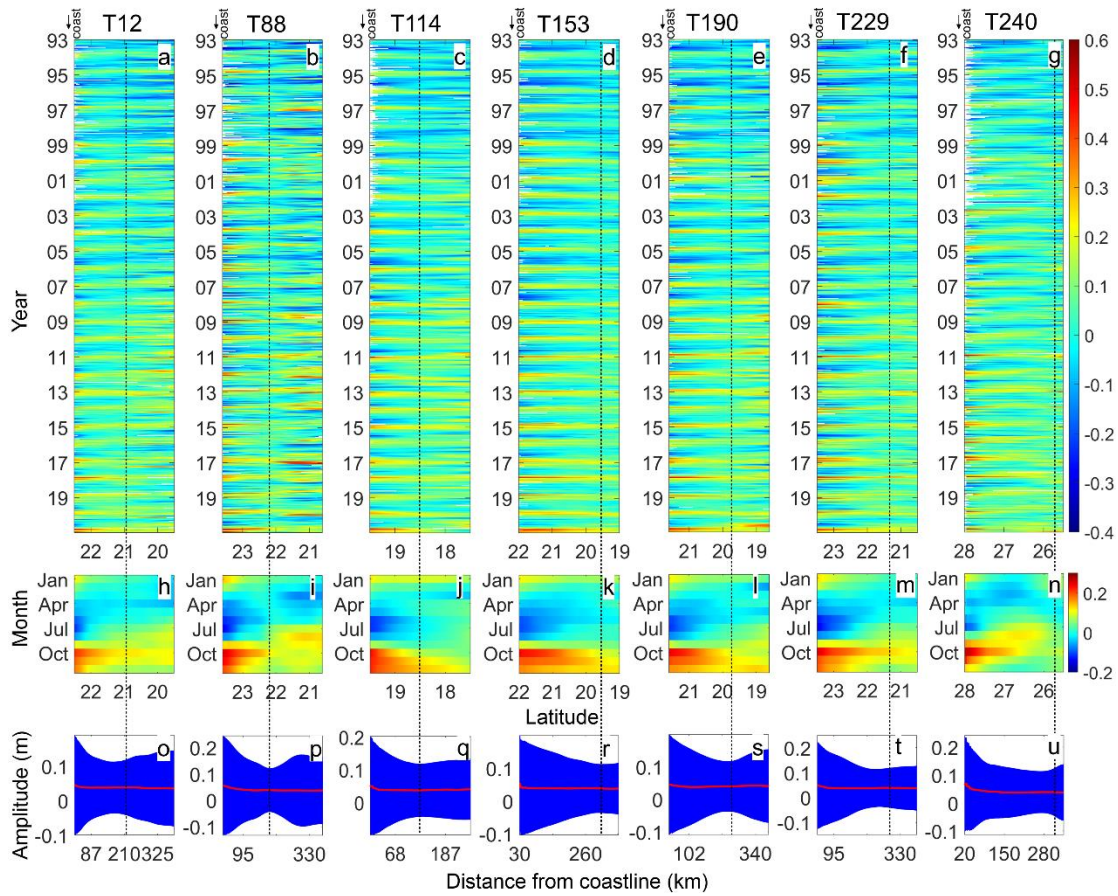
338

339 4.2 Along-track SLA

340 Fig. 6 shows the latitudinal distribution of the along-track SLA over the northern
341 continental shelf of the SCS from 1993 to 2020. One can see the clear annual signals
342 near the coast in Figs. 6a-g. Signals with periods shorter than one year could also be
343 seen. As shown in Figs. 6h-n, a clear seasonal cycle is discernable from the
344 climatological monthly mean of the along-track SLA. One can see that lower (higher)
345 sea levels over the shelf exist from March to August (September to February). The
346 trough ($\sim 0.2 \text{ m}$) and peak ($\sim 0.24 \text{ m}$) of sea level near the coast occur in July and
347 October, respectively. While in track 240, climatological monthly mean of the along-
348 track SLA on the shelf is smaller than that in track 88 especially in July. The 28-year

349 mean value of the along-track SLA is about 0.04 m, as shown in Figs. 6o-u.

350



351

352 **Fig. 6.** (a-g) Latitudinal distribution of the along-track SLA over the northern
 353 continental shelf of the SCS from 1993 to 2020 for tracks 12, 88, 114, 153, 190, 229,
 354 and 240, respectively. (h-n) Climatological monthly mean of the along-track SLAs for
 355 Tracks 12, 88, 114, 153, 190, 229, and 240. (o-u) Mean (red curves) and standard
 356 deviation (blue shadow) of the along-track SLA from 1993 to 2020 for tracks 12, 88,
 357 114, 153, 190, 229, and 240. Track numbers are shown on the tops of panels. Vertical
 358 dashed lines represent shelf break positions along the tracks, i.e., H_1 , listed in Table 1.
 359 The coastline position is marked by an arrow on top left of each panel.

