

Impact of permeability evolution in igneous sills on hydrothermal flow and hydrocarbon transport in volcanic sedimentary basins

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Abstract:

Sills emplaced in organic-rich sedimentary rocks trigger the generation and migration of hydrocarbons in volcanic sedimentary basins. Based on seismic and geological observations, numerical modeling studies of hydrothermal flow around sills show that thermogenic methane is channeled below the intrusion towards its tip, where hydrothermal vents nucleate and transport methane to the surface. However, these models typically assume impermeable sills and ignore potential effects of permeability evolution in cooling sills, e.g., due to fracturing. Here, we combine a geological field study of a volcanic basin (Neuquén Basin, Argentina) with hybrid FEM/FVM numerical modeling of hydrothermal flow around a sill, including hydrocarbon generation and transport. Our field observations show widespread veins within sills composed of graphitized bitumen and cooling joints filled with solid bitumen or fluidized shale. Raman spectroscopy indicates graphitization at temperatures between 350-500°C, suggesting fluid flow within the intrusions during cooling. This finding motivates our modeling setup, which investigates flow patterns around and through intrusions that become porous and permeable upon solidification. The results show three flow phases affecting the transport of hydrocarbons generated in the contact aureole: (1) Contact-parallel flow toward the sill tip prior to solidification, (2) upon complete solidification, sudden vertical “flushing” of overpressured hydrocarbon-rich fluids from the lower contact aureole towards and into the hot sill along its entire length, and (3) stabilization of hydrocarbon distribution and fading hydrothermal flow. In low-permeability host rocks, hydraulic fracturing facilitates flow and hydrocarbon migration toward the sill by temporarily elevating porosity and permeability. Up to 7.5% of the generated methane is exposed to temperatures >400°C in the simulations and may thus be permanently stored as graphite in or near the sill. Porosity and permeability creation within cooling sills may impact hydrothermal flow, hydrocarbon transport and venting in volcanic basins, as it considerably alters the fluid pressure configuration, provides vertical flow paths, and helps to dissipate overpressure below the sills.

43 **1 Introduction**

44 Sill intrusions emplaced in sedimentary rocks strongly influence generation and migration of
45 hydrocarbons and greenhouse gases in volcanic sedimentary basins. If sill intrusions are
46 emplaced in organic-rich strata, they trigger contact-metamorphic reactions (e.g., organic
47 matter transformation), overpressure generation and hydrothermal fluid flow in the
48 surrounding strata (Einsele et al., 1980; Aarnes et al., 2012). Many recent studies of volcanic
49 sedimentary basins investigate how such processes may cause the formation of hydrothermal
50 vent complexes, which facilitate greenhouse gas release to the atmosphere and can thus drive
51 global climate change (Svensen et al., 2004; Aarnes et al., 2010; Aarnes et al., 2012; Iyer et
52 al., 2013; Iyer et al., 2017; Galerne and Hasenclever, 2019). Additionally, the same processes
53 can be critical factors for hydrocarbon generation and migration in igneous petroleum
54 systems containing sills emplaced within shale formations (Senger et al., 2017; Spacapan et
55 al., 2020).

56 Hydrothermal flow in response to intrusion of magma into sedimentary host rocks has been
57 investigated for decades. The magmatic heat input leads to several temperature-dependent
58 processes that promote strong fluid pressure increase, which drive fluid flow (Einsele et al.,
59 1980; Delaney, 1982). These processes include for instance thermal fluid expansion, mineral
60 dehydration, organic matter transformation into hydrocarbon generation, and pore space
61 reduction due to mineral precipitation (Einsele et al., 1980; Delaney, 1982; Aarnes et al.,
62 2010; Townsend, 2018). When the rate of overpressure generation is larger than flow-driven
63 pressure dissipation, e.g., in low-permeable rocks like shale, hydraulic fractures form and
64 locally enhance fluid flow and pressure release (Jamtveit et al., 2004; Aarnes et al., 2012;
65 Kobchenko et al., 2014; Panahi et al., 2018; Rabbel et al., 2020). This process may lead to
66 explosive hydrothermal vents, which are present in several volcanic basins (Nermoen et al.,
67 2010; Aarnes et al., 2012; Iyer et al., 2017). In general, Ingebritsen et al. (2010) highlighted
68 the deciding role of permeability structure for magmatic hydrothermal systems, where
69 permeability of 10^{-16} m² represents the approximate boundary between convection and
70 conduction dominated systems.

71 Numerical models simulate these coupled processes to understand hydrothermal flow
72 dynamics and the associated hydrocarbon migration (usually represented as methane carried
73 in the hydrothermal fluids). Typically, these models investigate potential vent formation
74 around sills emplaced in organic-rich sediments. Such simulations usually assume that sills
75 are impermeable and show that vents seem to preferentially form at the inclined tips of large
76 “saucer-shaped” sills, because fluids get trapped under the sill and migrate towards their tips.
77 This situation favours both fast fluid pressure build-up below the sills and focussed fluid
78 migration towards the tips (Iyer et al., 2013, Iyer et al., 2017; Galerne and Hasenclever,
79 2019).

80 However, observations from several volcanic sedimentary basins indicate that the assumption
81 of impermeable sills is not generally valid. Sills often host fracture networks including
82 different fracture types. These form shortly after solidification of the magma and may include
83 columnar cooling joints or fractures related to thermal contraction or hydraulic fracturing
84 during hydrothermal activity (Senger et al. 2015, Witte et al. 2012, Rabbel et al., 2021).
85 Multiple studies have provided evidence that such fracture networks may be open and can
86 contain water (Chevallier et al., 2004) or hydrocarbons and act as fluid pathways or even
87 fractured reservoirs (Mark et al., 2018; Schofield et al., 2020; Spacapan et al., 2020).

