



Open-path measurement of stable water isotopologues using mid-infrared dual-comb spectroscopy

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Abstract.

We present an open-path mid-infrared dual-comb spectrometer (DCS) capable of precise measurement of the stable water
10 isotopologues H₂¹⁶O and HD¹⁶O. This system ran in a remote configuration at a rural test site for 3.75 months with 60 %
uptime and achieved a precision of <2 ‰ on the normalized ratio of H₂¹⁶O and HD¹⁶O (δD) in 1000 seconds. Here, we compare
the δD values from the DCS to those from the National Ecological Observatory Network (NEON) isotopologue point sensor
network. Over the multi-month campaign, the mean difference between the DCS δD values and the NEON δD values from a
15 similar ecosystem is <2 ‰ with a standard deviation of 18 ‰, which demonstrates the inherent accuracy of DCS measurements
over a variety of atmospheric conditions. We observe time-varying diurnal profiles and seasonal trends that are mostly
correlated between the sites on daily time scales. This observation motivates the development of denser ecological monitoring
networks aimed at understanding regional and synoptic scale water transport. Precise and accurate open-path measurements
using DCS provide new capabilities for such networks.

1 Introduction

20 A better understanding of water transport on different scales is necessary to understand the impacts of climate change on global
water use (Jury and Vaux, 2005). A continuous record of stable water isotopologues in atmospheric water vapor and
precipitation can provide a benchmark for models of water transport ranging from global circulation to field-scale mass balance
(Galewsky et al., 2016; Al-Oqaili et al., 2020; Welp et al., 2008; Good et al., 2015; Dansgaard, 1964; Craig et al., 1965). For
many of types of water isotopologue measurements, it is necessary to have a network of accurate sensors. Here, we demonstrate
25 accurate, long-term open-path dual-comb spectroscopic measurements of H₂¹⁶O (H₂O) and the deuterium-substituted
isotopologue HD¹⁶O (HDO) in atmospheric water vapor. The ratio of HDO to H₂O, referred to as δD, is characterized using
the standard definition (Galewsky et al., 2016):



$$\delta D = \frac{R_{measured}}{R_{VSMOW}} - 1, \quad (1)$$

30 where $R_{measured} = [\text{HDO}]/[\text{H}_2\text{O}]$ is the measured isotopologue ratio (ratio of absolute concentrations) and $R_{VSMOW} = 0.0003115$
is the standard isotope ratio according to the Vienna Standard Mean Ocean Water (VSMOW) scale. Normally, δD is expressed
as a “per mil” number (‰). Values of δD in nature fall in the range from -500 ‰ to 0 ‰ with spatial and temporal variation
driven by mass fractionation during condensation and evaporation processes. By monitoring δD , it is possible to
simultaneously track local evapotranspiration in the planetary boundary layer and synoptic scale water transport (Noone et al.,
35 2013; Galewsky et al., 2016).

Currently, most measurements of stable water isotopologues in atmospheric vapor rely on point sensor networks
(Finkenbiner et al., 2022; Fiorella et al., 2021; Wei et al., 2019; Aemisegger et al., 2012) that can be difficult to calibrate
(Bailey et al., 2015; Rambo et al., 2011) and utilize relatively expensive technology. Cryogenic discrete point sampling and
cavity-enhanced optical absorption point sensors achieve accuracy on the order of 2 ‰ or less (Galewsky et al., 2016) but
40 require careful, frequent calibration. Long open-path measurements could enhance the capabilities of sensor networks if the
measurements have high precision and accuracy. In general, open-path sensing techniques for atmospheric gases avoid
sampling biases inherent in extraction systems (inlets, tubing, pumps, filters, etc.), which is especially advantageous for
“sticky” gases such as water vapor or ammonia. In addition, the spatial averaging provided by open-path techniques enables
clearer comparisons between field measurements and atmospheric models (typically computed on kilometer-scale grids)
45 (Griffith et al., 2018). Open-path Fourier transform infrared (FTIR) spectroscopy for δD retrieval has been reported previously
(Wang et al., 2012; Griffith et al., 2018), but sources of error can arise in open-path FTIR related to the instrument lineshape
and limited spectral resolution of deployable systems (Lin et al., 2020), and have not been investigated for δD .

Open-path dual-comb spectroscopy (DCS) (Rieker et al., 2014; Waxman et al., 2017; Giorgetta et al., 2021; Cossel
et al., 2021) is promising for stable water isotopologue measurements due to the combination of broad spectral bandwidth,
50 high spectral resolution, high brightness, and spatial coherence, which enables kilometer-scale measurements of multiple trace
species simultaneously (Coddington et al., 2016). Open-path DCS is an emerging Fourier spectroscopy technique based on
sampling the atmosphere with laser frequency combs. This technique achieves high spectral resolution (here, the spectral point
spacing is 200 MHz or 0.0067 cm^{-1}) with no instrument lineshape while also spanning hundreds of wavenumbers (Waxman et
al., 2017). Several demonstrations have already shown the potential of laboratory DCS for isotopologue measurements
55 (Muraviev et al., 2018; Vodopyanov, 2020; Parriaux et al., 2022). Open-path DCS optical paths can be easily reconfigured
using remotely controlled telescopes to adjust to measurement conditions (Coburn et al., 2018) and can be spatially scanned
using mobile reflectors (Cossel et al., 2017), which could enable better tomographic characterization of evaporation and
transpiration in agricultural systems and natural ecosystems (Welp et al., 2008).

In this work, we demonstrate a mid-infrared (MIR) open-path DCS system capable of δD measurement with a
60 precision of 1.2 ‰ with 3000 seconds of averaging. This precision is sufficient to capture diurnal changes in δD at our test



measurement site (Platteville Atmospheric Observatory; PAO) which can range from 10 ‰ to 100 ‰. We demonstrate measurements over a 3.75-month long measurement period with 60 % uptime from fall into winter through a variety of meteorological conditions. We compare these DCS δD measurements to data from two nearby National Ecological Observatory Network (NEON) sites: Central Plains Experimental Range (CPER) and Niwot Ridge (NIWO) (Finkenbiner et al., 2022). The difference in δD values measured by DCS at PAO and by the point sensor at CPER have a mean of less than 2 ‰ and a standard deviation of 18 ‰. The NIWO site is in an alpine ecosystem whereas PAO is located in a great plains ecosystem, so one would not expect similar level of agreement. A comparison between δD at PAO and NIWO demonstrates that gradients of δD with altitude near mountainous regions are strong but not omnipresent. Finally, we look at some features of the data series including diurnal profiles and Keeling curves for a few time periods to demonstrate subtle differences between PAO and CPER meteorology. These results demonstrate that MIR open-path DCS can be a viable tool for water isotopologue measurements and can provide data that complement and extend current monitoring capabilities.