360

361 Moreover, the standard deviations (STD) of the along-track SLA from 1993 to
 362 2020 (blue shadow in Figs. 6o-u) show a bell-mouth-like structure over the shelf. The
 363 amplitudes of the along-track SLA reach the maxima near the coast. The minimum
 364 variance exists near the shelf edge. The cross-shelf structure of SLA indicates that sea-
 365 level signals depend on the shelf depth.

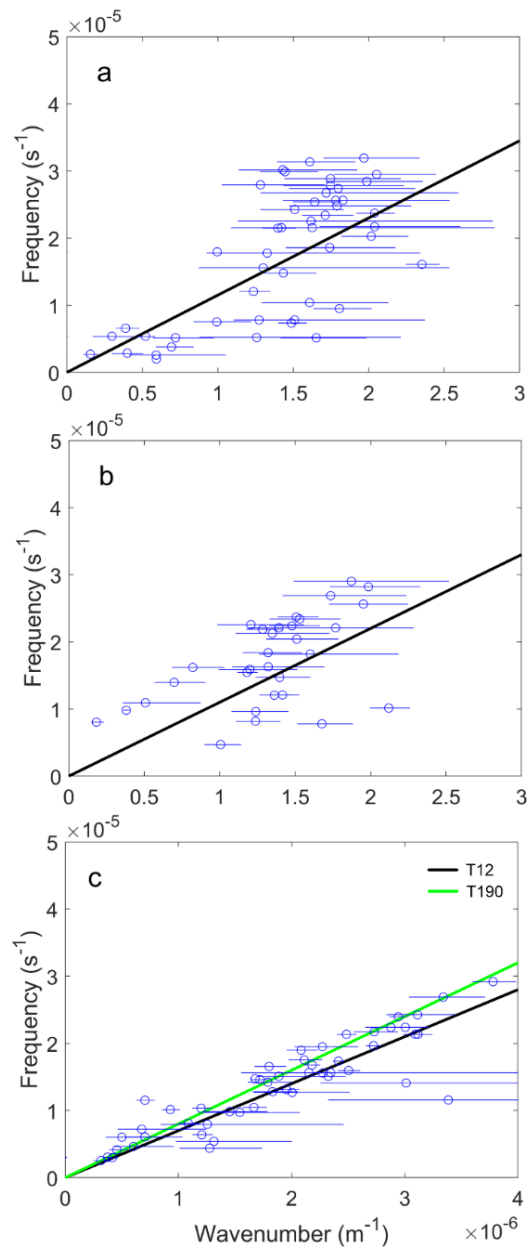
366

367 **5 Discussion**

368 5.1 Propagation of CSWs

369 Fig. 7 compares the data derived from the XWT of SLA with the dispersion
 370 relation of CSW. We overlay the dispersion relation curves of CSWs with the results
 371 (phase speed, c and period, T , $\lambda = c \cdot T$) from tide-gauge data analysis in Section 4.1.
 372 One can see that the data points derived from station pairs are distributed near the
 373 dispersion relation of CSWs, implying the signals with the periods shorter than 40 d are
 374 the CSWs propagating along the shelf.

375



376
 377 **Fig. 7.** Dispersion relation of the lowest mode of CSWs between (a) Kanmen and
 378 Xiamen, (b), Xiamen and Shanwei, as well as (c) Shanwei and Zhapo. The data points

379 are calculated from the XWT of SLA. The curves are the theoretical dispersion relation
380 for the mean depth profiles listed in Table 1. Black and green curves in (c) represent the
381 dispersion relation (from Fig. 2) for the topographic profiles along tracks 12 and 190
382

383 Figs. 7a-b deal with the signals between Kanmen and Shanwei. We use the
384 bathymetric profile near Track 240 and 88 of the altimeter satellite to calculate the
385 dispersion relation curve, respectively. However, since the topography between
386 Kanmen and Shanwei changes dramatically, the data points lie dispersedly around the
387 curve. The wavenumbers of CSWs range from 0.1 to $2.4 \times 10^{-5} \text{ m}^{-1}$ between stations
388 Kanmen and Shanwei. The CSWs propagate along the coast with a phase speed of 10
389 m s^{-1} .