88 Mainly based on evidence from field and subsurface data, several studies hypothesized that
89 cooling-related fracturing in sills creates an early migration pulse of fluids into the sill,

90 although the thermal regime of such a pulse is under debate (Witte et al., 2012, Spacapan et
91 al. 2020, Rabbel et al. 2021). In addition, Spacapan et al. (2019) noted the absence of
92 hydrothermal vents around the sills in the Neuquén Basin and suggested that pore pressures
93 were not high enough to create such features. A recent study on the Karoo Basin in South
94 Africa estimates that half of the thermogenic gas mobilized in the contact aureole of flat sills
95 may enter the sill through cooling joints (Lenhard et al., 2023). Although based on field
96 evidence, these geological models remain qualitative in terms of the physical process
97 dynamics, and a dedicated study to investigate quantitative hydrothermal flow and
98 hydrocarbon migration around and in fractured, permeable sills is currently missing.

99

100 In this study, we combine a field study from the Neuquén Basin, Argentina, with numerical
101 modelling to: (1) investigate if and in which thermal conditions opening of cooling joints may
102 trigger an early hydrocarbon migration pulse into the sill, and (2) assess the impact of
103 porosity and permeability generation in sills on the hydrothermal flow and hydrocarbon
104 migration and storage in comparison to systems with impermeable sills. The northern
105 Neuquén Basin provides well documented examples of sills with extensive cooling joint
106 networks emplaced in organic-rich shale, many of which are commercial oil reservoirs
107 (Rodríguez Monreal et al., 2009; Witte et al., 2012; Spacapan et al., 2020). Additionally, the
108 strong thermal impact of intrusions on host rock maturation in these systems is well
109 documented (Rodríguez Monreal et al., 2009; Spacapan et al., 2018). We first present
110 geological evidence for hydrocarbon transport through fractured sills in a hydrothermal
111 environment. Our field observations motivate a numerical modelling study to investigate the
112 influence of cooling joint formation, i.e., permeable sills, on the hydrothermal flow patterns
113 and hydrocarbon migration. We perform simulations for sills emplaced in host rocks of
114 different permeability to be able to discuss the effect for different geological settings. By
115 integrating the simulations results with geological evidence, we show how permeable sills
116 affect hydrothermal flow in volcanic sedimentary basins as well as the fate of hydrocarbons
117 generated by contact-metamorphism.

118

119 **2 Geological observations**

120 **2.1 Geological setting**

121 The study area is located around the Río Grande Valley (RGV) in the northern Neuquén
122 Basin, Argentina, about 100 km south of the town of Malargüe (Figure 1). The Neuquén
123 Basin initially formed as a series of isolated half-grabens during the late Triassic to early
124 Jurassic (Howell et al., 2005). During the middle Jurassic to early Cretaceous, these
125 depocenters coalesced during thermal subsidence, forming a large shallow-marine basin. This
126 phase included the deposition of the Vaca Muerta and Agrío formations, which comprise
127 several hundreds of meters of calcareous, organic-rich shale and form two important source
128 rock formations for presently exploited petroleum systems (Kietzmann et al., 2014). From the
129 early Cretaceous, the basin developed into a foreland basin in response to the compressive
130 tectonic regime of the Andean orogeny. This led to inversion of the Triassic normal faults and
131 generation of a series of fold-thrust belts along the western basin boundary (Maceda and
132 Figueroa, 1995; Yagupsky et al., 2008).

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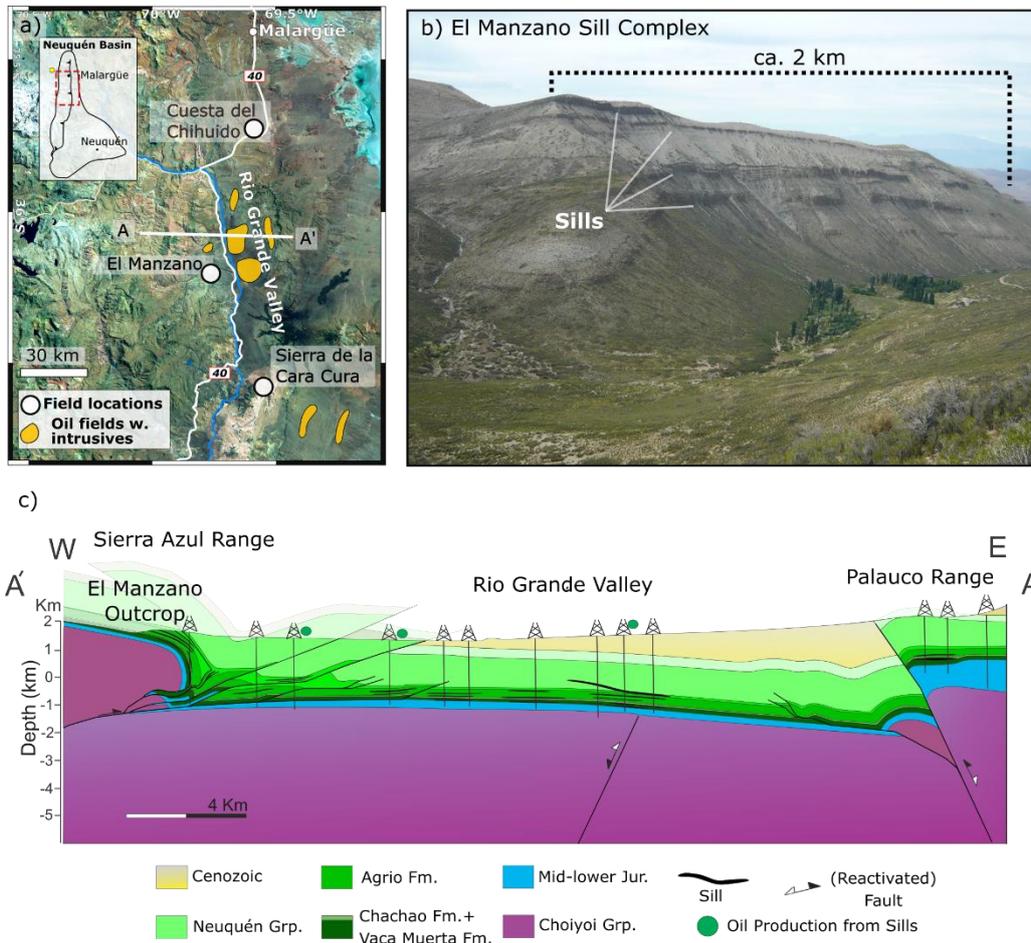


