1	Statistical analysis of dynamic behavior of continental shelf wave		
2	motions in the northern South China Sea		
3	Junyi Li ^{1,2,3} , Tao He ^{1,3} , Quanan Zheng ^{1,4} , Ying Xu ³ , Lingling Xie ^{1*}		
4	¹ Laboratory of Coastal Ocean Variation and Disaster Prediction, College of Ocean and		
5	Meteorology, Guangdong Ocean University, Zhanjiang 524088, China		
6	² Key Laboratory of Climate, Sources and Environments in Continent Shelf Sea and		
7	Deep Ocean, Zhanjiang 524088, China		
8	³ Key Laboratory of Space Ocean Remote Sensing and Application, MNR, Beijing,		
9	100081, China		
10	⁴ Department of Atmospheric and Oceanic Science, University of Maryland, College		
11	Park, MD 20742, USA		
12	*Corresponding author.		
13	E-mail address: L. Xie (xiell@gdou.edu.cn)		
14			
15	Abstract		
16	This study aims to analyze statistical behavior of the continental shelf wave		

motions, including continental shelf waves (CSWs) and arrested topographic wave 17 (ATW), in the northern South China Sea. The baseline consists of tide-gauge data from 18 stations Kanmen, Xiamen, Shanwei, Hongkong, and Zhapo, as well as along-track sea 19 level anomaly (SLA) data derived from multiple satellite altimeters from 1993 to 2020. 20 The subtidal signals propagating along the coast with periods shorter than 40 d and 21 phase speeds of about 10 m s⁻¹ are interpreted as CSWs. The cross-shelf structure of 22 23 along-track SLA indicates that mode 1 of CSWs is the predominant component trapped in the area shallower than about 200 m. The amplitudes of CSWs reach the maximum 24 0.6 m during July-September, and minimum 0.2 m during April-June. The inter-25 seasonal and seasonal signals represent ATWs. The amplitudes of ATWs reach 0.10 m 26 during October-December, twice of that during July-September. These observations are 27 well interpreted by the framework of linear wave theory. The cross-shelf structures of 28 CSWs and ATWs derived from along-track SLA illustrate that the methods are suitable 29

- 30 for observing dynamic behavior of the CSWs.
- 31

Keywords: Continental shelf wave; Arrested topographic wave; Cross-shelf SLA
structure; South China Sea

34

35 **1 Introduction**

Continental shelf wave (CSW) is a type of topographic Rossby wave (TRW) 36 trapped in the continental shelf with amplitudes ranging from several tens' centimeters 37 to more than one meter (Aydın and Beşiktepe, 2022; Clarke and Brink, 1985; Heaps et 38 al., 1988; Morey et al., 2006; Mysak, 1980; Robinson, 1964; Zheng et al., 2015). CSW 39 is a sub-inertial motion with a wavelength much greater than the depth (Li et al., 2015; 40 41 Schulz et al., 2011). It propagates along the shelf with the coast on its right (left) in the northern (southern) hemisphere (Clarke, 1977). During the impact of typhoon, an 42 excessive flooding in the coastal zone could be induced by a propagating CSW that 43 added to the locally wind-generated surge (Dukhovskoy and Morey, 2011; Han et al., 44 45 2012). Therefore, CSW is particularly important for coastal sea-level variations.

CSW is generally generated by large-scale weather systems moving across or 46 along the shelf (Thiebaut and Vennell, 2010). CSW events have been reported by 47 previous investigators lasted from 2 days to 2 weeks (Chen and Su, 1987; Li et al., 2015; 48 Li et al., 2021; Zheng et al., 2015). The phase speed of CSWs depends on the bottom 49 topography, ranging from 5 to 20 m s⁻¹ (Li et al., 2015; Li et al., 2016; Shen et al., 2021). 50 CSWs could be taken as barotropic motion in a homogeneous coastal area. While in a 51 stratified ocean, it could be classified into coastal trapped wave. If the bottom boundary 52 53 is flat, it propagates as a Kelvin wave. Overall, they are resulting from conserving potential vorticity over the shelf (Chen et al., 2022; Quan et al., 2021; Wang and Mooers, 54 1976). 55

The sea level variations in the South China Sea (SCS) are well depicted by these previous studies as the continental shelf occupies about its half area (Ding et al., 2012; Li et al., 2023b; Shen et al., 2021; Zhao et al., 2017; Zhou et al., 2023). However, two issues should be improved. The first is that the primary data are usually obtained from

the tide-gauge stations along the coastline. The data have high accuracies, but represent 60 the sea level at the coasts only. Thus, satellite-altimeter-observed the sea level variation 61 62 are often used to fill the data gaps in between the tide-gauge stations and on the continental shelf. The second issue is that the repeated period of the satellite altimeter 63 is 9.9 d, which is challenging to investigate the sea level variations with periods shorter 64 than 10 d. Previous studies used the along-track sea level anomaly (SLA) from satellite 65 altimeters to describe CSW (Chen et al., 2014; Li et al., 2016). However, the satellite 66 altimeter with the sparse tracks (as shown in Fig. 1) could only capture the cross-67 structure of CSW with one or two snapshot of one CSW. 68



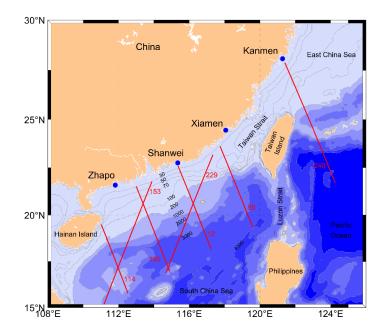




Fig. 1. Study area. Blue dots represent locations of tide-gauge stations Kanmen,
Xiamen, Shanwei, and Zhapo. Red lines represent segments of ground tracks 12, 88,
114, 153, 190, 229, and 240 for altimeter satellites over the continental shelf. Tracks 12,
88,114, and 190 are almost perpendicular to the coastline. Isobaths are in m.

75

Under the background of global warming, the sea level variation with very low frequency has been investigated by previous investigators. Ho et al. (2000) found seasonal sea level variability in the SCS using data from a satellite altimeter. Kajikawa and Yasunari (2005) investigated the interannual variability of the intra-seasonal variation over the SCS. Fang et al. (2006) analyzed low sea level along the eastern boundary of the SCS. Rong et al. (2007) investigated the relationship between ENSO
(El Niño and Southern Oscillation) and interannual variability of sea level in the SCS.
Zhuang et al. (2010) found strong intra-seasonal variability in the northern SCS. In
addition, the sea level variations are influenced by thermodynamic processes, e.g.,
eddies and thermal change of upper layer of the SCS (Cheng and Qi, 2007; Xie et al.,
2018; Zheng et al., 2014).