390 Fig. 7c shows the signals between Shanwei and Zhapo. The data points and the
391 theoretical curves agree quite well. Compared with the topography between Kanmen
392 and Shanwei, the topography changes slightly between stations Shanwei and Zhapo.
393 Thus, we conclude that the signals in tide-gauge data with the periods shorter than 40
394 d are the CSWs propagating along the shelf. The phase speed of CSWs is about 8 m s^{-1}
395 ¹ in the study area, which is close to that of previous studies (Li et al., 2015; Li et al.,
396 2016; Shen et al., 2021).

397 Previous studies present that the periods of CSWs range from two days to two
398 weeks in the SCS (Chen and Su, 1987; Li, 1993). The period of CSWs upstream, i.e.,
399 in the ECS, is often detected as several days (Ding et al., 2012; Ding et al., 2018; Hsueh
400 and Pang, 1989; Huang et al., 2015; Yin et al., 2014). Hsueh and Romea (1983) found
401 the sea level fluctuations with a period of more than 13 d along the West Korean coast.
402 Worldwide, these low-frequency CSWs are common along the coast of Chile and the
403 east coast of the Indian Ocean (Castro and Lee, 1995; Hormazabal et al., 2001; Marshall
404 and Hendon, 2013; Vialard et al., 2009), where the width of the continental shelf is
405 narrow. In this study, we define the CSW with a maximum period of about 30-40 d. The
406 main reason for the difference is the data length we used. The long-time series of DSLA
407 helps analyze the abnormal low-frequency CSWs in the SCS.

408 In addition, we should take care that CSWs in Fig. 7 are mixed with wind-forced
409 CSWs. It is difficult to separate the effect of wind-force and free propagating CSWs

410 clearly. Hsueh and Romea (1983) found that there is a clear coupling between surface
411 winds and coastal sea level in the northeast China Sea. Li et al. (2016) found that the
412 propagation time of the wind signals are much shorter than that of CSWs in the north
413 SCS. Even though one can see isolated cases which are away from the dispersion
414 relationship of CSW in Fig. 7, most of the data points are near the theoretical dispersion
415 relation. Moreover, the stratification in the continental shelf is important for the
416 characteristics of CSW. The sea level variation in this study should show coastal trapped
417 waves influenced by stratification.