Figure 1. (a) Satellite image (Landsat, courtesy USGS) of the study area in the northern Neuquén Basin including field localities (white dots) and producing igneous petroleum systems (green areas). (b) View of the El Manzano sill complex outcrop in the Sierra Azul range (Photo: F. Soto). (c) E-W structural geological section illustrating the relation between subsurface and outcropping sill complexes (courtesy. J. B. Spacapan)

134 In addition to tectonic deformation, the northern Neuquén Basin experienced intense
135 magmatic activity that formed a series of volcanic plateaus and widespread magmatic
136 intrusions within the sedimentary succession (Kay et al., 2006). In the study area, two main
137 eruptive cycles, termed Molle and Huincán eruptive cycles, occurred in the late Oligocene-
138 mid Miocene and the late Miocene-Pleistocene, respectively (Combina and Nullo, 2005).
139 These events led to the emplacement of extensive andesitic and basaltic sill complexes
140 predominantly in the Vaca Muerta and Agrío Formations, but also in overlying gypsum and
141 sandstone units (Spacapan et al., 2020).

142 In the RGV, heavily fractured sills emplaced in the Vaca Muerta and Agrío formations
143 constitute the reservoirs in an actively producing igneous petroleum system (Schiuma, 1994;
144 Witte et al., 2012). Spacapan et al. (2020) reported 2 to 27 m thick sills in these shale
145 formations, but up to 54 m within shallower clastic sediments of RGV. Other studies from the
146 northern Neuquén Basin report thick laccolith intrusions (>100 m) acting as reservoirs
147 (Rodríguez Monreal et al., 2009). Spacapan et al. (2018) showed that the heat input provided
148 by the sills matured the Vaca Muerta and Agrío shale formations, which otherwise show very
149 low thermal maturity at burial depths of ca. 2-2.5 km. The established model for the
150 formation of these igneous reservoirs includes that the intrusions developed interconnected
151 cooling joint networks, which subsequently stored the generated hydrocarbons (Witte et al.,
152 2012; Spacapan et al., 2020).

153 In addition to the subsurface sill complexes, thrust tectonics brought to surface exceptional
154 analogue outcrops in the surrounding mountain ranges including the Sierra Azul, Sierra Cara
155 Cura and Cuesta del Chihuido (Figure 1). Several studies describe these localities, which
156 offer easy access to sills emplaced in the Vaca Muerta and Agrío formations (Spacapan et al.,
157 2017; Rabbel et al., 2018; Rabbel et al., 2021). Especially the km-scale outcrops at El
158 Manzano (Figure 1b) and Sierra Cara Cura constitute direct analogues to the subsurface sill
159 complexes of RGV igneous petroleum system (Palma et al., 2019; Rabbel et al., 2021). These
160 three field localities are ideal case studies to reveal the interactions between igneous
161 intrusions and the petroleum system.

162

163 **2.2 Geological field observations**

164 During three field campaigns, we collected an extensive dataset at outcrops in El Manzano,
165 Sierra de la Cara Cura and Cuesta del Chihuido (Figure 1). We gathered ground-based and
166 drone digital photographs to document outcrop observations. Additionally, we collected over
167 100 rock samples from the intrusions, surrounding shale as well as various types of veins for
168 geochemical analyses. Note that a more comprehensive description of the field study is
169 presented by Rabbel et al. (2021).

170

171 **2.3 Observations of hydrocarbons inside and around sills**

172 Outcropping sills in all three localities feature solid bitumen and black shale inside the
173 fracture network of the sills. At Cuesta del Chihuido, both the side and roof of thin sills are
174 exposed, and the side view reveals upwelling dykelets of black shale (Vaca Muerta
175 Formation) entering the sill from the bottom contact (Figure 2a). The top view of the same
176 sill shows the entire polygonal cooling joint network with a black fill of the same material
177 (Figure 2a, b). Brecciated igneous material often surrounds the dykelets where they enter the
178 intrusion.

179 The larger sills at El Manzano (Sierra Azul) and Sierra Cara Cura also show widespread
180 bitumen in the fracture network of the sills, but at a much larger scale (Figure 3). We observe
181 arrays of 1-up to 50 cm thick and >10 m high bitumen dykes or veins (Figure 3a). Here, the
182 bitumen dyke cuts across the contact aureole and enters the sill intrusion. We find exposures
183 of similar structures where the sill interior is accessible (Figure 3b). The sill appears heavily
184 fractured in addition to preexisting cooling joints, and solid bitumen or calcite fill nearly all
185 fractures. On closer inspection, the bituminous material in these veins has a shiny and fibrous
186 texture.

187 At an exposed sill tip at El Manzano, we also observe that several cm thick bitumen veins
188 appear to be concentrated along the tip contact, where they mutually cross-cut with calcite
189 veins of at least similar thicknesses (Figure 3c, d). These calcite veins have cm-scale pores,
190 which occasionally contain solid bitumen themselves and release strong hydrocarbon smell
191 when the vein is broken up.

192

193 **2.4 Bitumen characterization**

194 The fibrous texture of the observed bitumen within the sills is intriguing and we hypothesized
195 that it may be much higher-grade bituminous material than that described in the Neuquén
196 basin and commonly attributed to regional burial (Parnell et al., 1995; Cobbold et al., 1999;
197 Zanella et al., 2015).

198 In the field, we tested this hypothesis by measuring the resistivity of the fibrous bitumen with
199 a hand-held multimeter. Graphitization of bitumen significantly changes the electric
200 resistivity of the material: amorphous solid bitumen is very resistive and used as an electric
201 insulator in industry applications (Hays et al., 1967), while graphite is an excellent conductor.
202 Qualitative on-site resistivity measurements showed that the fibrous bitumen conducts
203 electric currents well, i.e., within the detection limit of a standard multimeter, suggesting
204 significant graphitization.