Meanwhile, the upper layer thermal changes significantly influence the sea level 87 variations (Cheng and Qi, 2007). Wang et al. (2017) found that seasonal level anomalies 88 are closely related to ENSO events (Wang et al., 2022). Using sea surface height (SSH) 89 data from satellite altimeters, Xu et al. (2016) found that sea level variations in the 90 coastal area of the SCS are still strongly influenced by the coastal current system in 91 summer and winter. Seasonal circulation is mainly driven by the monsoon wind stress 92 (Gan et al., 2006). Lin et al. (2021) applied the arrested topographic waves (ATWs) 93 model to the coastal mean dynamic topography along the East China Sea (ECS) and 94 SCS. The mean circulation in a coastal zone of variable depth may be modeled by linear 95 96 equations (Wu, 2021). The result suggests that the mean dynamic topography is a counterbalance of contributions from the along-shelf wind and well predicted by the 97 ATW model. Therefore, the monsoon winds are a control factor for the sea level 98 variation. 99

This study aims to investigate the cross-shelf structures of sea level over the continental shelf. As the repeated period of the satellite altimeters, 9.9 d, is comparable with that of CSWs in the northern SCS, the statistical characteristics of the along-track SLA are applied to show the cross-shelf structure of CSWs using a long-term data set from 1993 to 2020. To figure out the cross-shelf structures of ATWs is another goal.

The rest of the paper is organized as follows: Section 2 describes the observed data. Section 3 presents theory of CSWs and ATWs. Section 4 presents the characteristics of the signals derived from the tide-gauge data and along-track SLA. Section 5 discusses the CSWs detected from tide-gauge data and the cross-shelf structure of sea level over the continental shelf. Section 6 gives summaries.

111 **2 Data**

112 2.1 Along-track sea level anomaly data

Satellite altimeter along-track SLA data are produced and distributed by the 113 Archiving, Validation, and Interpretation of Satellite Oceanographic Data (AVISO), 114 Centre National d'Etudes Spatiales (CNES) of France. The data from 1993 to 2020 are 115 derived from TOPEX/Poseidon, Jason-1, Jason-2, and Jason-3 measurements. The 116 satellite repetition period is 9.9 d, and the temporal resolution of the along-track data is 117 1 Hz. The along-track SLA is calculated by subtracting the twenty-year mean from the 118 SSH measured by the satellite altimeters. The along-track SLA is low pass filtered using 119 7-point moving average. The ground tracks in the study area, 12, 88, 114, 153, 229, and 120 240, are shown in Fig. 1. 121

Satellite altimetry provides a unique sea level dataset to the coastal sea level research. A few recent studies have stressed the importance of small-scale coastal processed on coastal sea-level variance (Cazenave and Moreira, 2022; Vignudelli et al., 2019). The along-track SLA has been successfully validated and applied to the coast zone by Birol et al. (2021). These studies present the availability of along-track SLA in the coastal zones.

128

129 2.2 Sea level anomaly from tide-gauge

The tide-gauge data at stations Kanmen, Xiamen, Shanwei, Hongkong and Zhapo
(as shown in Fig. 1) are obtained from the Global Sea Level Observing System
(GLOSS). The data cover a period from 1993 to 1997, with a temporal resolution of 1
h. De-tided sea level anomaly (DSLA) is calculated by removing tidal signals using a
Matlab toolbox (Pawlowicz et al., 2002).

The monthly sea level means at stations Xiamen, Shanwei, and Zhapo are obtained from the Permanent Service for Mean Sea Level (PSMSL). Monthly mean data cover periods of 1993-2003, 1993-1994, and 1993-2020 for stations Xiamen, Shanwei, and Zhapo, respectively.

139

140 2.3 Sea surface wind stress

Monthly sea surface wind stress is derived from the Copernicus Marine Environment Monitoring Service (CMEMS). The dataset covers a period from 1993 to 2020 with a spatial resolution of $0.25^{\circ} \times 0.25^{\circ}$. Sea surface wind stress data on the satellite altimeter ground tracks are decomposed into the cross-shelf and along-shelf components. The cross-shelf component is positive seaward and parallel to the satellite altimeter ground tracks (12, 88, 114 and 190). The along-shelf component is positive northward and perpendicular to these satellite altimeter ground tracks.

148

149 2.4 Topographic profile

The behavior of continental shelf waves is determined by the topography of the 150 continental shelf and slope, which has been well documented. The topographic profiles 151 along the satellite altimeter ground tracks are extracted from a dataset of ETOPO-2. 152 This study uses one-dimensional linear piecewise functions to fit the topographic 153 profiles along the satellite altimeter ground tracks. The width of the continental shelf, 154 depths of shelf break and deep basin along the tracks are listed in Table 1. The 155 156 continental shelf break is extracted as the location of maximum change in the gradient of continental slope. 157

158

Table 1. The width of the continental shelf, depths of shelf break and deep basin alongthe satellite ground tracks.

Track number	<i>l</i> (km)	$H_1(m)$	$H_2(m)$
12	200	-300	-3800
88	178	-200	-2500
114	123	-200	-1800
153	259	-110	-1600
190	247	-255	-3500
229	244	-225	-3500
240	273	-250	-5000

161

162 **3 Theory**

163 3.1 Momentum equation for CSWs

164 The linearized shallow-water equations governing a barotropic ocean on a rotating 165 earth are

$$\frac{\partial u}{\partial t} - fv = -g \frac{\partial \eta}{\partial x} + \frac{\tau_s^x - \tau_b^x}{\rho H}$$
(1a)

167
$$\frac{\partial v}{\partial t} + fu = -g \frac{\partial \eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H}$$
(1b)

166

168
$$\frac{\partial \eta}{\partial t} + \frac{\partial (uH)}{\partial x} + \frac{\partial (vH)}{\partial y} = 0$$
 (1c)

where cross-shelf and along-shelf velocities (u, v) are depth-averaged in cross-shelf and along-shelf coordinates (x, y). η is the sea surface height. The Coriolis parameter is f. The bathymetry H=H(x) is assumed to be a function of the cross-shelf variable, x only. τ_s and τ_b the surface and bottom stresses. g, ρ are the gravitational acceleration and the water density.