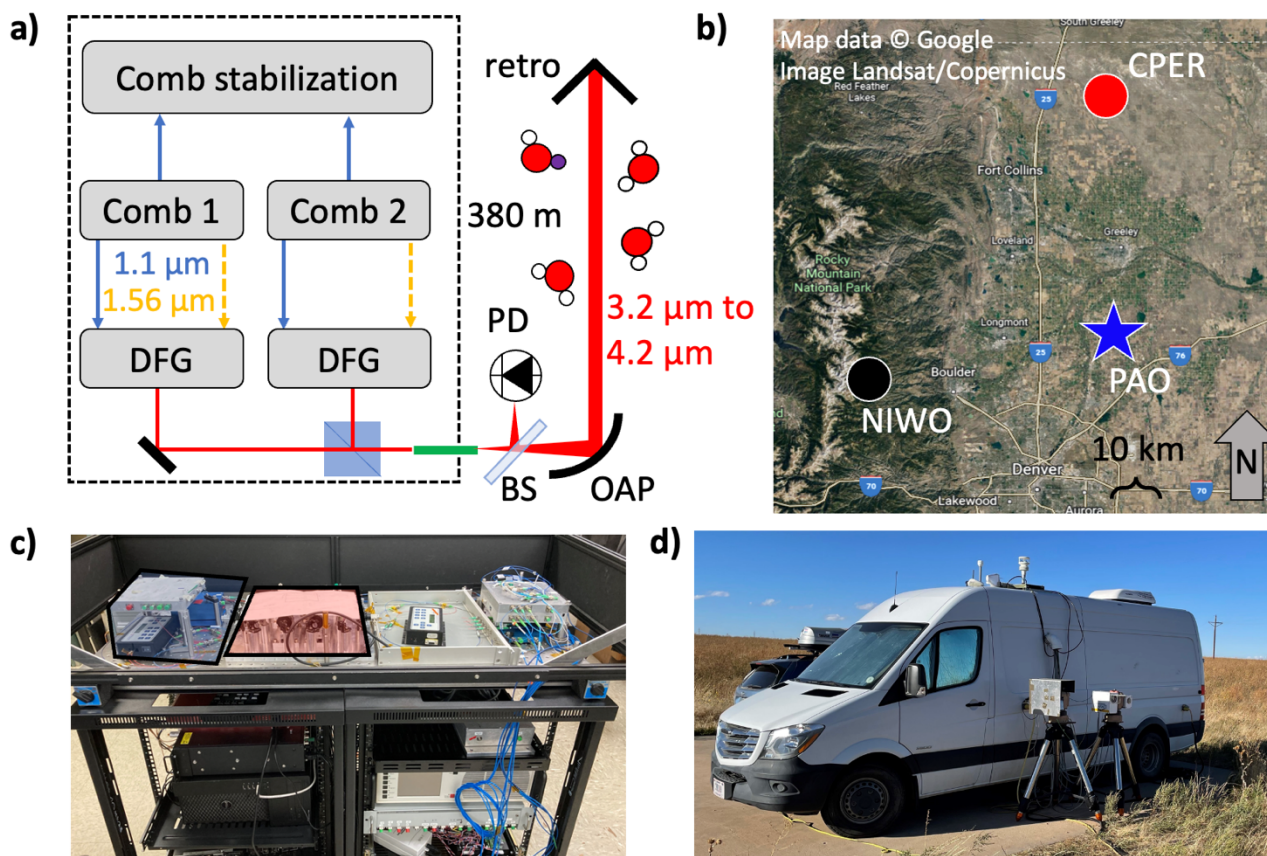
2 Experimental Setup

2.1 MIR DCS generation

The deployable MIR DCS system illustrated in Figure 1a follows the design from (Ycas et al., 2020) to generate laser light spanning 3.2 μm to 4.2 μm while leveraging a robust fiber laser architecture. We briefly review the design below along with minor changes that enabled improved field operation of the MIR DCS. The MIR light is produced using difference frequency generation (DFG) in aperiodically-poled lithium niobate (aPPLN) waveguides seeded by near-infrared fiber frequency combs (Erny et al., 2007; Maser et al., 2017; Ycas et al., 2018). The system is based on two fully-stabilized near-infrared combs (Sinclair et al., 2015) centered at 1560 nm with repetition frequencies of ~ 200 MHz and a repetition frequency difference of 208 Hz. By supplementing standard DCS locking schemes (Truong et al., 2016) with digital phase correction (Roy et al., 2012; Ycas et al., 2018), the MIR DCS achieves high mutual coherence in field environments.

To generate the MIR light through DFG, the output of each comb is first split into two branches, as discussed in (Ycas et al., 2020). One branch uses spectral broadening in highly nonlinear fiber (HNLF) to generate pulsed light spanning 1.05 μm to 1.2 μm , which, unlike the system reported in (Ycas et al., 2020), is directly coupled into free-space from the HNLF using an off-axis parabolic (OAP) mirror for improved efficiency. In the second branch, pulses are amplified at 1565 nm (~ 20 -nm bandwidth) and temporally stretched in 5 m of PM fiber to a pulse width of about 1 ps. The ~ 1.1 μm and 1.565 μm light from both branches is then combined on a long-pass dichroic mirror with a cut-off of 1180 nm and coupled into the aPPLN waveguide by a $f=25$ mm anti-reflection (AR) coated achromatic lens. The aPPLN waveguide is chirped to support broadband frequency conversion of the near infrared light to the 3.0 μm to 4.0 μm region (Suchowski et al., 2009). Slow thermal or other effects that cause pulse walk-off between the two branches will reduce the MIR light generation in the aPPLN waveguide. Placing the HNLF in the first branch in an insulated metal housing and stretching the second branch pulse duration to 1 ps

both greatly reduced pulse walk-off. Any residual pulse walk-off is compensated by a fiber-coupled computer-controlled optical delay line placed in the 1.565 μm branch which is adjusted to optimize the MIR DCS output power.



95 **Figure 1: MIR DCS experimental setup.** a) Schematic of MIR DCS with dotted line indicating the boundary of the mobile DCS van. DFG: difference frequency generation. PD: photodetector. BS: beamsplitter. OAP: off-axis parabolic mirror. Red lines indicate free-space MIR beam paths. Blue lines indicate 1560 nm light. Dashed orange lines indicate supercontinuum light near 1150 nm. Green line indicates InF₃ single-mode fiber. b) Satellite image of Northern Colorado Front Range Urban Corridor with PAO (blue star), CPER (red circle) and NIWO (black circle) sites marked. The NIWO and CPER sites are both ~75 km away from the PAO site. Map Data: Google. c) Image of the field-deployable MIR DCS system. Optical platform containing combs (Comb 1 shaded in blue box) and DFG optics (shaded in red box) is placed on top of two electronics racks containing comb control and acquisition systems. d) Image of mobile DCS van at PAO site with MIR telescope system in the foreground. (The second telescope supports parallel measurements from a near-infrared DCS system that are not discussed here.)

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105 The MIR output of the aPPLN waveguide from one comb is collimated, filtered using an AR-coated germanium wedge, combined with the MIR output of the other comb on a broadband AR-coated CaF₂ beamsplitter and coupled into InF₃ fiber for transport to the telescope (see Figure 1d). The spectrum coupled into fiber covers a continuous region from 3.2 μm to 4.2 μm , up to the CO₂ absorption band. More than 1 mW of fiber-coupled power is available in the 3.5 μm to 4.0 μm region where there is a strong HD¹⁶O absorption band centered at 3.6 μm . A waveguide redesign or modification of the DFG approach



110 will enable coverage from 3.0 μm to 5.0 μm simultaneously (Zhou et al., 2020). Such bandwidth would give access to cleaner H_2^{18}O and H_2^{16}O absorption bands in addition to the strong HDO and H_2^{16}O bands covered with the current spectral bandwidth.

2.2 Field Deployment of the MIR DCS

The entire MIR DCS optical system (free-space and fiber optics) is placed on a 1.2 m by 0.5 m platform which is then
115 stacked on two electrical rack units (each measuring 0.5 m by 0.8 m by 0.84 m) containing the pump diodes, temperature controllers, detectors, locking electronics, DCS acquisition system, computers and power supplies (Figure 1c). Many of these components could be further reduced in size to make the system more compact. To limit sensitivity to environmental perturbations, the actual free-space optics for dual-comb DFG only occupy a 0.35 m by 0.35 m custom breadboard. The platform and breadboard are covered by a laser eye-safe box and the entire double rack unit (with a volume of 0.75 m^3) is
120 installed in the back of a van, which is then driven to the PAO site near Platteville, Colorado [40.181782 N, 104.725054 W] as shown in Figure 1b. This rural site is in the Denver-Julesburg basin, which is a major oil and gas production area.