418

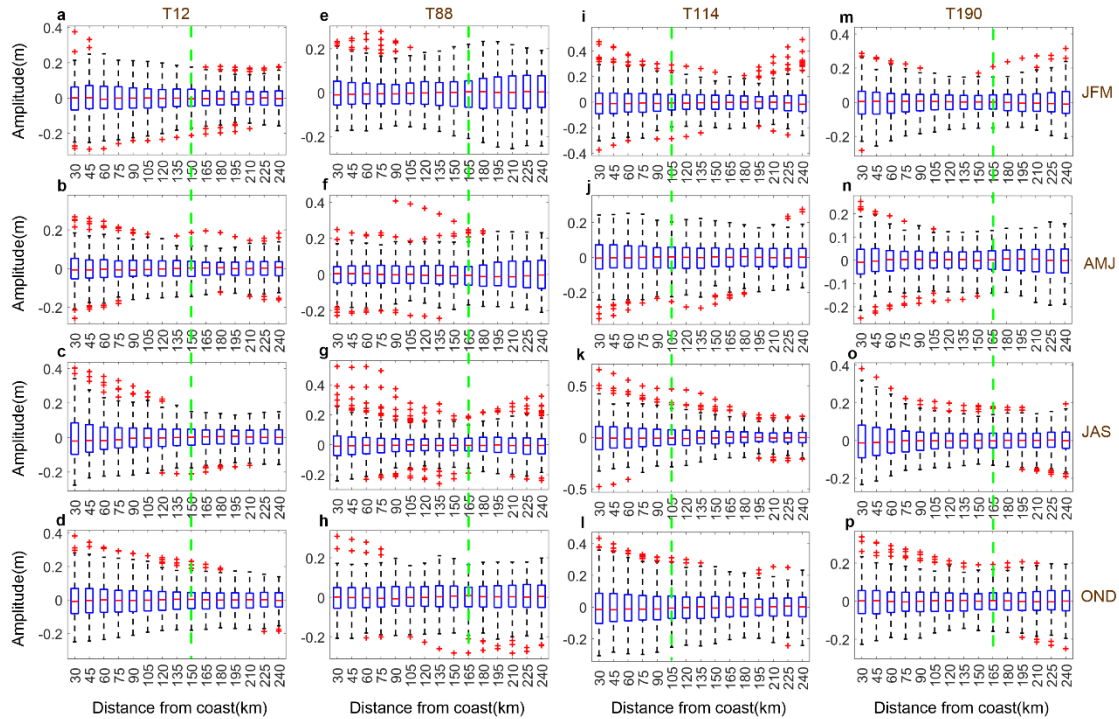
419 5.2 Trapped cross-shelf structure

420 5.2.1 CSWs

421 As the sampling period of the altimeter satellites is about 10 d, the signals with
422 periods shorter than 20 d could not be distinguished from the along-track SLA based
423 on the Nyquist sampling theorem. Fortunately, the cross-shelf structure of CSWs could
424 be sampled as fragments by repeat satellite observations.

425 In this case, the CSWs with a period less than 40 d were abundant from 1993 to
426 2020 on the northern continental shelf of the SCS, especially in winter (Fig. 5). Even if
427 the significance level is larger than 5%, the power is universal and continuous in this
428 period band. Therefore, it could be considered that a large number of repeat
429 observations by altimeter satellites were executed during CSW events.

430 Fig. 8 shows a boxplot of the along-track SLA over the continental shelf for tracks
431 12, 88, 114, and 190. The maxima, minima, and outliers derived from Fig. 8 are listed
432 in Table 2. The whisker dashed black line outside of the box extends to the most extreme
433 data points. One can see that the trapped characteristics as the maximum amplitude
434 (75th percentile), interquartile range (IQR), and outliers (red cross sign) occur at the
435 coastal side. In contrast, the minima occur at the edge of the continental shelf (150, 165,
436 105, and 165 km offshore for tracks 12, 88, 114, and 190). The amplitude of the along-
437 track SLA for maxima and IQR decreases gradually from the coastline (0.2-0.4 m) to
438 the edge of the continental shelf (~0.1 m). The largest outlier is 0.65 m over the shelf
439 30 km offshore.



441

442 **Fig. 8.** Boxplot of the along-track SLA on the shelf for tracks 12, 88, 114, and 190 from
 443 1993 to 2020 (in the column). Seasonal means of SLA are plotted in the rows. Green
 444 dashed lines present the minimum STD of the along-track SLA. The climatological
 445 seasonal mean is removed. The along-track SLA is averaged for each 15 km offshore.
 446 In each box, the central red line indicates the median, and the bottom and top edges of
 447 the blue box indicate the 25th (Q1) and 75th (Q3) percentiles, respectively. The upper
 448 (Q3+1.5IQR) and lower (Q1-1.5IQR) whiskers extend to the most extreme data points
 449 not considered outliers. The outliers are the most extreme data points (larger than upper
 450 whisker or smaller than lower whisker), plotted individually using the red cross marker.
 451 IQR= Q3-Q1.

452

453 Table 2. Parameters extracted from along-track SLA in Fig. 8.

Track		T12 (150 km)				T88 (165 km)				T114 (105 km)				T190 (165 km)			
Position	(Distance from coast (km))	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND
*	e)**																
	Month	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND	JFM	AMJ	JAS	OND
	Outlier(m)	0.28	0.25	0.38	0.34	0.47	-0.35	0.65	0.43	0.23	0.24	0.52	0.31	0.38	0.27	0.41	0.38
Coast	Max (m)	0.26	0.19	0.32	0.24	0.29	0.24	0.42	0.36	0.21	0.17	0.25	0.18	0.21	0.19	0.34	0.28
	Min (m)	-0.23	-0.22	-0.24	-0.23	-0.3	-0.24	-0.33	-0.31	-0.17	-0.18	-0.24	-0.2	-0.25	-0.21	-0.28	-0.25

Edge	Outlier(m)	0.17	-0.15	0.19	0.19	0.23	-	0.29	-	0.21	0.39	0.28	0.16	-0.2	0.2	-0.2	0.21
	Max (m)	0.15	0.12	0.13	0.16	0.16	0.2	0.2	0.22	0.17	0.18	0.15	-0.18	0.16	0.12	0.14	0.17
	Min (m)	-0.15	-0.14	-0.13	-0.15	-0.19	-0.2	-0.2	-0.21	-0.16	-0.14	-0.15	-0.2	-0.17	-0.13	-0.16	-0.17

454 *Coast is the position with the highest STD for each track, about 15 km offshore. Edge
455 is the position with the minimum STD of the along-track SLA as shown in Fig. 8.
456 ** Distance between coastline and edge.