The scales of the along-shelf length of CSW ($L = 2\pi/k \approx 2 \times 10^3$ km), and cross-shelf length ($l \approx 200$ km) are subject to the long-wave assumptions ($l/L \ll 1$), i.e., $\partial u/\partial t = 0$ (Li et al., 2016; Schulz et al., 2011). Under this approximation, Eqs. (1a-b) become

178
$$v = \frac{g}{f} \frac{\partial \eta}{\partial x} - \frac{\tau_s^x - \tau_b^x}{f\rho H}$$
(2a)

179
$$u = -\frac{g}{f}\frac{\partial\eta}{\partial y} - \frac{\partial}{\partial t}\left(\frac{g}{f^2}\frac{\partial\eta}{\partial x} - \frac{\tau_s^x - \tau_b^x}{f^2\rho H}\right) + \frac{\tau_s^y - \tau_b^y}{f\rho H}$$
(2b)

180 Substitute Eqs. (2a-b) into Eq. (1c), we obtain the equation governing SSH of 181 CSWs

182
$$\frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y - \tau_b^y}{f\rho} \right) + \frac{\partial^2}{\partial t \partial x} \left(\frac{\tau_s^x - \tau_b^x}{f^2 \rho} \right) - H \frac{\partial}{\partial y} \left(\frac{\tau_s^x - \tau_b^x}{f\rho H} \right) = 0$$
183 (3)

Assume that CSWs are forced by along-shelf wind stress, i.e., $(\tau_s^x, \tau_s^y) = (0, \tau_s^y)$. The bottom friction is neglected to simplify the calculation in Eq. (3). Therefore, the equation governing the SSH of CSWs becomes

187
$$\frac{\partial}{\partial t} \left(\eta - \frac{g_H}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} + \frac{\partial}{\partial x} \left(\frac{\tau_s^y}{f\rho} \right) = 0 \tag{4}$$

188 The change in the SSH of CSWs is balanced by the variation of along-shelf wind 189 stress in cross-shelf direction. Assume a periodic along-shelf wind stress, $\tau_s^y =$ 190 $\tau_0 \exp[i(\alpha y + \omega t)]$ (where α is wavenumber of wind stress, and $\tau_0 = \text{constant}$). Eq. 191 (4) becomes

192

$$\frac{\partial}{\partial t} \left(\eta - \frac{gH}{f^2} \frac{\partial^2 \eta}{\partial x^2} - \frac{g}{f^2} \frac{\partial \eta}{\partial x} \frac{dH}{dx} \right) - \frac{g}{f} \frac{\partial \eta}{\partial y} \frac{dH}{dx} = 0$$
(5)

193 which means that the change in SSH is independent of along-shelf wind stress.

194 The assumption of a wave solution

195
$$\eta = \phi(x) \exp[i(ky + \omega t)]$$
(6)

196 yields an equation for $\phi(x)$,

$$H\frac{d^2\phi}{dx^2} + \frac{dH}{dx}\frac{d\phi}{dx} + \left(\frac{fk}{\omega}\frac{dH}{dx} - \frac{f^2}{g}\right)\phi = 0$$
(7)

198

197

199 3.1.1 Over the shelf

200 For
$$0 \le x \le l$$
, $H = H_1 x/l$, the equation for $\phi(x)$ is

201
$$x\frac{d^2\phi}{dx^2} + \frac{d\phi}{dx} + \left(\frac{fk}{\omega} - \frac{f^2l}{gH_1}\right)\phi = 0$$
(8)

which is subject to the following boundary conditions: $\phi(0) = a$, and $\phi(l) = A$. *a* and *A* should be arbitrary. We could take *a* as the amplitude of fluctuation in sea level from the tidal gauge station.

The solution to Eq. (8) is expressed as the sum of the first and second kinds of Bessel functions (Robinson, 1964; Schulz et al., 2011)

207
$$\phi(x) = aJ_0 \left(2\left(\frac{fk}{\omega} - \frac{f^2l}{gH_1}\right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) + bY_0 \left(2\left(\frac{fk}{\omega} - \frac{f^2l}{gH_1}\right)^{\frac{1}{2}} x^{\frac{1}{2}} \right)$$
(9)

where *a* is arbitrary constant. As the solution for $\phi(x)$ is finite, therefore

209
$$\phi(x) = a J_0 \left(2 \left(\frac{fk}{\omega} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right)$$
(10)

210

211 3.1.2 In the deep basin

For
$$l < x$$
, $H = H_2$, the equation for $\phi(x)$ is

213
$$\frac{d^2\phi}{dx^2} - \frac{f^2}{gH_2}\phi = 0 \tag{11}$$

which is subject to the following boundary conditions: $\phi(x \to l) = A$, and $\phi(\infty) = 0$.

The solution is

216
$$\phi(x) = A \exp\left(-\frac{fl}{\sqrt{gH_2}}\left(\frac{x}{l}-1\right)\right) + B \exp\left(\frac{fl}{\sqrt{gH_2}}\left(\frac{x}{l}-1\right)\right)$$
(12)

As the solution for $\phi(x)$ is also finite in the deep basin, i.e.,

218
$$\phi(x) = A \exp\left(-\frac{fl}{\sqrt{gH_2}}\left(\frac{x}{l} - 1\right)\right)$$
(13)

219 where *A* is arbitrary constant.

220

221 3.1.3 Dispersion relation

As the fluid is continuous at the edge of continental shelf, therefore

223
$$H_1 \cdot u|_{x \to l_-} = H_2 \cdot u|_{x \to l_+}$$
(14a)

224
$$aJ_0\left(2\left(\frac{fkl}{\omega} - \frac{f^2l^2}{gH_1}\right)^{\frac{1}{2}}\right) = A$$
(14b)

225 From Eq. (14a), we have

226
$$H_1 \cdot a\left(kJ_0 + \frac{\omega}{f}\left(\frac{fk}{\omega l} - \frac{f^2}{gH_1}\right)^{\frac{1}{2}}J_0'\right) = H_2 \cdot A\left(k - \frac{\omega}{\sqrt{gH_2}}\right)$$
(15)

where J_0' is the derivation of the zero-order of the first kind Bessel function,

228
$$J'_0(x) = -\frac{x}{2} (J_0(x) + J_2(x))$$
(16)

- 229 where J_2 is the second order of J_0 .
- Substituting Eq. (16) into Eq. (15) yields a dispersion relation for CSWs

231
$$\left(\frac{cfl}{gH_2} - 1 + \frac{c}{(gH_2)^{\frac{1}{2}}}\right) J_0 - \left(\frac{H_1}{H_2} - \frac{cfl}{gH_2}\right) J_2 = 0$$
(17)

where $c \ (= \frac{\omega}{k})$ is phase speed of CSWs. We solve Eq. (17) using the zero-finding function in MATLAB. The solution of phase speed for CSWs is shown as Fig. 2. The zero-crossing points for each curve present the phase speed of CSWs. The first zerocrossing point on the right hand of each curve points out the phase speed of mode-1 CSW, e.g., c = 7.6 m s⁻¹ for track 190.