205 In addition, we applied Raman spectroscopy to better constrain the nature of the solid
206 bitumen and its thermal history. Raman spectroscopy provides positions and relative
207 intensities of spectral peaks characterizing carbonaceous materials like bitumen, including D
208 (“disorder”) and G (“graphite”) peaks at 1345 cm^{-1} and 1585 cm^{-1} , respectively (Potgieter-

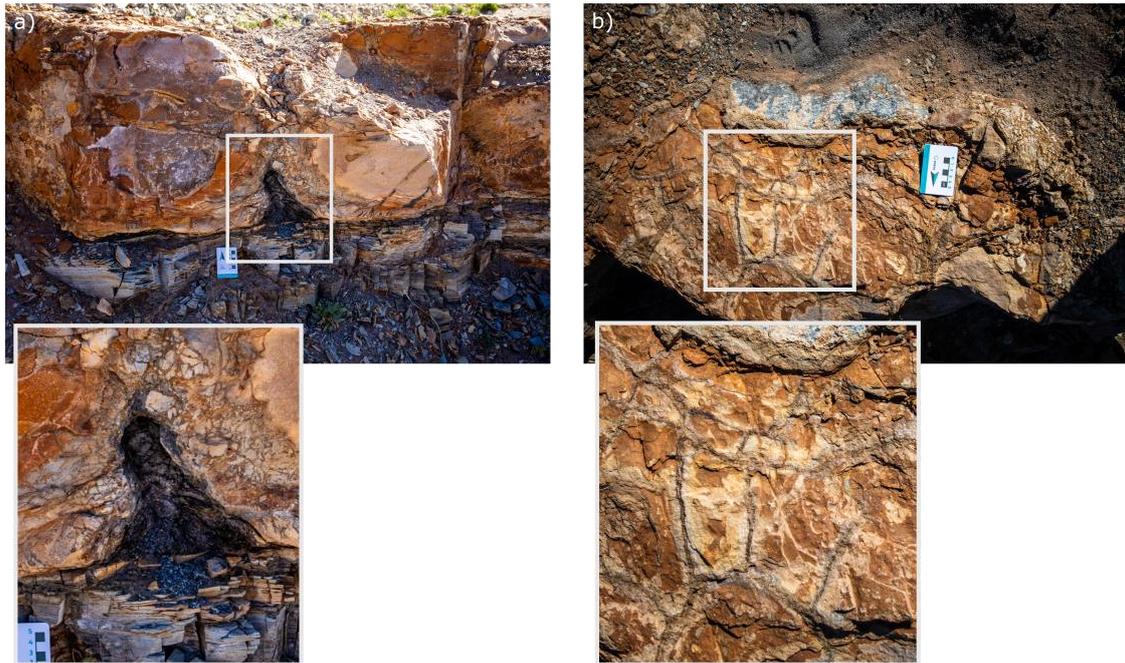
209 Vermaak et al., 2011; Rantitsch et al., 2016). The shape of the spectra and the D/G peak and
210 area ratios allow a classification of high-grade alteration of the bitumen to anthracite or
211 (semi-)graphite and may serve as a geothermometer for high-temperature regimes (Beysac et
212 al., 2002; Rantitsch et al., 2016). Due to the high temperatures within and around igneous
213 intrusions, we expect this method to give an indication on the degree of thermal alteration and
214 thus temperatures that the hydrocarbons experienced. Since Raman spectra can show varying
215 absolute intensities, we normalized each spectrum to the intensity of the respective G peak
216 (I_G) for visualization purposes.

217 Raman spectrograms of the sampled bitumen veins show very clearly developed G and D1
218 peaks and I_{D1}/I_G ratios of ca. 1 and 0.9, respectively (Figure 4). The D3 band between the
219 peaks is nearly absent in the sample from Sierra de la Cara Cura (from vein in Figure 3b),
220 while it is visible at low intensity in the presented sample from El Manzano (from vein in
221 Figure 3a). Note that both vein samples stem from the intrusion-host contact, and each veins
222 penetrate about 10 m into around 20 m thick sills.