The fiber-coupled MIR DCS light is sent to a telescope located outside of the van (Figure 1d). In the telescope (Figure 1a), the diverging beam from the InF_3 fiber is collimated using an $f = 178.55$ mm OAP, which results in a 6.8-cm diameter beam that can be projected over long paths. The telescope is placed on a tripod-mounted azimuth-elevation gimbal and is
125 aimed at a 12.7-cm diameter retroreflector located 380 m away. The DCS light reflected from the retroreflector returns to the telescope and is then reflected by a broadband AR-coated CaF_2 beamsplitter placed between the fiber end and the collimating OAP onto a TEC-cooled HgCdTe detector. A co-axial cable transmits the detected dual-comb interferogram voltage signal back to the van where it is digitized, phase-corrected and coherently co-added to generate one DCS spectrum every 2 minutes (Figure 2a). Each of these averaged spectra is derived from $\sim 25,000$ raw unaveraged interferograms. In addition to the DCS
130 spectrum, several other datasets are available for analysis. Meteorological data are provided by a sensor at the PAO site operated by the National Oceanic and Atmospheric Administration (NOAA, site code: PVL). This data includes wind direction and speed, solar radiation, relative humidity, temperature and precipitation. Water vapor levels are also recorded at the site using two different commercial cavity-ringdown spectroscopy instruments. We note that the MIR DCS described here is also capable of precise detection of multiple species in addition to water isotopologues including primary greenhouse gases (CO_2 ,
135 CH_4 , N_2O), air pollutants (*e.g.* CO) and volatile organic compounds (*e.g.* HCHO and C_2H_6).

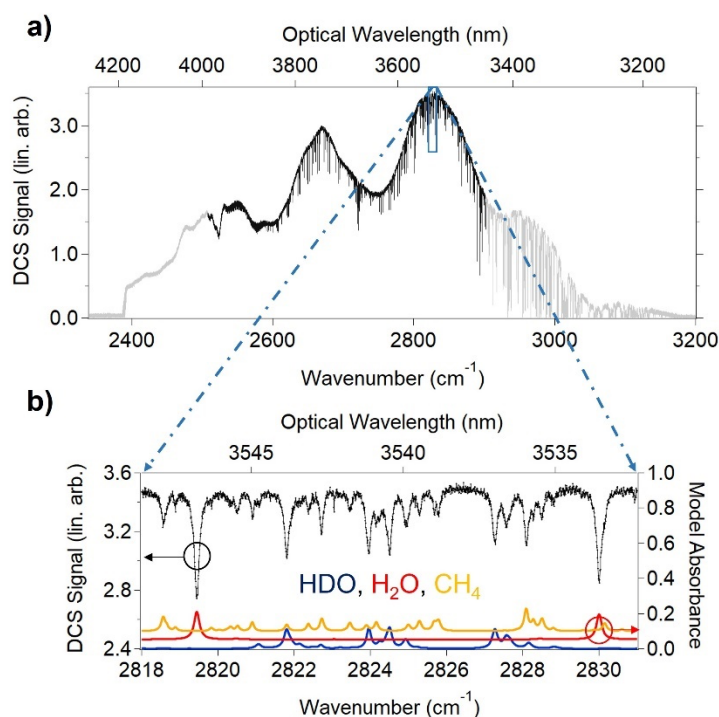
3 Data Analysis

3.1 Broadband Cepstral Domain Fitting of MIR DCS data

The deployable open-path MIR DCS generates broadband, high-resolution atmospheric spectra shown in Figure 2a that cover rovibrational absorption features from many different atmospheric species. In order to isolate water isotopologue absorption
140 from these congested broadband spectra, one could select several “micro-windows” where traditional frequency-domain fitting



techniques can successfully extract concentrations in a reliable manner (Griffith et al., 2003; García et al., 2022). In this work, we instead choose to utilize a broadband fit using cepstral-domain analysis of molecular free-induction decay (Cole et al., 2019). The cepstral transform starts with an intensity spectrum in the frequency domain which is then converted to an absorbance spectrum by taking the logarithm. This absorbance spectrum is then Fourier-transformed to yield a “cepstrum” which represents the modified molecular free-induction decay (FID) in the time (“cepstral”) domain. The cepstral domain technique simplifies the fitting of broadband DCS data by separating the baseline structure, etalon effects, and FID into different cepstral regions. The data are fit in the cepstral domain using a 10 ps to 350 ps bandpass filter. The 10 ps cut-on is applied to the DCS spectra to remove baseline effects and the 350 ps cut-off avoids etalons due to the PPLN waveguides themselves. No other significant etalons occur between these two end points. The effect of the cepstral filter choice on the fit retrievals is discussed in Section 3.3. To avoid complications of fitting deep absorption features that approach the spectrum noise floor, the fit region is restricted to the weaker absorption section of the DCS spectrum ranging from 3.5 μm to 4.0 μm (see Figure 2a). This region includes most of the ν_1 absorption band of HDO as well as the long wavelength end of the $2\nu_2$ band of H_2O . The large number and strength of HDO transitions within this fit region ensures a very precise measurement of this secondary isotopologue. We note that the peak absorption strengths of HDO and H_2O in this region are similar, a desirable trait for any optical isotopologue detection technique. Importantly, the 200 MHz point spacing of the DCS spectra fully resolves the absorption lines of both HDO and H_2O as well as cluttering species like CH_4 (see Figure 2b).



160 **Figure 2:** a) MIR DCS spectrum (2-minute average) taken over the 760 m round trip path at PAO (grey trace). Black trace indicates region selected for spectral analysis in this work. b) Zoomed-in MIR DCS spectrum from 2818 cm^{-1} to 2831 cm^{-1} . Black dots indicate DCS signal (left axis). The red and blue traces are the H_2O and HDO absorption models, respectively, from HITRAN2020 (right



axis with H₂O offset by 0.05). The orange trace is the CH₄ absorption model from HITRAN2020 (right axis and offset by 0.1). Although absorption from other molecules contribute additional clutter in this region, CH₄ is the major source.

We model the FID signals using absorption cross-sections generated using line-by-line data from HITRAN2020 (Gordon et al., 2021). We include H₂¹⁶O, HD¹⁶O and H₂¹⁸O absorption, but the retrieved H₂¹⁸O concentration does not have sufficient precision to be useful due to the weak absorption in the fit spectral region. In addition to water, the fit model includes CH₄, C₂H₆, N₂O, CO₂, HCHO and HCl. Concentrations of these species can be inferred from the DCS spectra with useful precisions for analysis of ambient air. In the cepstral fit, temperature is allowed to vary while the pressure is fixed to a measured value at PAO, allowing for calculation of the atmospheric number density for air according to the ideal gas law (although the isotopologue ratio $R_{measured}$ and δD are independent of the total air mass).

Data were acquired from October 5th, 2021, to January 27th, 2022; time series data is plotted in Figure 3. The quality of the data is characterized by a DCS figure-of-merit (FOM), equal to the signal-to-noise ratio at 1-second averaging multiplied by number of spectral elements (Coddington et al., 2016). Over the three-month data acquisition, the highest FOM values achieved were in the 1.0×10^6 to 1.5×10^6 range, comparable to recent high-quality laboratory DCS results (Ycas et al., 2018). We only analyze data above a defined FOM threshold of 2.0×10^5 , which effectively removes only extremely noisy data points with poor return power. For this threshold, our uptime over the multi-month measurement period was 60 %. The main cause of dropouts was telescope misalignment, not loss of DCS stabilization or degradation of MIR light generation. The daily average FOM decreased by about 50 % over the course of the 3-month long measurement campaign, mostly due to contamination build-up on the retroreflector.