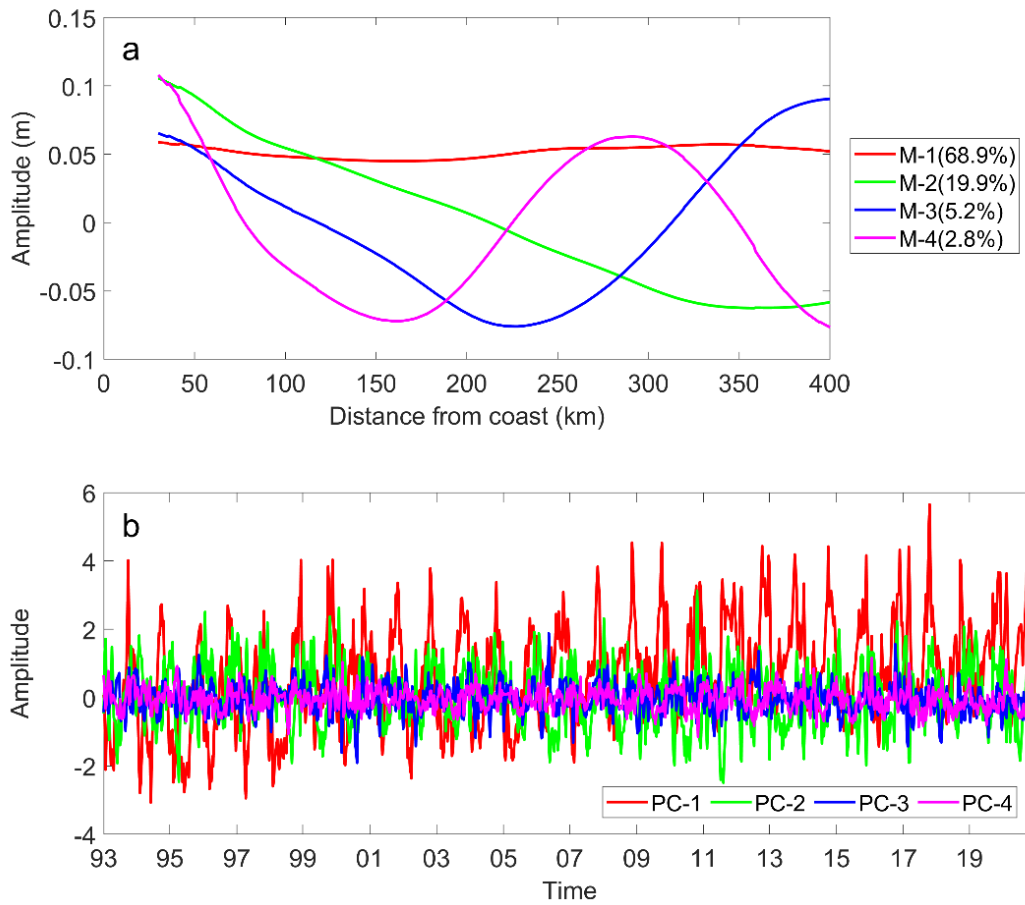
457

458 The largest SLAs over the shelf occur from July to September. For example, the
459 75th percentile of SLAs near the coastline is about 0.25 m for track 144, and 0.42 m for
460 track 88. The largest outlier of 0.65 m also occurs from July to September for track 88.
461 The smallest SLA occurs from April to June when the wind is weak during the monsoon
462 transition period (Wang et al., 2009). The extreme data from July to September show
463 the occurrences of storm surges over the shelf (Chen et al., 2014).

464 Overall, the maximum amplitude, IQR, and extreme data of the along-track data
465 over the shelf show the trapped wave characteristics, which shows that SLA decreases
466 gradually from the coastline to the edge of the continental shelf. The trapped
467 characteristics from the along-track SLA are similar to the cross-shelf structure of
468 normalized amplitudes of the mode 1 CSWs, as shown in Fig. 3a. The along-track SLA
469 should contain higher modes. The along-track SLA along tracks 153 and 229 show
470 similar characteristics (not shown). That along track 240 (as shown in Fig. 6n) presents
471 a differentiated pattern in the coast side and shelf edge during May-July. The main
472 reason should be the existence of cold eddy in the north of Taiwan Island.

473 To further reveal the variations in the along-track SLA in track 12 on the shelf, the
474 empirical orthogonal function (EOF) analysis results are shown in Fig. 9 (EOF analysis
475 for data along other tracks is similar, not shown here). The first four EOF modes of the
476 along-track SLA explain 96.8% of the total variance. Mode 1 explains the seasonal
477 variance of SSH (red curve in Figs. 6o-u). The seasonal cycle could be characterized
478 clearly from Fig. 9b, with peaks in October and troughs in May. Mode 2 and Mode 3
479 are similar to the cross-shelf structure as shown in Fig. 3a, which explains 25.1% of the
480 total variance. Mode 3 is influenced by the background of SSH in the open sea side.