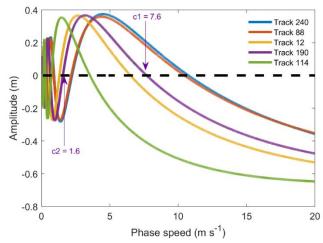


Fig. 2. Solution of phase speed for CSWs from Eq. (17). The zero-crossing values 239 represent the phase speed of CSWs for different modes. Colorful curves are amplitude 240 of zero-order of the first kind Bessel function in different modes for the idealized depth 241 profile of tracks 12, 88, 114, 190, and 240, respectively. Arrows point out the phase 242 speed of mode-1 (7.6 m s⁻¹) and mode-2 (1.6 m s⁻¹) CSWs along track 190. 243

244

3.1.4 Cross-shelf structure 245

With Eqs. (10), (13) and (14b), the wave solution of SSH, i.e., Eq. (3), is 246

$$\eta = \begin{cases} \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{fk_i}{\omega_i} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) \exp[i(k_i y + \omega_i t)] & x \le l \end{cases}$$

247
$$\eta = \begin{cases} \sum_{i=1}^{\infty} a_i \cdot J_0 \left(2 \left(\frac{fk_i l}{\omega_i} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) \exp[i(k_i y + \omega_i t)] & x > l \end{cases}$$
248 (18)

248

where $\exp[i(ky + \omega t)]$ is the waveform propagating along the shelf. $a_i \cdot J_0$ is the 249 cross-shelf structure of the waveform for mode *i* over the shelf. 250

In the cross-shelf direction, the SSH becomes 251

252
$$\eta = \begin{cases} \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{fk_i}{\omega_i} - \frac{f^2 l}{gH_1} \right)^{\frac{1}{2}} x^{\frac{1}{2}} \right) & x \le l \\ \\ \sum_{i=1}^{\infty} \eta_{0i}(t) \cdot J_0 \left(2 \left(\frac{fk_i l}{\omega_i} - \frac{f^2 l^2}{gH_1} \right)^{\frac{1}{2}} \right) \exp \left(-\frac{fl}{\sqrt{gH_2}} \left(\frac{x}{l} - 1 \right) \right) & x > l \end{cases}$$
(19)

where $\eta_{0i}(t) (= a_i \exp[i(k_i y + \omega_i t)])$ is time series of SSH at the coastline for mode 253 i. 254

Fig. 3 shows the SSH evolution over the continental shelf. The SSH mode 1 looks 255 like a bell mouth. For Modes 2 and 3, nodes and antinodes appear on the shelf. Node 256 for Mode 2 appears at 50 km off the coast over the shelf. Nodes for Mode 3 appear at 257

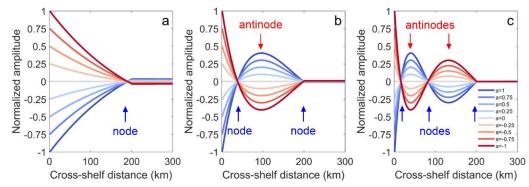


Fig. 3. (a-c) Cross-shelf structure of normalized amplitudes of the first three CSW
modes for idealized depth profile of track 12. Gradient color curves represent amplitude
evolution over time for CSWs in the cross-shelf direction.

260

265 3.2 Steady situation

The governing equations in the steady situation for Eqs. (1a-c) are

267
$$-fv = -g\frac{\partial\eta}{\partial x}$$
(20a)

$$fu = -g\frac{\partial\eta}{\partial y} + \frac{\tau_s^y - \tau_b^y}{\rho H}$$
(20b)

268

$$\frac{\partial(uH)}{\partial x} + \frac{\partial(vH)}{\partial y} = 0$$
(20c)

A linear drag is used for the bottom friction, i.e., $\vec{\tau_b} = \rho \lambda \vec{u}$ (Hsueh and Pang, 1989; Lin and Yang, 2011). The SSH Eqs. (20a-c) becomes (Csanady, 1978)

272
$$\eta = B\tau_0 e^{-\frac{x}{L}} \sin\left(\alpha y + \frac{x}{L} - \frac{\pi}{4}\right)$$
(21)

where $L = \sqrt{2\lambda l/\alpha f H_1}$ is the scale width of trapped sea level in the cross-shelf direction, *B* is an arbitrary constant. A drag coefficient, $\lambda = O(5 \times 10^{-4})$ m s⁻¹, is used for the linear bottom friction (Chapman, 1987; Lin and Yang, 2011). The typical magnitude for winter wind stress over the northern continental shelf of the SCS is O(0.1) N m² (Lin and Yang, 2011; Lin et al., 2011). $f = 5 \times 10^{-5}$ s⁻¹, and $\alpha =$ $2\pi/4000$ km evaluated by Lin et al. (2021).

The normalized SSH in the cross-shelf direction of ATW for tracks 12, 88, 114, and 190 are shown in Fig. 4. One can see that the trapped sea level in the cross-shelf direction decays quickly from 1 at the coastline to ~ 0.2 at the edge of the continental shelf (~200 km), and 0.1 at a distance of 300 km. The ATW amplitude decays offshore with a scale equal to the deformation radius, and L = ~100 km in the study area, which is much less than the local Rossby radius of deformation (~600 km). Under different wind stresses, the amplitude of ATW evolves similarly to that in Fig. 3a. As Track 240 is not perpendicular to the coastline, it is beyond the scope.



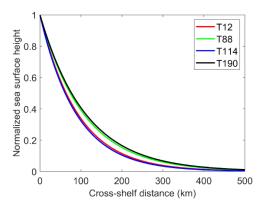




Fig. 4. Normalized SSH in cross-shelf direction for tracks 12, 88, 114, and 190.