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3.2 Temperature Correction

As discussed above, the DCS fit is used to extract a path-averaged temperature. We chose this method over using a separate temperature measurement since a point temperature measurement does not accurately represent the temperature along the path due to potential horizontal and vertical temperature gradients (Burns et al., 2012). In particular, the NOAA temperature measurement at PAO is at a different height above ground than the optical path, leading to differences due to vertical gradients that oscillate over the course of a day. However, because the path-averaged temperature is retrieved from the cepstral fit, biases in the spectral database parameters will bias the temperature retrieval. To correct for this, we assume that the DCS and NOAA measurement should agree on average for daily timescales even though there are oscillations on shorter time scales. Figure 4 shows the correlation between the DCS temperature (T_{DCS}) and the PAO NOAA sensor (T_{PAO}). The two sensors demonstrate a strong linear correlation albeit with a non-unity slope and a small offset: $T_{DCS} = c_0 T_{PAO} + c_1$, where $c_0 = 0.88^\circ\text{C}/^\circ\text{C}$ and $c_1 = 5.33^\circ\text{C}$. This bias is not surprising as the HDO linestrengths in this region are only known to ~ 3 % accuracy and the broadening coefficients may have up to 10 % errors (Devi et al., 2017). A bias in the retrieved temperature due to spectral database errors will also result in a potential bias in δD . By refitting the same DCS data over a range of fixed temperatures,

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we find the bias in the extracted δD value, $\Delta\delta D$, for a given temperature bias, ΔT , has a linear form with $\Delta\delta D = c_2 \Delta T + c_3$,
 195 where $c_2 = 10.4 \text{ ‰/°C}$ and $c_3 = -1.2 \text{ ‰}$. By substituting the linear relationship between DCS temperature and NOAA PAO
 temperature into this equation, we can globally correct the DCS data for the HITRAN temperature dependence errors. The
 final correction takes the form:

$$\delta D_{corrected} = \delta D_{uncorrected} - \left\{ c_2 \left[\left(1 - \frac{1}{c_0} \right) T_{DCS} + \frac{c_1}{c_0} \right] + c_3 \right\}. \quad (2)$$

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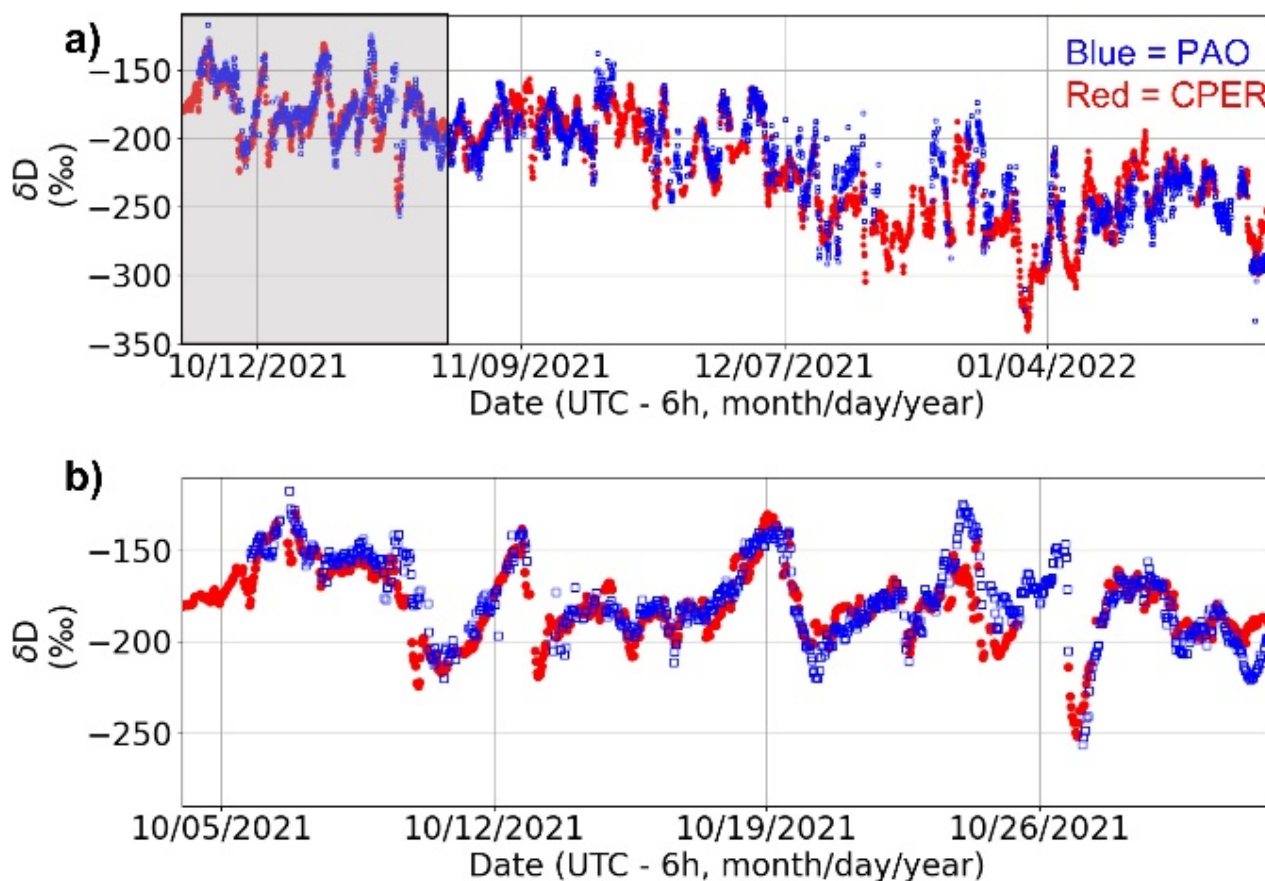
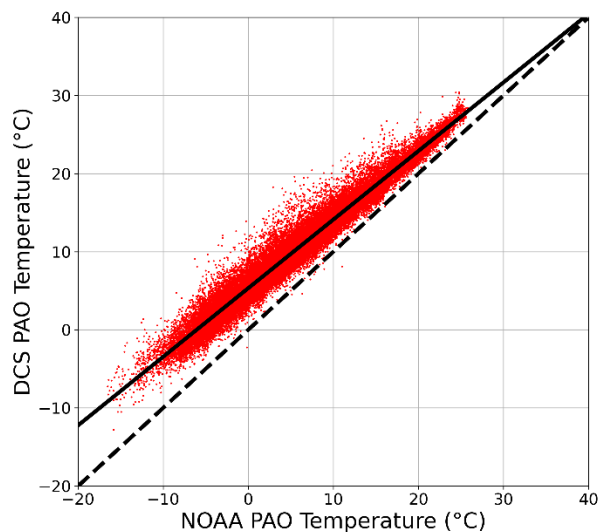


Figure 3: a) Time series of δD over the campaign. Blue squares: open-path DCS at PAO. Red dots: NEON point sensor at CPER.
 b): Zoom-in of time series of δD for October 2021 (section highlighted in grey in a)).



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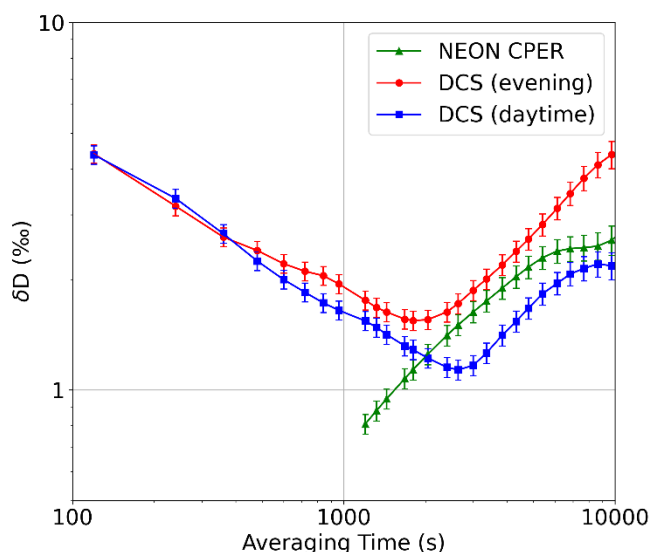
Figure 4: Accuracy of DCS temperature retrieval. Red dots indicate the DCS temperature at PAO versus NOAA temperature at PAO for 2-minute windows. A linear fit to this data (solid line) returns a slope of $0.88^{\circ}\text{C}/^{\circ}\text{C}$ and an intercept of 5.33°C ($R^2 = 0.96$). The dashed line indicates 1-to-1 correlation.