481 The amplitude is 0.3-0.4 m, which is comparable to the outlier data as shown in Table
 482 2. Mode 4 is similar to the cross-shelf structure of mode-2 CSW as shown in Fig. 3b,
 483 which only explains 2.8% of the total variance. The EOF analysis indicates that CSW
 484 could explain <30% of the total variance of SSH in the study area, which is comparable
 485 to the IQR in boxplot of the along-track SLA (Fig. 8).
 486



487
 488 **Fig. 9.** The first four EOFs for the along-track SLA on the shelf along track 12. Cross-
 489 shelf amplitude (a) and time series (b) of EOFs. The variance explained by each mode
 490 is labeled in (a).

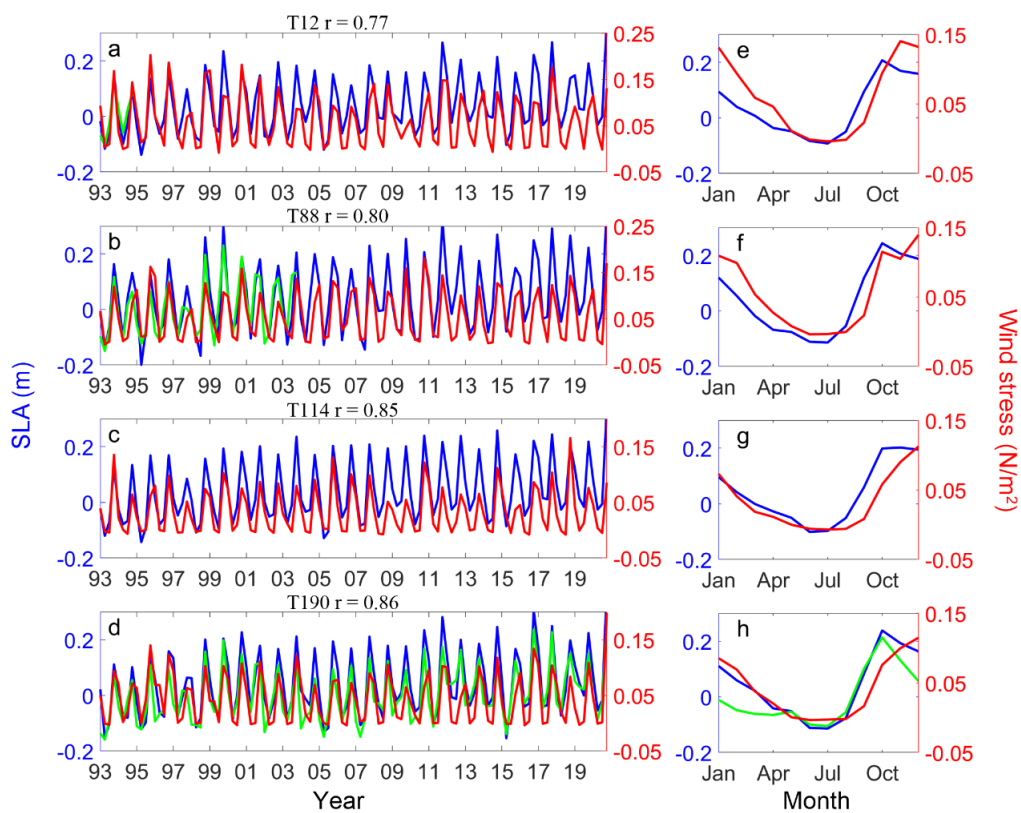
491

492 5.2.2 ATWs

493 As shown in Fig. 6, the inter-seasonal and seasonal signals could be distinguished
 494 from along-track SLA. Fig. 10 shows the seasonal mean of the along-track SLA and
 495 along-shelf sea surface wind. One can see that the time series of the seasonal mean of

496 the along-track SLA show the seasonal variation with the amplitude of 0.1-0.2 m at 15
 497 km offshore. The time series of the seasonal mean of DSLA show similar characteristics
 498 with the along-track SLA. The correlation relationship between the seasonal mean of
 499 the along-track SLA and along-shelf sea surface winds reaches >0.77 . The monthly
 500 mean of SLA during April-September is negative, while positive in the other months. It
 501 is attributed to that the local wind stress substantially influences the coastal sea level
 502 (Lin et al., 2022).