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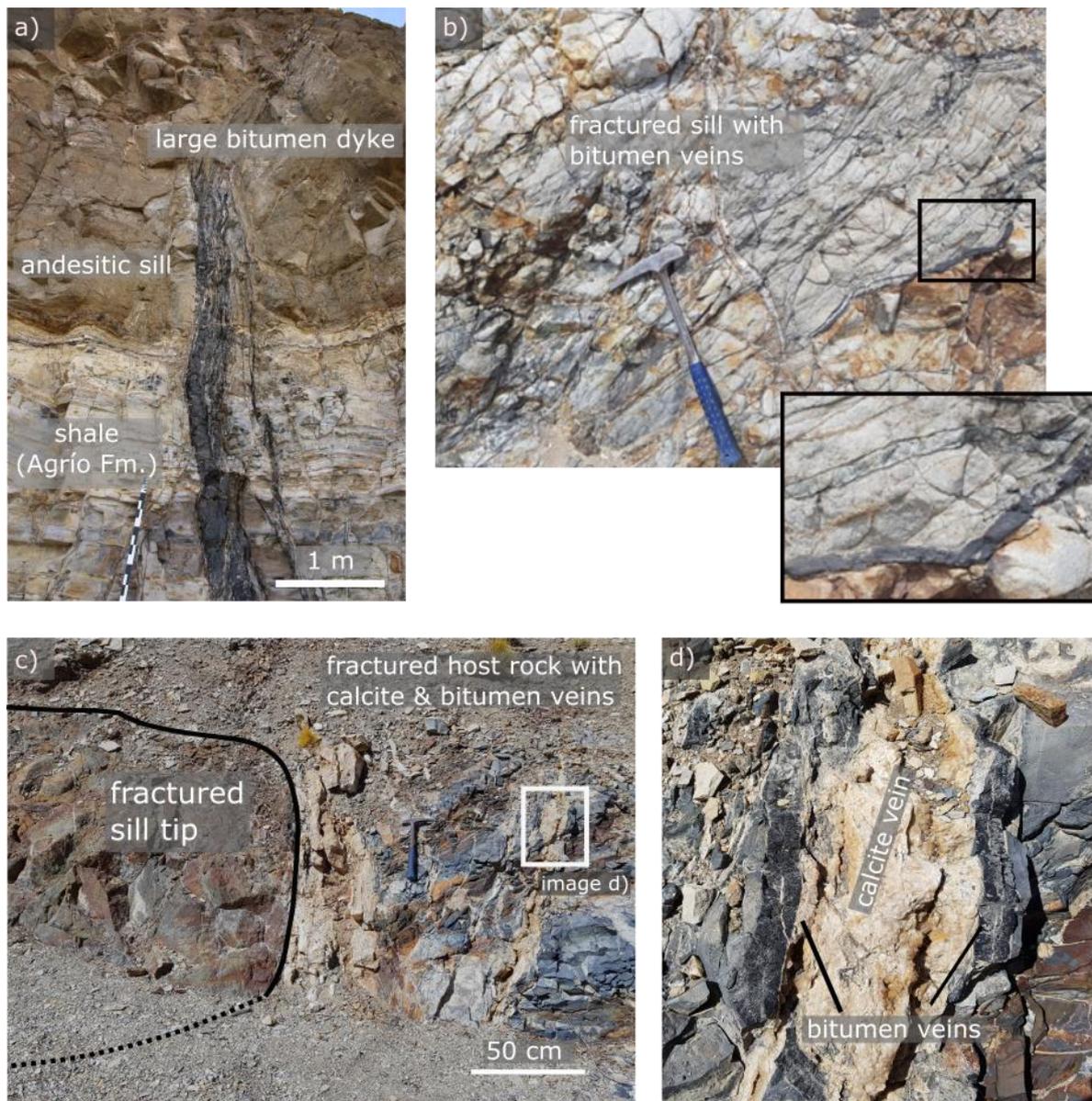
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226

227 Figure 2. Field observations of upwelling dykelets of liquefied shale and bitumen entering the
228 cooling joint network of a thin sill at Cuesta del Chihuido (ca. 30 cm thick). (a) Side view
229 showing the sediment-intrusion contact and dyke (Photo: D. Michelon), (b) top view
230 demonstrating black bituminous fill in the polygonal cooling joints.

231



233

234

235 Figure 3. Examples of dykes or veins of solid bitumen associated the sill intrusions. (a)
 236 Bitumen dyke at El Manzano of >10 m height and up to 0.5 m thickness originating in the
 237 aureole of Agrío Fm. and entering the sill through the bottom contact. (b) Fractured zone
 238 inside a sill at Sierra de la Cara Cura exhibiting many cm-scale veins of solid bitumen. (c, d)
 239 Exposed sill tip at El Manzano showing high concentration of fibrous bitumen and calcite
 240 veins in the contact in front of the tip.

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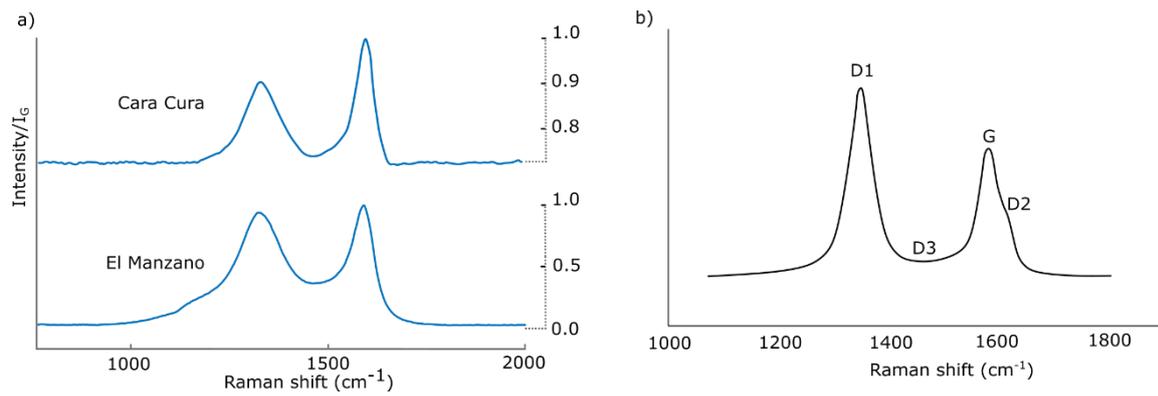
243 2.5 Composition and thermal implications of bitumen samples

244 We compare our Raman results with those measured in carbonaceous material from several
245 studies, where increased graphitization and metamorphism lead to well-developed, narrow
246 graphite (G) and disordered carbon (D1) peaks, weak or absent D3 bands, and I_G/I_D peak
247 ratios of <1 (Beysac et al., 2002; Kwiecinska et al., 2010; Rantitsch et al., 2016). Although
248 we did not perform quantification via peak-fitting, a qualitative comparison of our results
249 with highly metamorphosed sediments presented by Beysac et al. (2002, Fig. 6 and 11) leads
250 to estimated temperatures of 350-500°C for our samples (Figure). Hydrothermal
251 graphitization can occur along intrusion-sediment contacts at relatively shallow crustal levels
252 and requires temperatures of $\geq 400^\circ\text{C}$ (Buseck and Beysac, 2014). Hydraulic fracturing
253 focuses the flow of hydrothermal fluids oversaturated with CH_4 and/or CO_2 from which
254 crystalline graphite may precipitate (Rumble, 2014). This fits well with the observations that
255 the bitumen dykes in our study area consist of pure, often crystalline graphitic material and
256 occupy fractured zones in the aureole, around the intrusion tip, or within the sills themselves
257 (Figure 3). In a previous summary of fracture types present in the sills of the study area,
258 Rabbel et al. (2021) interpreted these features as hydraulic fractures. Thus, evidence from
259 graphitized bitumen in the fractures in the aureole and in the sills, themselves points to
260 hydrocarbon transport in a high-pressure, high-temperature environment in which at least part
261 of the mobilized carbon transforms to graphitic carbonaceous material.