210 3.3 Time Series and Precision Analysis

The time series of δD values taken at PAO and CPER is shown in Figure 3. As will be discussed in depth in Section 4.1, the PAO values closely match the CPER values for the entirety of the measurement campaign. We notice significant fluctuation of the δD value over the course of the day as well as slow changes in the average value over the long measurement period. This behaviour is consistent with ground-based and satellite-derived δD columns (Schneider et al., 2020; Fiorella et al., 2018; 215 Noone et al., 2013) in the central United States where variations in synoptic weather patterns can bring water vapor from many different climates over the Northern Colorado area. To estimate the precision of our open-path DCS δD measurement, we perform an Allan analysis (Werle et al., 1993) of the data taken during the day on October 31st, 2021, from 5:15 am to 3:30 pm. On this day, the CH_4 levels were stable and near the minimum-observed level ($2.02 \pm 0.01 \mu\text{mol}/\text{mol}$) for >10 hours, indicating well-mixed atmospheric conditions. As shown in blue squares in Figure 5, the Allan deviation reaches a minimum 220 of 1.2 ‰ in less than 3000 seconds, after which the true atmospheric variations drive the deviation to higher values. The NEON point sensor at CPER sees a similar level of drift at the hour timescales (green triangles), well above the measurement error of commercial cavity-enhanced systems (Bailey et al., 2015). Furthermore, the Allan deviation of the open-path DCS measurements from a time period with more atmospheric variability (red circles in Figure 5) shows a minimum at an earlier time, which further indicates that the long-time deviations are driven by atmospheric variability. As we will demonstrate in 225 Section 5, single-per-mil precision in less than an hour makes it possible to use open-path DCS to track diurnal and seasonal



trends at PAO that match the observed trends from an established point sensor network. This analysis proves that open-path DCS is a viable alternative to extractive point sampling techniques for continuous ecological monitoring of isotopes.



230 **Figure 5: Allan analysis.** The overlapping Allan deviation (OADEV) for open-path DCS data from 10/31/2021 when the atmosphere was well-mixed (5:20 am on 10/31/2021 to 3:30 pm on 10/31/2021) is shown with blue squares. The OADEV is 4 ‰ at 120 seconds and averages down with the square root of time to 1.1 ‰ at 3000 seconds. After 3000 seconds, the OADEV rises steeply, likely due to real atmospheric fluctuations. For comparison, the OADEV for the NEON CPER point sensor is shown for the same period (green triangles). The common upward trend in both indicates variation due to atmospheric fluctuations. Finally, the OADEV for a DCS measurement under more typical mixing conditions (from 3:30 pm on 10/31/2021 to 1:30 am on 11/1/2021) is shown with red dots and exhibits the expected higher instability at longer times due to greater atmospheric fluctuation.

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3.4 Sources of Error in Open-Path Isotopologue Sensing with DCS

An evaluation of the accuracy of the DCS δD retrieval by comparison with the CPER NEON site is discussed in Section 4.2. In this section, we review the main potential sources of error in the open-path DCS measurement: detection and digitization nonlinearities (Guay et al., 2021; Malarich et al., 2022), concentration retrieval algorithm bias, and spectral model error. To limit the effect of detection and digitization nonlinearities, the received power was limited to $<30 \mu W$ and high-linearity amplifiers were used. We have previously confirmed that the detection is linear at these power levels. We also do not observe correlations between retrieved gas concentrations and received power, which indicates linear response.

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As with any absorption spectroscopy, the retrieval algorithm can introduce biases depending on the representation of the spectral intensity baseline and treatment of instrumental effects such as etalons. Here, the cepstral technique was used to separate baseline terms from the molecular absorption via filtering in the cepstral domain. The 10-ps cepstral filter was chosen to minimize mean-squared error of the residuals while maximizing the signal-to-noise ratio (SNR). Moving the filter to 5 ps

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visibly distorts the fit residuals, indicating a poor quality of fit. Moving the cepstral filter to 15 ps causes a 1 ‰ to 2 ‰ shift in the retrieved δD and decreases the SNR. Therefore, we can reasonably attribute up to a 2 ‰ error to the cepstral fit technique.

250 As discussed in Section 3.2, temperature retrieval error has a considerable effect on the retrieved DCS δD value. This correlation is likely driven by the relative temperature dependence of HDO and H₂O absorption cross-sections in HITRAN. After performing the linear correction from Section 3.2, the residual temperature difference between the DCS and the NOAA PAO sensor is small (± 0.5 K, 1σ width). Therefore, the residual systematic error is estimated to be limited at the 5 ‰ level based on temperature sensitivity of the fitted δD value (*i.e.*, c_2 from Section 3.2). Some of this ± 0.5 K temperature difference
255 is likely driven by near-surface temperature gradients (Burns et al., 2012) and therefore the quoted 5 ‰ error is most likely an overestimate. Although the temperature may also vary slightly across the path, we extract the path-averaged temperature and are mostly insensitive to these gradients (Malarich and Rieker, 2021). In principle, an error in the HITRAN database for the relative pressure dependence of HDO and H₂O absorption cross-sections would similarly lead to an error in δD . However, by fitting at different pressures, we find a 1 ‰ error in pressure leads to a negligible 0.1 ‰ shift in the retrieved isotopologue
260 ratio. Finally, relative error in the HITRAN linestrengths for H₂O and HDO concentration will correspond directly to an error in δD . As noted earlier, estimated errors for HITRAN are 3 ‰ for HDO and <1 ‰ for H₂O in this region (Devi et al., 2017; Loos et al., 2017), giving an upper limit of 30 ‰ for the model error in δD . We do compare the absolute H₂O retrieval from the DCS to the co-located CRDS point sensor, finding agreement to within ~ 1 ‰; however, the CRDS calibration has not been verified, so we cannot quantify this error exactly. Finally, we note that HITRAN isotope ratios are slightly different than the
265 VSMOW scale, so we applied a small shift (~ 1 ‰) from the HITRAN-extracted δD value to scale to VSMOW (Griffith, 2018).

4 Comparison between DCS and NEON network sites

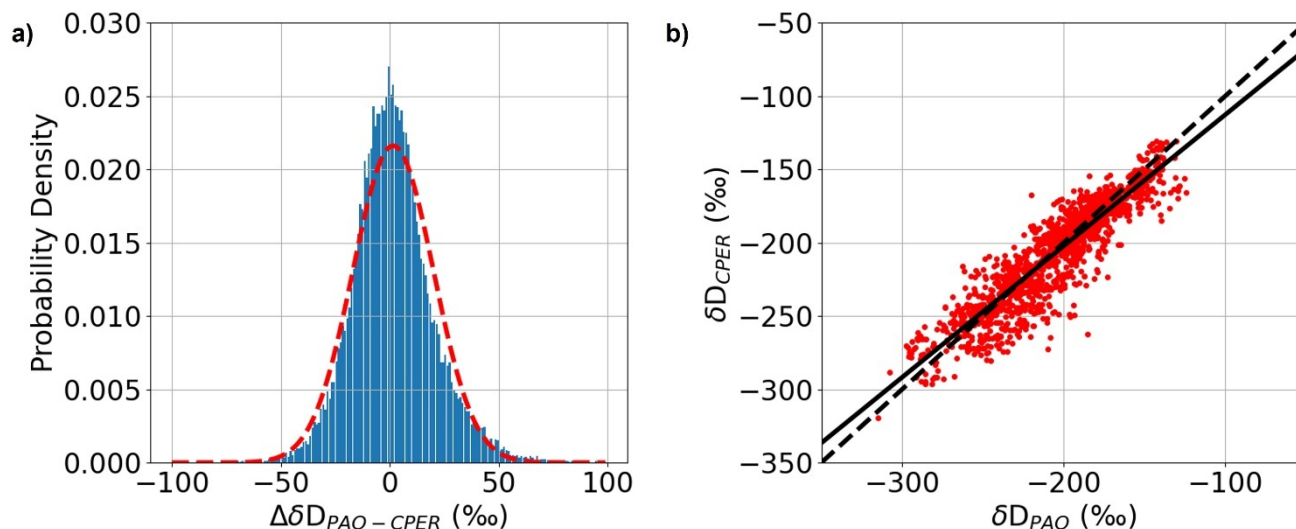
4.1 Comparison of δD at PAO and CPER

As shown in Figure 1, the closest and most similar NEON site to PAO is CPER. CPER is about 73 km north of PAO, located
270 near the Colorado-Wyoming border on the northern slope of the South Platte River Basin. Both PAO and CPER are located at ~ 1500 m elevation in high plains ecosystems (National Land Cover Database grassland/herbaceous class) with cold semi-arid Köppen classification. PAO is surrounded by oil and gas and agricultural development while CPER is situated far away from most industries and farms.