503



504

505 **Fig. 10.** Time series of the seasonal mean of the along-track SLA and along-shelf sea
 506 surface wind stress for tracks (a)12, (b) 88, (c) 114, (d)190, and (e-h) climatological
 507 monthly mean of the along-track SLA and along-shelf sea surface wind stress. Red
 508 curves represent the along-track SLA at 15 km offshore. Green curves represent the
 509 seasonal mean of sea level data at tide-gauge stations (a) Xiamen, (b) Shanwei, and (d)
 510 Zhapo. Green curve in (h) is climatological monthly mean of sea level data at tide-
 511 gauge station Zhapo.

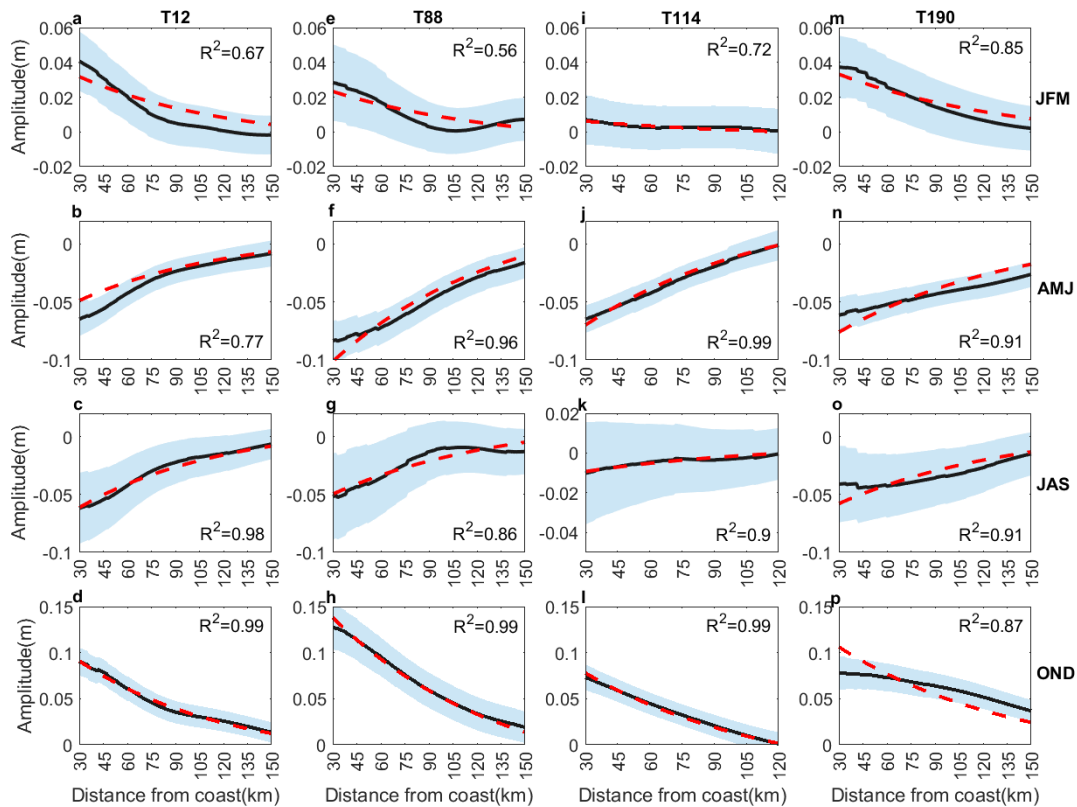
512

513 In addition, one can also see a characteristic of out-of-sync between the seasonal
 514 mean of the along-track SLA and wind. The maximum mean sea surface wind stress
 515 occurs in November and December. While, that of monthly along-track SLA occurs in

516 October. Ding et al. (2020) investigated the seasonality of coastal circulation in the
 517 north SCS using a numerical model. The result indicates that the maximum of water
 518 transport is 1.93 Sv occurred in autumn. Li et al. (2023a) found the along-shelf current
 519 is strongest in October at approximately 0.17 m s^{-1} . The along-shelf current involves
 520 ATW, which is the reason why there is a difference between monthly climatological
 521 mean of the along-track SLA and sea surface wind stress as shown in Figs. 10e-h.

522 Fig. 11 shows the cross structure of the seasonal mean of the along-track SLA.
 523 One can see that the SLA on the coastline side is lower than on the ocean side from
 524 April to September. In the other seasons, the slope of the along-track SLA is the
 525 opposite. The along-track SLA presents similar characteristics as shown in Fig. 10, i.e.,
 526 the variation of SLA is controlled by the sea surface winds over the shelf.