262 Our field and sample results strongly suggest that significant volumes of hydrocarbons
263 circulated through the sills when the temperature at their margins were 350-500°C. The
264 temperature was thus likely much higher in the interior of the sills. We infer that
265 hydrothermal flow also occurred along cooling fractures within the sill, i.e., the sill developed
266 some permeability while still hot (cf. Figure 2). This interpretation challenges the common
267 assumption of previous models of hydrothermal circulation around cooling sills that
268 intrusions remain impermeable during cooling (Aarnes et al., 2012; Iyer et al., 2017, Galerne
269 & Hasenclever, 2019). An exception is the study of Iyer et al. (2013), who tested a model that
270 included a linear permeability increase inside a cooling sill. While the author highlighted the
271 importance of this mechanism in potentially facilitating the upward migration of thermogenic
272 gas generated beneath sills, this test remained exploratory. Hence, no great detail analyses on
273 the effect on the hydrothermal flow is provided nor supported by field evidence. Thus, a
274 dedicated modeling study on the effect of permeability creation in cooling sills on
275 hydrothermal flow and hydrocarbon transport is still missing. In the following section, we
276 therefore present numerical simulations to test the effects of permeability increase associated
277 with fracturing within the sills on hydrothermal circulations, and how this affects
278 hydrothermal transport of hydrocarbons around sills.

279

280



281

282 Figure 4. (a) Raman spectra from two bitumen vein samples at Sierra de la Cara Cura and El
 283 Manzano, respectively. Both samples include well-developed and narrow graphite and
 284 disordered carbon (G, D1) peaks as well as weak or absent D3 band. (b) Reference Raman
 285 spectrum to illustrate spectra decomposition of carbonaceous material after Beyssac et al.
 286 (2002).

287

288 **3 Numerical Simulations**

289 **3.1 Model description**

290 We employ the two-dimensional (2D) finite element model of Galerne and Hasenclever
291 (2019), who applied it to quantify degassing through sill-related hydrothermal vents. The
292 model is presented in detail in Galerne and Hasenclever (2019). Here we will limit our
293 discussion to the main features and point out the key adjustments made for this study. In
294 short, the model simulates hydrothermal flow around a cooling sill, but is also coupled to a
295 model for the heat-driven chemical transformation of organic matter into hydrocarbons
296 (represented by methane). This allows us to investigate not only hydrothermal circulation
297 around sills as such, but also how this affects transport of the hydrocarbons generated in the
298 contact aureole.

299 The model considers single-phase hydrothermal flow of a compressible fluid in a porous
300 medium following Darcy's law. Temperature calculations comprise heat diffusion, heat
301 advection and heat sources/sinks related to latent heat of magma crystallization, mineral
302 dehydration, thermal cracking of organic matter as well as internal fluid friction and pressure-
303 volume work. Fluid density varies with temperature and pressure according to the equation of
304 state of pure water. The pore pressure equation also contains source terms representing fluid
305 release due to temperature-dependent, irreversible contact metamorphic reactions, including
306 (i) organic matter transformation into methane and (ii) clay mineral dehydration.

307 To calculate organic matter transformation to light hydrocarbon, the model uses the
308 EASY%Ro method (Sweeney and Burnham, 1990), which quantifies the converted fraction
309 of organic matter and thermal maturity through vitrinite reflectance Ro . Here we assume that
310 all organic matter transformation is converted to methane. We monitor transport and
311 accumulation of the released methane due to hydrothermal flow using a finite volume
312 advection scheme, but do not consider buoyancy effects resulting from the addition of
313 methane to the pore fluid. Clay mineral dehydration follows the maximum storable weight
314 fraction of water in the stable mineral assemblage at a given temperature, which is predicted
315 by phase equilibria (Connolly, 2009). This process not only produces additional pore fluid,
316 but also causes a permanent porosity increase in the affected host rock to ensure mass
317 conservation of rock and fluid. The brittle-ductile transition for the host rock is assumed to
318 happen at 500-750°C and linearly decreases permeability (Galerie and Hasenclever, 2019).
319 Note here that while the model calculations are conducted with methane properties, the
320 results can be used to understand transport of (light) hydrocarbons around sills in general. We
321 thus use “methane” and “hydrocarbons” interchangeably in the context of this study.

322 The model provides time series of the 2D fields of all relevant rock properties in the model
323 domain and physical quantities related to metamorphism and hydrothermal flow. Since we
324 investigate the impact of permeable sills on the fluid and hydrocarbon circulation, our
325 analysis focuses on visualization of the temperature, permeability, fluid flow and methane
326 accumulation during cooling of the sill.

327 **3.2 Adjustments for this study**

328 We adjusted three aspects of the original model to honor geological observations in the study
329 area. First, we limit rock failure to tensile hydraulic fracturing and do not consider shear
330 failure. We assume that hydrofracturing occurs at sufficiently high pore fluid overpressure,
331 i.e., if pore pressure exceeds the sum of lithostatic stress and tensile strength. Here, we
332 consider hydrofracturing of the host rock and not within the sill.

333 Second, we assume that hydrofracturing increases not only permeability but also porosity,
334 which is often neglected in numerical models for simplicity. However, the additional space
335 provided by the opening of hydraulic fractures, which we approximate by the porosity
336 increase, is an important storage buffer during thermal expansion of fluids and hydrocarbon
337 generation. While a transient permeability increase during hydrofracturing can easily be
338 defined without affecting the numerical stability and physical plausibility of the model,
339 prescribing a porosity increase associated with hydrofracturing is not straightforward. A
340 prescribed too large porosity increase, for instance, would create a strong suction effect
341 leading to unrealistically low pressures or even underpressure. We solve this problem by
342 iteratively increasing porosity in regions where overpressure exceeds the failure criterion and
343 solving again for the pore pressure field until a consistent solution establishes, which on
344 average requires 10 – 15 iterations. Hydrofracturing is treated reversible and its effects on
345 porosity and permeability vanish once pore pressure drops below the failure criterion. In this
346 study, we limit the hydrofracturing-related maximum porosity and permeability increase to
347 1% and a factor 100, respectively.