In general, temperatures at PAO and CPER are well correlated with both sites ranging from -20° C to 28° C during
275 the measurement period. Absolute concentrations of H₂O at PAO and CPER are also well correlated, falling between a range from 900 $\mu\text{mol/mol}$ to 12500 $\mu\text{mol/mol}$ during the measurement period. Because PAO and CPER are located on opposite sides of the South Platte River Basin, the two sites do experience differences in microscale and mesoscale meteorology, which leads to a difference in average surface wind patterns (Johnson and Toth, 1981). At the synoptic scale (~ 300 -3000 km scale), the



280 sites are close enough to experience atmospheric flow from mostly similar sources, which was verified using atmospheric trajectory simulation in HYSPLIT (Stein et al., 2015). Thus, we expect that the δD values at PAO and CPER should mostly agree, albeit with some potential differences due to the different micro- and mesoscale meteorology as well as occasional differences in synoptic flow.



285 **Figure 6:** a) Histogram of $\Delta\delta D_{PAO-CPER} = \delta D_{PAO} - \delta D_{CPER}$. Red dashed line: fit to Gaussian model (mean = 1.6 ‰, 1σ width = 18 ‰). b): PAO δD versus CPER δD . Black solid line: linear fit with slope = 0.90 ‰/‰, Intercept = -23 ‰, $R^2 = 0.83$. Black dashed line: 1 to 1 correlation line.

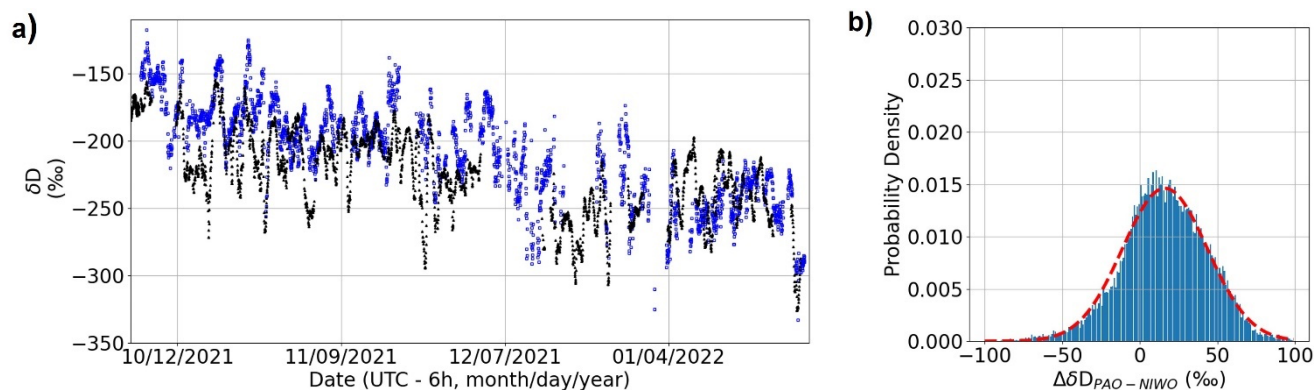
The histogram of $\Delta\delta D_{PAO-CPER} = \delta D_{PAO} - \delta D_{CPER}$ over the measurement period is shown in Figure 6a and a correlation plot between δD_{PAO} and δD_{CPER} is shown in Figure 6b. Both analyses validate the accuracy of the DCS data. A Gaussian fit to the histogram data of Fig. 6a yields a mean for $\Delta\delta D_{PAO-CPER}$ of <2 ‰ and a standard deviation of 18.2 ‰. The slope of the correlation in Fig. 6b is 0.90 with an $R^2 = 0.84$. This level of agreement between isotope ratios measured by a point sensor and an active open-path sensor has not been demonstrated previously to our knowledge. The agreement in terms of mean and standard deviation is comparable or better than observed for vertical column densities measured using solar-looking FTIR and satellite reflectance measurements (Schneider et al., 2020). We note that the $\Delta\delta D_{PAO-CPER}$ distribution does not perfectly follow a normal distribution but is better approximated by a generalized logistic distribution indicating that the distribution is rather heavy-tailed compared to a normal distribution. This could be from an uncorrected systematic in one of the measurements or a true difference between isotope values at PAO and CPER.

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4.3 Comparison of δD at PAO and NIWO

300 To illustrate the large spatial variability in δD that is possible near the intersection of mountain and plains ecosystems, we also compare the measurements at the PAO site to the NIWO NEON site as shown in Figure 7. The NIWO site sits at ~ 3500 m in the Front Range near the Continental Divide about ~ 20 km west of Boulder and ~ 75 km west of PAO. The NIWO site has a subarctic Köppen classification and an evergreen forest and grassland/herbaceous class according to the National Land Cover Database. Thus, even though NIWO is located a similar distance from PAO as CPER, there is a ~ 2000 m difference in elevation
 305 between PAO/CPER and NIWO, which should lead to higher δD values at PAO relative to NIWO. In addition, subarctic NIWO is more strongly influenced by the free troposphere (which has significantly lower δD) leading to time periods with large disagreement between NIWO and PAO. This behaviour appears in both the histogram of differences (Figure 7a) and the raw time series (Figure 7b). Overall, we observe a ~ 15 ‰ offset in $\Delta\delta D_{PAO-NIWO}$. However, the $\Delta\delta D_{PAO-NIWO}$ distribution deviates from a normal distribution with a mode that is closer to zero offset than the estimated mean. It is likely that the skew and offset of the distribution arises from the varying influence of the free troposphere on NIWO driven by events such as
 310 mountain upslope flow (Baumann et al., 1997). This is also observed in the time series where there are time periods with correlated δD values between PAO and NIWO. In the future, DCS can help provide a denser network of measurements at varying altitudes across the Front Range in order to better characterize the atmospheric interactions between the higher altitude ecosystems and the plains ecosystems below.



315

Figure 7: a) Time series of δD at PAO (blue) and NIWO (black). b) $\Delta\delta D_{PAO-NIWO}$ histogram. Red dotted line: fit to Gaussian model (mean = 16 ‰, 1σ width = 27 ‰).



5 Temporal variability of δD at PAO and CPER

320 5.1 Diurnal and Seasonal Patterns of δD

Diurnal profiles can provide information about local sources of water vapor as well as persistent atmospheric flow patterns such as katabatic winds (Bréant et al., 2019). Here, we show that the open-path DCS measurements have sufficient precision/accuracy to track changes in diurnal patterns of δD . Because significant day-to-day mean variability occurs in δD due to synoptic flow, a stacked or ensemble diurnal profile is generated by first subtracting the daily mean δD and then looking
325 at the diurnal variation in mean-subtracted δD . These profiles are shown in Figure 8 for each month of the measurement period (October 2021 through January 2022) for both PAO and CPER. As with the histogram, the diurnal profiles generally agree between the two sites, which highlights the similarity between the two locations and shows that the open-path DCS can resolve diurnal patterns. There are some small differences in the diurnal patterns during October and December with PAO showing slightly lower δD values in the morning compared to CPER. This trend could be due to differences in downslope flows at PAO
330 versus CPER (Baumann et al., 1997).