- 290
- 291 **4 Signals in sea level anomaly**
- 292 4.1 Tide-gauge data

Fig. 5 shows the wavelet transform (WT) of SLA at station Xiamen from 1993 to 293 1997. One can see abundant signals with periods from several days to one year in Figs. 294 5a-b. A dozen signals with periods from several days to about one month could be seen 295 from the WT of SLA every year. The duration of these signals, with a period of several 296 days, is about ten days (e.g., in Fig. 5d). Moreover, the signals with a period of about 297 one month sustain for about three months (Figs. 5c-d). Even if significance level against 298 red noise is larger than 5%, the power is universal and continuous in the period bands 299 300 of 10-60 d. The main reason is that these signals are not so significant compared with the signals with large amplitude. The wavelet analysis also exhibits a significant inter-301 seasonal and seasonal variation of SLA. The characteristics of the signals at stations 302 Kanmen, Shanwei, and Zhapo are almost the same (not shown here). Therefore, the 303 variation of SLA along the northern SCS coast is universal. 304

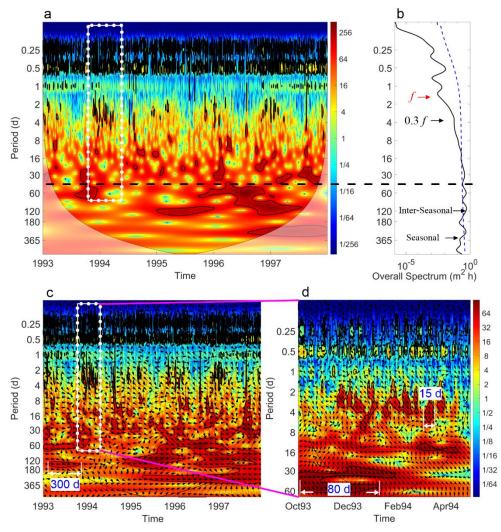


Fig. 5. Temporal variability of SLA at the coastline. (a) WT of SLA at station Xiamen. 307 (b) Overall spectrum of SLA at station Xiamen. (c) XWT of SLA between stations 308 Kanmen and Xiamen. (d) XWT of SLA between stations Kanmen and Xiamen from 309 310 October 1993 to April 1994. White line and arrows are auxiliary lines pointed out duration. The thick line is a 5% significance level against red noise, and the cone of 311 influence (COI) is shown as the thin line. Color codes of power spectra normalized by 312 variance are in arbitrary units. White dotted rectangles in (a) and (c) show the temporal 313 domain of (d). The Blue dashed curve in (b) is a 5% significance level. Arrows in (b) 314 point out the characteristic frequencies of inertial oscillation and CSWs. The black 315 dashed lines in (a) and (b) represent signal period boundary (40 d). The arrows in (c) 316 and (d) show the relative phase relationship between SLA at Kanmen and Xiamen with 317 in-phase (anti-phase, leading, and lagging) pointing right (left, down, and up). 318

319

Figs. 5c-d show the cross wavelet transform (XWT) of SLA between stations Kanmen and Xiamen. One can see that the period band and the occurrence time of the significant cross wavelet power are consistent with that in Fig. 5a. The signals with periods shorter and longer than 40 d show remarkably different characteristics. In the period band shorter than 40 d, one can see the arrows point bottom left uniformly. The uniform phase lag indicates the fixed time delay of the signal between two tidal gauge station. While, the direction of arrows is in disorder in the lower period band, which indicates there is no evidence for propagating.

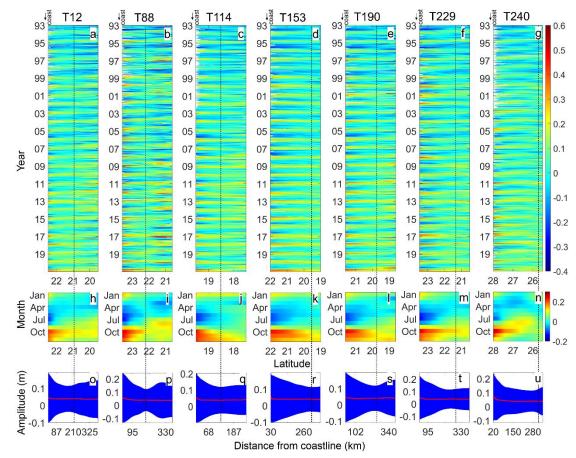
In the period band shorter than 40 d, the signals at station Xiamen lag that at station Kamen about 15 h. The propagation phase speed of the sea level signal could be calculated by the lag time of sea level propagation between stations Kanmen and Xiamen. The result is about 9 m s⁻¹, which is very close to that reported by a number of recent studies (Ding et al., 2012; Li et al., 2015; Li et al., 2016; Zhao et al., 2017).

In the period band longer than 40 d, the phase of signals between Kanmen and Xiamen is a little complicated. The signal phase with the period of 360 d from 1995 to 1997 is almost 0. However, the seasonal signals at station Xiamen lag (lead) that at station Kanmen in winter (summer) about $\pi/4-\pi/2$, implying that the signal propagates very slowly along the coast. Csanady (1978) concluded that these signals are a kind of ATWs.

339

340 4.2 Along-track SLA

Fig. 6 shows the latitude distribution of along-track SLA over the northern 341 continental shelf of the SCS from 1993 to 2020. One can see the clear annual signals 342 near the coast in Figs. 6a-g. Signals with periods shorter than one year could also be 343 seen. As shown in Figs. 6h-n, a clear seasonal cycle is discernable from the 344 climatological monthly mean of along-track SLA. One can see that lower (higher) sea 345 levels over the shelf exist from March to August (September to February). The trough 346 (~-0.2 m) and peak (~0.24 m) of sea level near the coast occur in July and October, 347 respectively. While in track 240, climatological monthly mean of along-track SLA on 348 the shelf is smaller than that in track 88 especially in July. The 28-year mean value of 349 along-track SLA is about 0.04 m, as shown in Figs. 6o-u. 350



352

353 Fig. 6. (a-g) Latitudinal distribution of along-track SLA over the northern continental shelf of the SCS from 1993 to 2020 for tracks 12, 88, 114,153, 190, 229, and 240, 354 respectively. (h-n) Climatological monthly mean of along-track SLAs for Tracks 12, 88, 355 114,153, 190, 229, and 240. (o-u) Mean (red curves) and standard deviation (blue 356 shadow) of along-track SLA from 1993 to 2020 for tracks 12, 88, 114,153, 190, 229, 357 and 240. Track numbers are shown on the tops of panels. Vertical dashed lines represent 358 359 shelf break positions along the tracks, i.e., H₁, listed in Table 1. The coastline position is marked by an arrow on top left of each panel. 360

Moreover, the standard deviations (STD) of along-track SLA from 1993 to 2020 (blue shadow in Figs. 60-u) show a bell-mouth-like structure over the shelf. The amplitudes of along-track SLA reach the maxima near the coast. The minimum variance exists near the shelf edge. The cross-shelf structure of SLA indicates that sea-level signals depend on the shelf depth.