348 Third, we approximate the process of cooling joint formation within the sill through a linear
349 temperature-dependent permanent increase of permeability, similar to Iyer et al. (2013), but
350 additionally consider the corresponding permanent porosity increase. Note that this is likely a
351 strong simplification of fracture flow through cooling joints networks, which is still poorly
352 constrained. Bulk volume reduction of the cooling and crystallizing magma induces thermal
353 stresses that lead to the formation of a cooling joint network, creating primary porosity and
354 permeability inside the intrusion (e.g., Petford, 2003; Hetényi et al., 2012). The overall pore-
355 space gained in our model is set to equate 8% volume loss, occurring during the transition
356 from a melted to crystallized magma (between the liquidus and solidus temperature), based
357 on reported fracture porosities from fractured sill reservoirs in the study area (Witte et al.,
358 2012; Spacapan et al., 2020). To be consistent with the crystal-mush model described by
359 Marsh (2002), the onset of the pore opening should be when 50-55% of the magma has
360 crystallized. Here we take a value of 1000°C as the onset of the brittle-ductile-transition
361 (BDT) temperature. Using a linearized, temperature-dependent definition of the melt fraction,
362 $(T-T_S)/(T_L-T_S)$, with $T_L = 1100^\circ\text{C}$ and $T_S = 900^\circ\text{C}$ being liquidus and solidus temperatures,
363 respectively, the set value for the BDT in our simulations implies that cooling joint creation
364 starts at 1000°C when, at any distance from the sill margins, 50% of the sill has crystallized
365 (Figure 5c).

366 Note that we limit the model's representation of fracturing in the sill to cooling joints and thus
367 perform a strong simplification compared to the complex interplay of thermal and hydraulic
368 fracturing mechanisms observed in the field (cf. section 3 and Rabbel et al, 2021). However,
369 the goal is to study the general impact of porous and permeable sills on hydrothermal flow and
370 the associated hydrocarbon transport and storage, which is possible even with this limitation.

371

372 **3.3 Modeling setup**

373 Our modeling setup consists of a single flat sill of 50 m thickness and 1 km length emplaced
374 at 3 km depth in a homogenous host rock (Figure 5a). We performed a parameter sensitivity
375 study to investigate the impact of porosity and permeability development in sills emplaced in
376 either low-permeability (e.g., shale) and high-to-medium permeability (e.g., silt-/sandstone)
377 host rocks. This allows us to investigate the sills of our study area (sills in Agrío and Vaca
378 Muerta shales), but also compare to sills in the overlying Neuquén Group (silt-/ sandstones)
379 or other relevant geological settings. Figure 5 and Table 1 show the model setup and the list

380 of important model parameters, respectively. Note that we also carried out simulations with
381 thinner sills and provide those as supplementary data for completeness, but do not address
382 their results in detail. We first conducted a series of reference setups including permanently
383 impermeable sills emplaced in either high- or low-permeable host rock. Subsequently, we
384 use the same setups but activate temperature-dependent porosity and permeability generation
385 for the sills. Hydraulic fracturing of the host rock is activated in all simulations.

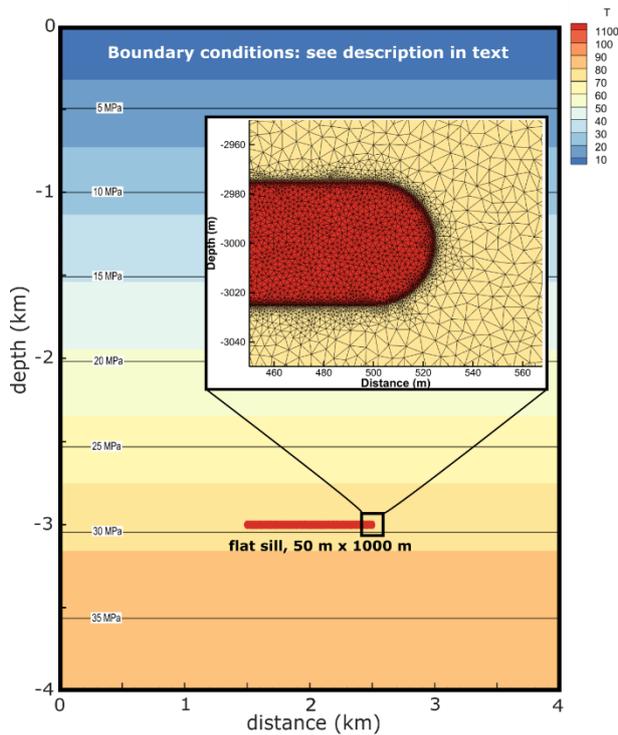
386 The model domain (Figure 5a) is 4000 m wide, extends from the surface to 1000 m below the
387 emplacement depth of the sill and is discretized with a triangular mesh with variable element
388 sizes between 0.5 - 50 m (smallest around and within the sill, see Figure 5c). We assume
389 instantaneous sill emplacement at 1100°C, corresponding to the inferred liquidus
390 temperatures of andesitic magma in the study area (Spacapan et al., 2018). The left and right
391 boundaries are insulating and impermeable, and we calibrated the fixed temperature at the
392 impermeable bottom boundary to create a geothermal gradient of around 25 °C/km. The top
393 boundary mimics the behaviour of a shallow seafloor with temperature set to 10 °C, pressure
394 set to 0.1 MPa and free in- and outflow. Otherwise, initial conditions consider no basin
395 history such as uplift and erosion, and no pre-existing thermal maturation prior to sill
396 emplacement. We justify this by low background maturity values reported in the study area
397 (Spacapan et al., 2018; Palma et al., 2019; Rabbel et al., 2021).