527



528

529 **Fig. 11.** Fitting climatologic seasonal mean of the along-track SLA (black curves) using
 530 cross-shelf ATWs (red dashed curves) for tracks (a-d) 12, (e-h) 88, (i-l) 114, and (m-p)
 531 190. Dashed shadow is the STD of the seasonal mean of the along-track SLA.

532

533 Meanwhile, from Fig.11 one can see the cross-shelf structure of the along-track

534 SLA. The fitting curves show that theoretical ATWs (Fig. 4) explain the cross-shore
535 structure of the along-track SLA very well. The amplitudes of ATWs in track 12 are
536 0.04 m, -0.06 m, -0.05 m, and 0.10 m in January-March, April-June, July-September,
537 and October-December. That along track 88 is relatively larger, e.g., 0.13 m in October-
538 December. While, the minimum amplitude occurs in track 114. Differently from Fig. 8,
539 the amplitudes of ATWs during October-December are larger than that during July-
540 September. It should be attributed that the monsoon winds in winter are stronger than
541 that in summer.

542 Lin et al. (2021) investigated the tilt of mean dynamic topography along the coast
543 of the Chinese mainland. Wu (2021) used a nondimensional parameter ($Pe_\beta = D_\beta/\alpha$) to
544 describe the influence of open ocean forcing on shelf circulation, which is determined
545 by the ratio of long-wave-limit planetary to TRW speeds (D_β) and linear Ekman
546 number (α). In this study, $Pe_\beta < 1$, which indicates shelf currents decayed rapidly toward
547 the coast. Their results confirm that the ATWs predict the coastal dynamic topography
548 over the continental shelf of the SCS well.

549

550 **6 Summary**

551 Using sea level data derived from the tide-gauge stations Kanmen, Xiamen,
552 Shanwei, Hongkong, and Zhapo, this study analyzes statistical features of the CSWs,
553 inter-seasonal and seasonal signals. Meanwhile, along-track SLA data derived from
554 multiple satellite altimeters from 1993 to 2020 are applied to detect the cross-shelf
555 structures of the signals. The major results are summarized as follows.

- 556 1) CSWs of periods shorter than 40 d propagate along the coast with a phase speed of
557 about 10 m s^{-1} in the ECS and 8 m s^{-1} in the SCS. The dispersion relation indicates
558 that the waves belong to mode 1 CSWs.
- 559 2) Owing to the fact that the repeat observation period of the satellite altimeters is
560 comparable with that of CSWs, we combine fragments of the numerous repeat
561 observations of the along-track SLA to reconstruct the cross-shelf structure of
562 CSWs. The results show that the maximum amplitudes of CSWs have remarkable
563 seasonal variability, about 0.6 m during July-September, while only 0.2 m during

564 April-June. The reconstructed cross-shelf structures of CSWs confirm the property
565 of mode 1 CSWs. Moreover, the energy is trapped within the partial continental
566 shelf shallower than 200 m.

567 3) The inter-seasonal and seasonal signals present as ATWs on the continental shelf.
568 The amplitudes of ATWs have remarkable seasonal variability, ~ 0.10 m during
569 October-December, twice larger than 0.04 m, 0.05 m, and -0.06 m during January-
570 March, July-September, and April-June, respectively. These results reveal that the
571 local wind stress substantially influences ATWs on the continental shelf.

572 4) The results derived from the observation data of the along-coast tide-gauge stations
573 combined with cross-shelf tracks of satellite altimeters are interpreted well by the
574 framework of linear wave theory. It implies that the technological approaches
575 developed in this study are suitable for constructing the cross-shelf structures of
576 CSWs and ATWs on the continental shelf.

577 However, owing to the neglective wind stress and baroclinicity, higher modes of
578 waves are not discussed in this paper. Observations from moorings and numerical
579 models will be used in our future studies to obtain the characteristics of baroclinic
580 coastal trapped waves.

581

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589

590 **Data Availability Statement**

591 The tide gauge data are available at <https://psmsl.org/data/obtaining/>. The along-track
592 SLAs are obtained at <ftp-access.aviso.altimetry.fr>.

593

594 **Author contributions**

595 JYL were responsible for writing the original draft. Review and editing were conducted
596 by QAZ. Conceptualization was handled by JYL, QAZ and LLX. TH and YX were
597 responsible for data curation. LLX acquired funding.

598

599 **Competing interests**

600 The contact author has declared that none of the authors has any competing interests.

601

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