367

368 5 Discussion

369 5.1 Propagation of CSWs

Fig. 7 compares the data derived from the XWT of SLA with the dispersion

relation of CSW. We overlay the dispersion relation curves of CSWs with the results (phase speed, *c* and period, T, $\lambda = c \cdot T$) from tide-gauge data analysis in Section 4.1. One can see that the data points derived from station pairs are distributed near the dispersion relation of CSWs, implying the signals with the periods shorter than 40 d are the CSWs propagating along the shelf.

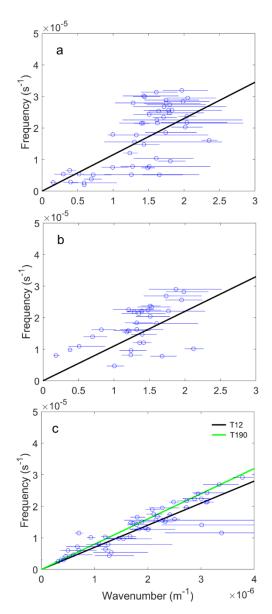


Fig. 7. Dispersion relation of the lowest mode of CSWs between (a) Kanmen and
Xiamen, (b), Xiamen and Shanwei, as well as (c) Shanwei and Zhapo. The data points
are calculated from the XWT of SLA. The curves are the theoretical dispersion relation
for the mean depth profiles listed in Table 1. Black and green curves in (c) represent the
dispersion relation (from Fig. 2) for the topographic profiles along tracks 12 and 190

Figs. 7a-b deal with the signals between Kanmen and Shanwei. We use the bathymetric profile near Track 240 and 88 of the altimeter satellite to calculate the dispersion relation curve, respectively. However, since the topography between Kanmen and Shanwei changes dramatically, the data points lie dispersedly around the curve. The wavenumbers of CSWs range from 0.1 to 2.4×10^{-5} m⁻¹ between stations Kanmen and Shanwei. The CSWs propagate along the coast with a phase speed of 10 m s⁻¹.

Fig. 7c shows the signals between Shanwei and Zhapo. The data points and the theoretical curves agree quite well. Compared with the topography between Kanmen and Shanwei, the topography changes slightly between stations Shanwei and Zhapo. Thus, we conclude that the signals in tide-gauge data with the periods shorter than 40 d are the CSWs propagating along the shelf. The phase speed of CSWs is about 8 m s⁻¹ in the study area, which is close to that of previous studies (Li et al., 2015; Li et al., 2016; Shen et al., 2021).

Previous studies present that the periods of CSWs range from two days to two 398 399 weeks in the SCS (Chen and Su, 1987; Li, 1993). The period of CSWs upstream, i.e., in the ECS, is often detected as several days (Ding et al., 2012; Ding et al., 2018; Hsueh 400 and Pang, 1989; Huang et al., 2015; Yin et al., 2014). Hsueh and Romea (1983) found 401 the sea level fluctuations with a period more than 13 d along the West Korean coast. 402 Worldwide, these low-frequency CSWs are common along the coast of Chile and the 403 east coast of the Indian Ocean (Castro and Lee, 1995; Hormazabal et al., 2001; Marshall 404 and Hendon, 2013; Vialard et al., 2009), where the width of the continental shelf is 405 narrow. In this study, we define the CSW with a maximum period of about 30-40 d. The 406 407 main reason for the difference is data length we used. The long-time series of DSLA helps analyze the abnormal low-frequency CSWs in the SCS. 408

In addition, we should take care that CSWs in Fig. 7 are mixed with wind-forced CSWs. It is difficult to separate the effect of wind-force and free propagating CSWs clearly. Hsueh and Romea (1983) found that there is clear coupling between surface winds and coastal sea level in the northeast China Sea. Li et al. (2016) found that the propagation time of the wind signals are much shorter than that of CSWs in the north SCS. Even one can see isolated cases which are away from dispersion relationship of CSW in Fig. 7, most of the data points are near the theoretical dispersion relation. Moreover, the stratification in the continental shelf is important for the characteristics of CSW. The sea level variation in this study should present baroclinic and barotropic CSWs in the shelf area together.

419

420 5.2 Trapped cross-shelf structure

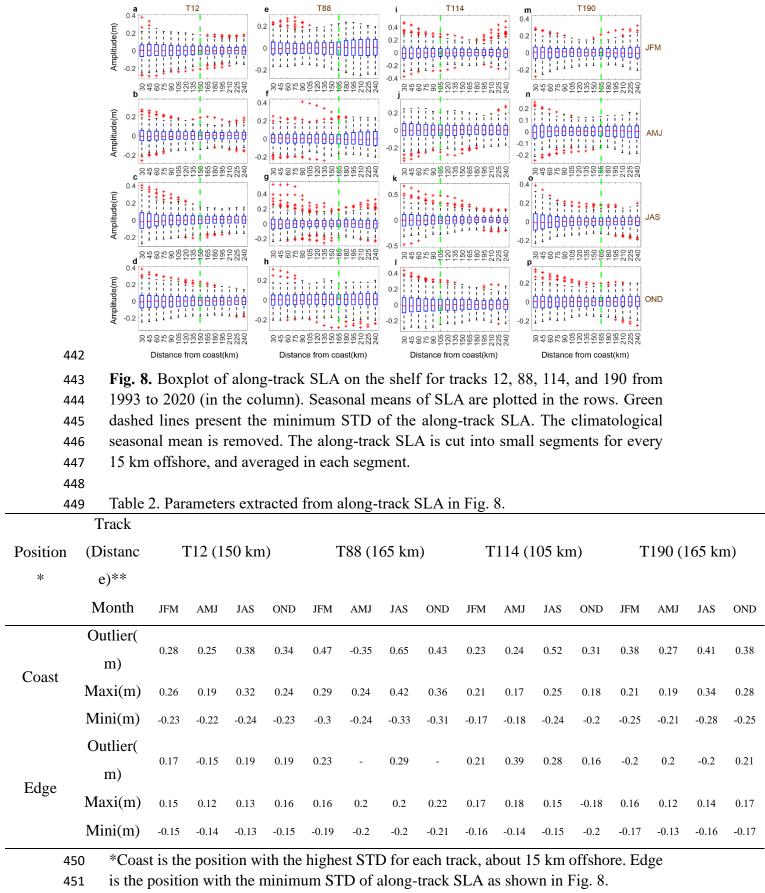
421 5.2.1 CSWs

As the sampling period of the altimeter satellites is about 10 d, the signals with the periods shorter than 20 d could not be distinguished from the along-track SLA based on the Nyquist sampling theorem. Fortunately, the cross-shelf structure of CSWs could be sampled as fragments by repeated satellite observations.