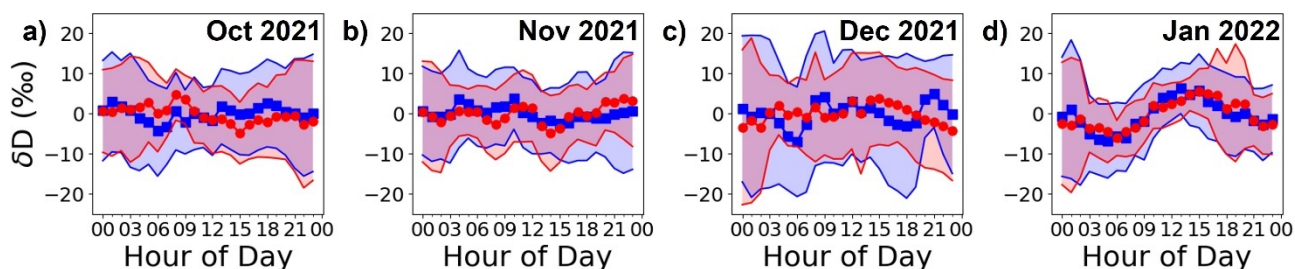
Interestingly, the stacked diurnal profiles show a distinct change between the first three months and January 2022. In particular, the January 2022 period shows a clearer diurnal trend that is positively correlated with ambient air temperature. A detailed analysis of diurnal profiles is complicated: in general, the profiles are determined by an interplay between local fluxes (evapotranspiration, sublimation, etc), atmospheric flow patterns such as mountain flow, and boundary layer dynamics leading
335 to mixing/entrainment with the free troposphere (Welp et al., 2008; Wen et al., 2010; Noone et al., 2013; Su et al., 2022; Steen-Larsen et al., 2013; Bréant et al., 2019). Diurnal patterns with zero or negative temperature correlations like those shown from October 2021 through December 2021 have been observed in both rural and urban environments (Welp et al., 2008; Wen et al., 2010) and were previously hypothesized to be the result of vertical mixing in planetary boundary layer (PBL) coupled with entrainment from the free troposphere. Diurnal profiles with a positive temperature correlation are often observed in colder
340 climates (Steen-Larsen et al., 2013; Bréant et al., 2019). We do not have enough auxiliary measurements to conclusively identify the cause of this change, but we will note that the period up to the end of December 2021 was extremely dry and somewhat windy and was followed by a snowy and colder January 2022, which suggests that the changes in diurnal profiles might be caused by some combination of meteorological changes related to the colder/snowier weather or larger local water vapor fluxes due to snow cover.

The accuracy of the DCS measurements also allows us to track seasonal changes in δD . Unlike the hourly-averaged δD , the daily-averaged δD trends strongly with daily-averaged temperature (see Figure 9). This correlation holds over the whole measurement period and yields a slope of 3.7 ± 0.2 ‰/°C which is consistent with global observations of precipitation in the Global Network of Isotopes in Precipitation (GNIP) network (Gat et al., 2000). A rolling ordinary least squares regression was also performed on the unaveraged δD and temperature data from PAO to confirm the dependence of δD on temperature.
350 This analysis shows that on average there is a 4.1 ± 0.9 ‰/°C dependence of δD on ambient air temperature at PAO. The average δD value at 0°C was -227 ± 10 ‰ which is much lower than expected based on temperature effects alone for coastal areas but



could possibly be explained by a combination of the altitude and continental effects (Gat et al., 2000; Schneider et al., 2020). Our multi-month measurement campaign has allowed us to view the transition from autumn-type diurnal cycles in the Front Range Urban Corridor to winter-type diurnal cycles with the cycles at PAO and CPER agreeing well both qualitatively and quantitatively without any recalibration necessary for the DCS measurements. The close agreement between the sites suggests boundary layer dynamics and temperature effects are the dominant drivers of δD variation in Northern Colorado plains ecosystems.

355



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Figure 8: Stacked diurnal cycles for δD at PAO (blue) and CPER (red) with mean value subtracted. Shaded regions indicate standard deviation of values in each hourly bin. a) October 2021. b) November 2021. c) December 2021. d) January 2022.

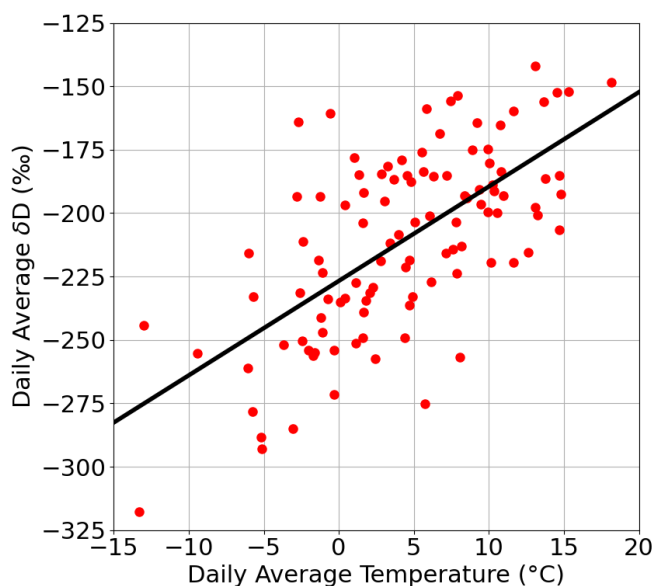


Figure 9: Correlation between daily average temperature and δD at PAO site. Slope = 3.7 ± 0.2 ‰/°C.

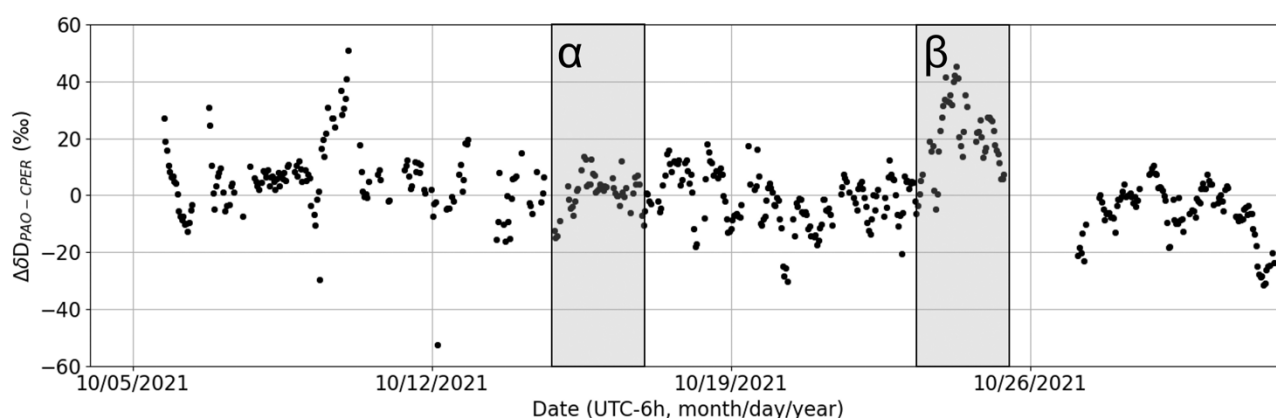
5.2 Temporal differences between PAO and CPER

The high precision of the DCS measurements also provides an opportunity to look for meteorological differences in δD between PAO and CPER. Figure 10 shows $\Delta\delta D_{\text{PAO-CPER}}$ over October 2021 resampled to 3000 s time resolution. This month

365



370 corresponds to the most continuous high-SNR section of PAO data and thus is the best section of data for probing temporal dynamics. In this time series, we see periods with high agreement (*e.g.*, the period marked α spanning 10/15/2021 to 10/16/2021) as well as periods with more marked disagreement (such as the period marked β spanning 10/23/2021 to 10/24/2021). This disagreement between PAO and CPER is likely indicative of differences in synoptic flow at the two sites during this period as discussed below.