In this case, the CSWs with a period less than 40 d were abundant from 1993 to 2020 on the northern continental shelf of the SCS, especially in winter (Fig. 5). Even if the significance level is larger than 5%, the power is universal and continuous in this period band. Therefore, it could be considered that a large number of repeated observations by altimeter satellites were executed during CSW events.

Fig. 8 shows a boxplot of along-track SLA over the continental shelf for tracks 12, 431 88, 114, and 190. The maxima, minima, and outliers derived from Fig. 8 are listed in 432 Table 2. The whisker dashed black line outside of the box extends to the most extreme 433 data points, not considering the outliers at the 5 % significance level. One can see that 434 the trapped characteristics as the maximum amplitude (75th percentile), interquartile 435 range (IQR), and outliers (red plus sign) occur at the coastal side. In contrast, the 436 minima occur at the edge of the continental shelf (150, 165, 105 and 165 km offshore 437 for tracks 12, 88, 114 and 190). The amplitude of along-track SLA for maxima and IQR 438 decreases gradually from the coastline (0.2-0.4 m) to the edge of the continental shelf 439 $(\sim 0.1 \text{ m})$. The largest outlier is 0.65 m over the shelf 30 km offshore. 440

441



****** Distance between coastline and edge.

The largest SLA over the shelf occur from July to September. For example, the 75th percentile of SLAs near the coastline is about 0.25 m for track 144, and 0.42 m for track 88. The largest outlier of 0.65 m also occurs from July to September for track 88. The smallest SLA occurs from April to June when the wind is weak during the monsoon transition period (Wang et al., 2009). The extreme data from July to September show the occurrences of storm surges over the shelf (Chen et al., 2014).

Overall, the maximum amplitude, IQR, and extreme data of the along-track data 460 over the shelf show the trapped wave characteristics, which shows that SLA decreases 461 gradually from the coastline to the edge of the continental shelf. The trapped 462 characteristics from the along-track SLA are similar to the cross-shelf structure of 463 normalized amplitudes of the mode 1 CSWs, as shown in Fig. 3a. It should contain 464 higher modes in the along-track SLA. The along-track SLA along tracks 153 and 229 465 show similar characteristics (not shown). That along track 240 (as shown in Fig. 6n) 466 presents a differentiated pattern in the coast side and shelf edge during May-July. The 467 468 main reason should be the existence of cold eddy in the north of Taiwan Island.

To further reveal the variations in the along-track SLA in track 12 on the shelf, the 469 empirical orthogonal function (EOF) analysis results are shown in Fig. 9 (EOF analysis 470 for data along other tracks is similar, not shown here). The first four EOF modes of the 471 along-track SLA explain 96.8% of the total variance. Mode 1 explains the seasonal 472 variance of SSH (red curve in Figs. 6o-u). The seasonal cycle could be characterized 473 clearly from Fig. 9b, with peaks in October and troughs in May. Mode 2 and Mode 3 474 are similar to the cross-shelf structure as shown in Fig. 3a, which explains 25.1% of the 475 total variance. Mode 3 is influenced by the background of SSH in the open sea side. 476 477 The amplitude is 0.3-0.4 m, which is comparable to the outlier data as shown in Table 2. Mode 4 is similar to the cross-shelf structure of mode-2 CSW as shown in Fig. 3b, 478 which only explains 2.8% of the total variance. The EOF analysis indicates that CSW 479 could explain <30% of the total variance of SSH in the study area, which is comparable 480 to the IQR in boxplot of along-track SLA (Fig. 8). 481

453

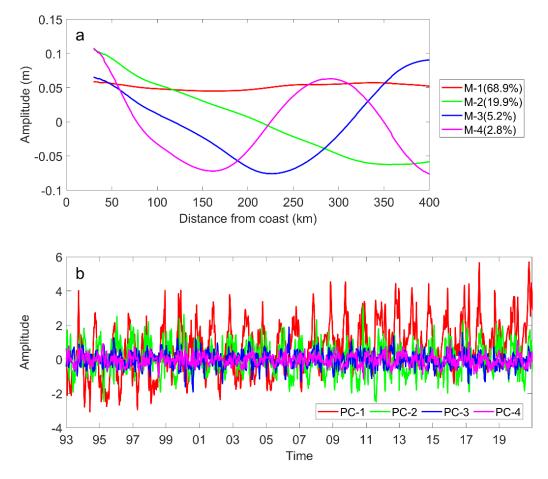
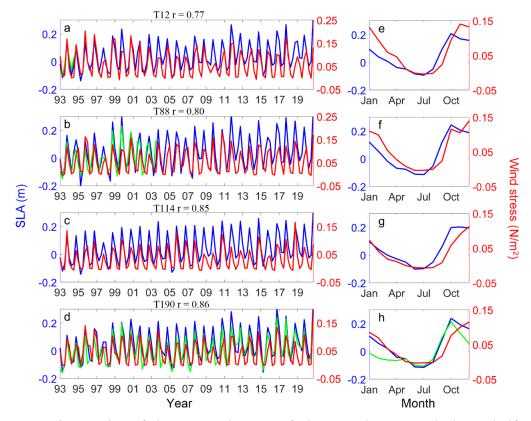


Fig. 9. The first four EOFs for along-track SLA on the shelf along track 12. Cross-shelf
amplitude (a) and time series (b) of EOFs. The variance explained by each mode is
labeled in (a).

483

488 5.2.2 ATWs

As shown in Fig. 6, the inter-seasonal and seasonal signals could be distinguished 489 from along-track SLA. Fig. 10 shows the seasonal mean of along-track SLA and along-490 shelf sea surface wind. One can see that the time series of the seasonal mean of along-491 492 track SLA show the seasonal variation with the amplitude of 0.1-0.2 m at 15 km offshore. The time series of the seasonal mean of DSLA show similar characteristics 493 with along-track SLA. The correlation relationship between the seasonal mean of 494 along-track SLA and along-shelf sea surface winds reaches >0.77. The monthly mean 495 of SLA during April-September is negative, while positive in the other months. It is 496 attributed to that the local wind stress substantially influences the coastal sea level (Lin 497



500

Fig. 10. Time series of the seasonal mean of along-track SLA and along-shelf sea surface wind stress for tracks (a)12, (b) 88, (c) 114, (d)190, and (e-h) monthly climatological mean of along-track SLA and along-shelf sea surface wind stress. Red curves represent the along-track SLA at 15 km offshore. Green curves represent the seasonal mean of sea level data at tide-gauge stations (a) Xiamen, (b) Shanwei, and (d) Zhapo.

507

In addition, one can also see a characteristic of out of sync between seasonal mean 508 of along-track SLA and wind. The maximum of mean sea surface wind stress occurs in 509 November and December. While, that of monthly along-track SLA occurs in October. 510 Ding et al. (2020) investigated the seasonality of coastal circulation in the north SCS 511 using a numerical model. The result indicates that the maximum of water transport is 512 1.93 Sv occurred in autumn. Li et al. (2023a) found the along-shelf current is strongest 513 in October at approximately 0.17 m s⁻¹. The along-shelf current involves with ATW, 514 which is the reason why there is a difference between monthly climatological mean of 515 along-track SLA and sea surface wind stress as shown in Figs. 10e-h. 516

Fig. 11 shows the cross structure of the seasonal mean of along-track SLA. One can see that the SLA on the coastline side is lower than on the ocean side from April to September. In the other seasons, the slope of along-track SLA is the opposite. The along-track SLA presents similar characteristics as shown in Fig. 10, i.e., the variation of SLA is controlled by the sea surface winds over the shelf.

522

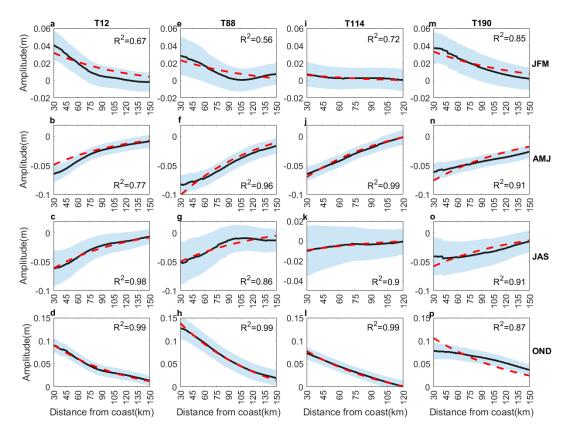




Fig. 11. Fitting climatologic seasonal mean of along-track SLA (black curves) using
cross-shelf ATWs (red dashed curves) for tracks (a-d) 12, (e-h) 88, (i-l) 114, and (m-p)
190. Dashed shadow is the STD of the seasonal mean of along-track SLA.

527

Meanwhile, from Fig.11 one can see the cross-shelf structure of along-track SLA. The fitting curves show that theoretical ATWs (Fig. 4) explains the cross-shore structure of along-track SLA very well. The amplitudes of ATWs in track 12 are 0.04 m, -0.06 m, -0.05 m, and 0.10 m in January-March, April-June, July-September, and October-December. That along track 88 is relatively larger, e.g., 0.13 m in October-December. While, the minimum amplitude occurs in track 114. Differently from Fig. 8, the amplitudes of ATWs during October-December are larger than that during JulySeptember. It should be attributed to that the monsoon winds in winter are stronger thanthat in summer.

Lin et al. (2021) investigated the tilt of mean dynamic topography along the coast of the Chinese mainland. Wu (2021) used a nondimensional parameter (Pe_{β}) to describe the influence of open ocean forcing on shelf circulation, which is determined by the ratio of long-wave-limit planetary to TRW speeds. In this study, $Pe_{\beta} < 1$, which indicates shelf currents decayed rapidly toward the coast. Their results confirm that the ATWs predict the coastal dynamic topography over the continental shelf of the SCS well.

543

544 6 Summary

Using sea level data derived from the tide-gauge stations Kanmen, Xiamen, Shanwei, Hongkong, and Zhapo, this study analyzes statistical features of the CSWs, inter-seasonal and seasonal signals. Meanwhile, along-track SLA data derived from multiple satellite altimeters from 1993 to 2020 are applied to detect the cross-shelf structures of the signals. The major results are summarized as follows.

1) CSWs of periods shorter than 40 d propagate along the coast with the phase speed
 of about 10 m s⁻¹ in the ECS and 8 m s⁻¹ in the SCS. The dispersion relation indicates
 that the waves belong mode 1 CSWs.

2) Owing to the fact that the repeated observation period of the satellite altimeters is 553 comparable with that of CSWs, we combine fragments of the numerous repeated 554 observations of along-track SLA to reconstruct the cross-shelf structure of CSWs. 555 The results show that the maximum amplitudes of CSWs have remarkable seasonal 556 variability, about 0.6 m during July-September, while only 0.2 m during April-June. 557 The reconstructed cross-shelf structures of CSWs confirm the property of mode 1 558 CSWs. Moreover, the energy is trapped within the partial continental shelf 559 shallower than 200 m. 560

3) The inter-seasonal and seasonal signals present as ATWs on the continental shelf.
The amplitudes of ATWs have remarkable seasonal variability, ~0.10 m during
October-December, twice larger than 0.04 m, 0.05 m and -0.06 m during JanuaryMarch, July-September, and April-June, respectively. These results reveal that the

local wind stress substantially influences ATWs on the continental shelf.

4) The results derived from the observation data of along-coast tide-gauge stations
combined with cross-shelf tracks of satellite altimeters are interpreted well by the
framework of linear wave theory. It implies that the technological approaches
developed in this study are suitable for constructing the cross-shelf structures of
CSWs and ATWs on the continental shelf.

However, owing to the neglective wind stress and baroclinicity, higher modes of waves are not discussed in this paper. Observations from moorings and numerical models will be used in our future studies to obtain the characteristics of baroclinic coastal trapped waves.

575

576 Acknowledgments

This research was funded by the National Key Research and Development Program of
China (2022YFC3104805); National Natural Science Foundation of China (42276019,
41976200, 41706025); Innovation Team Plan for Universities in Guangdong Province

(2019KCXTF021); First-class Discipline Plan of Guangdong Province (080503032101,
231420003); Observation and Research Station for Tropical Ocean Environment in

582 coastal water west of Guangdong.

583

584 Data Availability Statement

The tide gauge data are available at https://psmsl.org/data/obtaining/. The along-track
SLAs is obtained at ftp-access.aviso.altimetry.fr.

587

588 Author contributions

589 JYL were responsible for writing the original draft. Review and editing were conducted

by QAZ. Conceptualization was handled by JYL, QAZ and LLX. TH and YX were

responsible for data curation. LLX acquired funding.

592

593 **Competing interests**

594 The contact author has declared that none of the authors has any competing interests.

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