375 **Figure 10: $\Delta\delta D_{\text{PAO-CPER}}$ versus time. Very good agreement is observed during the period denoted α (first grey band), while fairly strong disagreement is seen during the period denoted β (second grey band), reaching more than 20 ‰ for about a day.**

Differences in synoptic flow can be probed using a two-source mixing model or Keeling analysis (Noone, 2012; Keeling, 1958). When two air masses with different δD and H_2O content mix (*e.g.*, moist ocean air mixing with dry free tropospheric air), δD follows a linear relationship when plotted against $1/[\text{H}_2\text{O}]$ according to the following form (Noone, 2012):

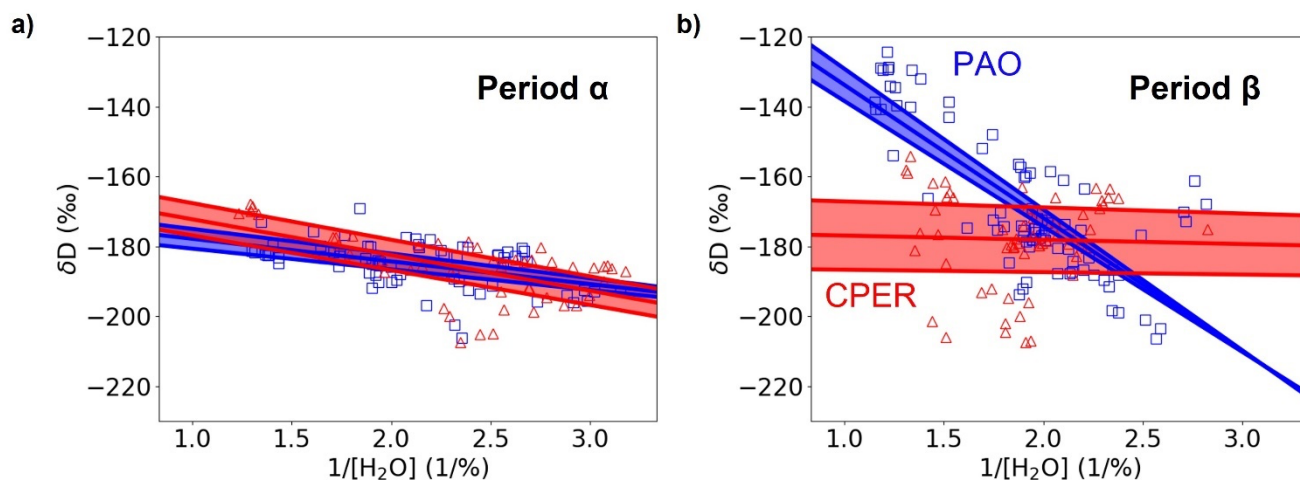
380

$$\delta D = [\text{H}_2\text{O}]_o (\delta D_o - \delta D_{\text{source}}) \frac{1}{[\text{H}_2\text{O}]} - \delta D_{\text{source}}, \quad (3)$$

385 where $[\text{H}_2\text{O}]_{\text{source}}$ and δD_{source} are the H_2O concentration and δD for the source air mass (*e.g.*, moist ocean air) and $[\text{H}_2\text{O}]_o$ and δD_o are the H_2O concentration and δD for the dry air mass. We can extract δD_{source} either by plotting δD versus $1/[\text{H}_2\text{O}]$ and finding the y-intercept or by fitting the hyperbolic relation described in Equation 3. Figure 11 shows comparisons between δD versus $1/[\text{H}_2\text{O}]$ measured at two times (labelled periods α and β in Figure 10). During period α , both CPER and PAO show the same mixing behaviour; however, during period β , the mixing behaviour is significantly different with notably different slopes and intercepts. A limited analysis of atmospheric transport at PAO and CPER suggests that synoptic flow typically comes from



the Pacific Northwest, from the Southwest, or a from combination of these two. A “latitude effect” has been previously
 390 observed where air from higher latitudes tends to be more depleted of heavy isotopologues (Araguás-Araguás et al., 2000;
 Schneider et al., 2020). During period β , the δD_{source} at PAO was 75 ‰ higher than at CPER, which suggests that the air at
 PAO more likely came from the Southwest (higher δD_{source}) whereas the air at CPER more likely came from the Pacific
 Northwest or a combination of sources. This provides further evidence that the synoptic-scale atmospheric transport at PAO
 and CPER can be different.



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Figure 11: δD versus $1/[H_2O]$ over shorter time periods at PAO (blue squares) and CPER (red triangles) suggest intermittent
 meteorology differences between the sites. a) Data over the time period α of good agreement. The fit to Eq. (3) gives $\delta D_{\text{source}} = -168 \pm 3$
 ‰ for the PAO data and $\delta D_{\text{source}} = -162 \pm 5$ ‰ for the CPER data. b) Data over the time period β of disagreement. The fit to Eq. (3)
 400 gives $\delta D_{\text{source}} = -93 \pm 7$ ‰ for the PAO data and $\delta D_{\text{source}} = -176 \pm 10$ ‰ for the CPER data. Shaded regions indicate confidence interval
 linear fit.

6 Discussion/Conclusion

The observed differences in δD between sites separated by 75 km motivates the construction of denser isotopologue
 measurement networks for benchmarking models of water transport. However, accurate/precise measurement of isotopologues
 405 in lower tropospheric water vapor is a challenging task for traditional open-path techniques such as FTIR and CW laser
 spectroscopy. Point sensors can provide accurate measurements but lack the operational flexibility of open-path techniques
 and require regular recalibration. We have shown that MIR DCS is capable of accurate open-path measurement of H_2O and
 HDO concentrations in real world settings for extended periods of time. Our DCS δD precision of 1.2 ‰ at 3000 seconds
 allows us to characterize shifts in ambient δD levels. After correcting the data for a small linear error in retrieved temperature
 410 due to errors in the spectral database, the DCS-based results taken at PAO matches those taken at CPER with a mean difference
 of 1.6 ‰ and a root mean square error of 18 ‰ over a three-month period. It is likely that some of the differences are true
 differences driven by meteorology. In the future, DCS laboratory spectroscopy can facilitate improved accuracy for the water



absorption lineshape models used to extract isotopologue ratios in the MIR with DCS (Hayden et al., 2019). With the current precision, it should be possible to measure over multiple vertically or horizontally separated paths with one DCS instrument
415 (Coburn et al., 2018; Herman et al., 2021) to understand influences on isotopologue ratios from local evaporative flux (Welp et al., 2008). Such experiments allow for precise extraction of the local isotopologue ratio for the evapotranspirative flux (δD_{ET}) and can be used to better understand the water cycle of natural environments like estuaries or man-made environments like reservoirs and agricultural operations (Welp et al., 2008; Barrie et al., 2015).

In addition to $H_2^{16}O$ and HDO, open-path DCS simultaneously provides concentration for other atmospheric species
420 and can help answer a variety of questions related to atmospheric transport. Future ultra-broadband DCS systems based on ultra-short pulsed nonlinear optics (Lesko et al., 2021; Zhou et al., 2020) will enable simultaneous coverage of $H_2^{16}O$, $H_2^{18}O$ and HDO for characterization of d-excess (defined as the deviation from global meteoric water line) over open paths (Dansgaard, 1964; Craig et al., 1965). When combined with information on combustion sources from DCS measurements of ambient CO , CO_2 and HCHO, the d-excess record may help confirm carbon inventories using open-path DCS systems (Xing
425 et al., 2020). Also, water isotopologue measurements at NIWO and PAO combined with DCS multispecies measurements of trace gases including ammonia, greenhouse gases and volatile organic compounds could improve our understanding of upslope pollution events that contribute to nitrification and ozone production along the Front Range Urban Corridor and the Rocky Mountain National Park area (Piña et al., 2019; Baumann et al., 1997). We also highlight the important role of ground-based sensors in the calibration of satellite remote sensing measurements of HDO vertical columns (Schneider et al., 2020). In the
430 future, MIR DCS could provide a mobile, reconfigurable platform to supplement the current fixed Total Carbon Column Observing Network (TCCON) and Network for the Detection of Atmospheric Composition Change (NDACC) measurements used for calibration of inversion algorithms used to extract height-dependent water isotopologue ratios (Wunch et al., 2011; Zhou et al., 2021).

Data Availability.

435 Water vapor isotopologue and meteorological data is available in CSV format through the NIST MIDAS Database (DOI to be provided).

Author Contributions.

D.I.H., F.R.G., E.B., G.M., N.M., B.R.W., I.C., and K.C.C. helped develop the hardware and set up the field measurement. D.I.H., G.M., and K.C.C. collected the data. D.I.H., F.R.G., G.M., N.M., and K.C.C. analysed the data. D.I.H., N.R.N., I.C.
440 and K.C.C. wrote the manuscript. All co-authors provide comments on the manuscript.

Competing Interests.

The contact author has declared that none of the authors has any competing interests.



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Review Statement.

455 TBD

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