398 For the sediments, we chose a homogenous material with 5% TOC and ca. 5 weight percent
399 bound water, as well as exponential decay of porosity with depth (Figure 5b). Permeability is
400 porosity-dependent and follows a Kozeny-Carman relationship. This relationship is calibrated
401 to values of low-permeability Agrio and Vaca Muerta shale formations at 2-3 km depth in the
402 northern Neuquén Basin, yielding 10^{-18} m² at 3 km. This corresponds to emplacement depths
403 for the igneous petroleum systems present in both subsurface and outcrop (Figure 5b). To
404 compare with settings with more permeable lithologies (silt-/sandstone), we increased the
405 host rock permeability by two orders of magnitude for another set of simulation (Figure 5b).
406 Yet, each host rock setup remains a simplification, as we do not include lithological
407 variations. Sill permeability starts at 10^{-20} m² (impermeable) and increases to 10^{-15} m² at
408 900°C. In lack of macroscopic permeability measurements, we chose the maximum
409 permeability value to approximate the upper range for Neuquén sill reservoirs as reported by
410 Spacapan et al. (2020). These values were obtained from (micro-)fractured sill matrix
411 samples and therefore likely underestimate bulk permeability, as they do not include
412 macroscopic cooling joints.

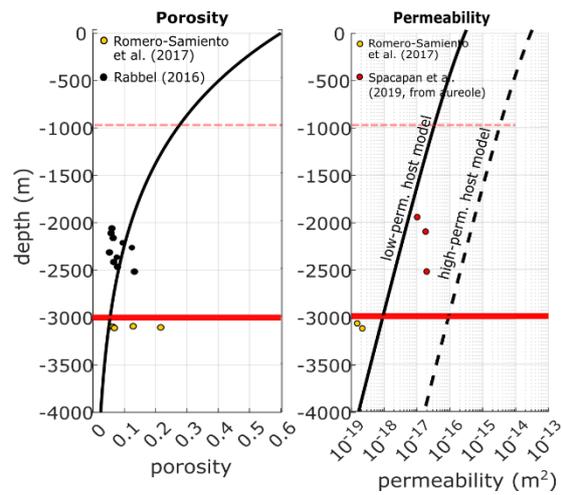
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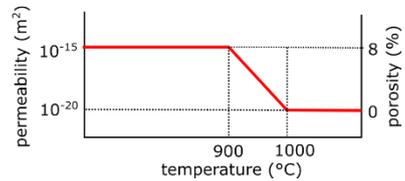
a) Simulation Setup Sketch



b) Porosity-Permeability models for host rock



c) Temperature-permeability/porosity evolution in sill



415

416 Figure 5. Illustration of the numerical modelling setup. (a) Model domain, with initial
 417 temperature and pore pressure field and close-up for mesh illustration. (b) Porosity-
 418 permeability-depth relationships in the models alongside reported data from the Vaca Muerta
 419 formation from various depths (Rabbel 2017, Romero-Samiento et al. 2017, Spacapan et al.
 420 2019). Note that permeabilities from Spacapan et al. (2019) are the smallest values reported
 421 from the fractured and altered aureole of sills and thus likely overestimating background
 422 values. (c) Illustration of porosity-permeability-temperature function for sill with assumed
 423 cooling joint formation.

424

425

426 Table 1. Material properties used for hydrothermal simulations

Fixed rock parameters	Value	Unit
Fluid Properties	Equation of state for pure water (IAPS-84)	
Magma / Sill Properties (representing andesite)		
Specific heat capacity ^{1,2}	900	J kg ⁻¹ K ⁻¹
Latent heat of crystallization ^{1,2}	320	kJ kg ⁻¹
Initial sill temperature ^{3,4}	1100	°C
Liquidus temperature ³	1100	°C
Solidus temperature ⁴	900	°C
Brittle-Ductile-Transition (permeability creation)	1000-900	°C
Sill thickness ⁴	50 (+ 10 in supplement)	m
Density ^{2,6}	2830	kg m ⁻³
Thermal conductivity ⁵	2.51	W m ⁻¹ K ⁻¹
Permeability range in sill ⁷	10 ⁻²⁰ – 10 ⁻¹⁵	m ²
Porosity range in sill ⁷	0 – 0.08	1
Host Rock Properties (representing shale)		
Density ⁶	2600	kg m ⁻³
Specific heat capacity ⁵	960	J kg ⁻¹ K ⁻¹
Thermal conductivity ⁵	2.55	W m ⁻¹ K ⁻¹
Porosity ⁶	Ref fig	Ref fig
Permeability	Ref fig	Ref fig
Initial TOC ⁴	0.05	1
Initial bound water content ²	0.048	1
Tensile strength ⁸	3	MPa
Enthalpy of mineral dehydration ^{1,2}	2800	kJ kg(H ₂ O) ⁻¹
Enthalpy of organic cracking ^{1,2}	375	kJ kg(TOC) ⁻¹
Brittle-ductile-transition ²	500-750	°C
Sources: ¹ Aarnes et al. (2010), ² Galerie and Hasenclever (2019), ³ Stern et al. (1975), ⁴ Spacapan et al. (2018), ⁵ Angenheister et al. (1982), ⁶ Rabbel (2017), ⁷ Spacapan et al. (2020), ⁸ Schön (2015)		

427

428

429 3.4 Numerical Simulation: results

430 3.4.1 Impermeable sill: Flow around sill tip and methane plume

431 We first present and compare the results of the reference simulations of a 50 m thick,
432 impermeable sill emplaced at 3 km depth in either a low-permeability (Figure 6) or high
433 permeability (Figure 7) host rock. Figure 6 and 7 display the evolution of temperature,
434 vitrinite reflectance (R_o) as proxy for thermal maturity (first row of images), pore fluid
435 pressure (second row), permeability and transiently opened fracture porosity to highlight
436 regions of active hydrofracturing (third row) and the methane fraction of the fluid (fourth
437 row) in the model at 1, 60 and 1000 years after the sill is emplaced. Additionally, flow
438 vectors colored by fluid velocity in the pore space are shown. We highly recommend to also
439 view the movies supplied in the additional materials of this paper to get a better sense of the
440 process dynamics in space and time.

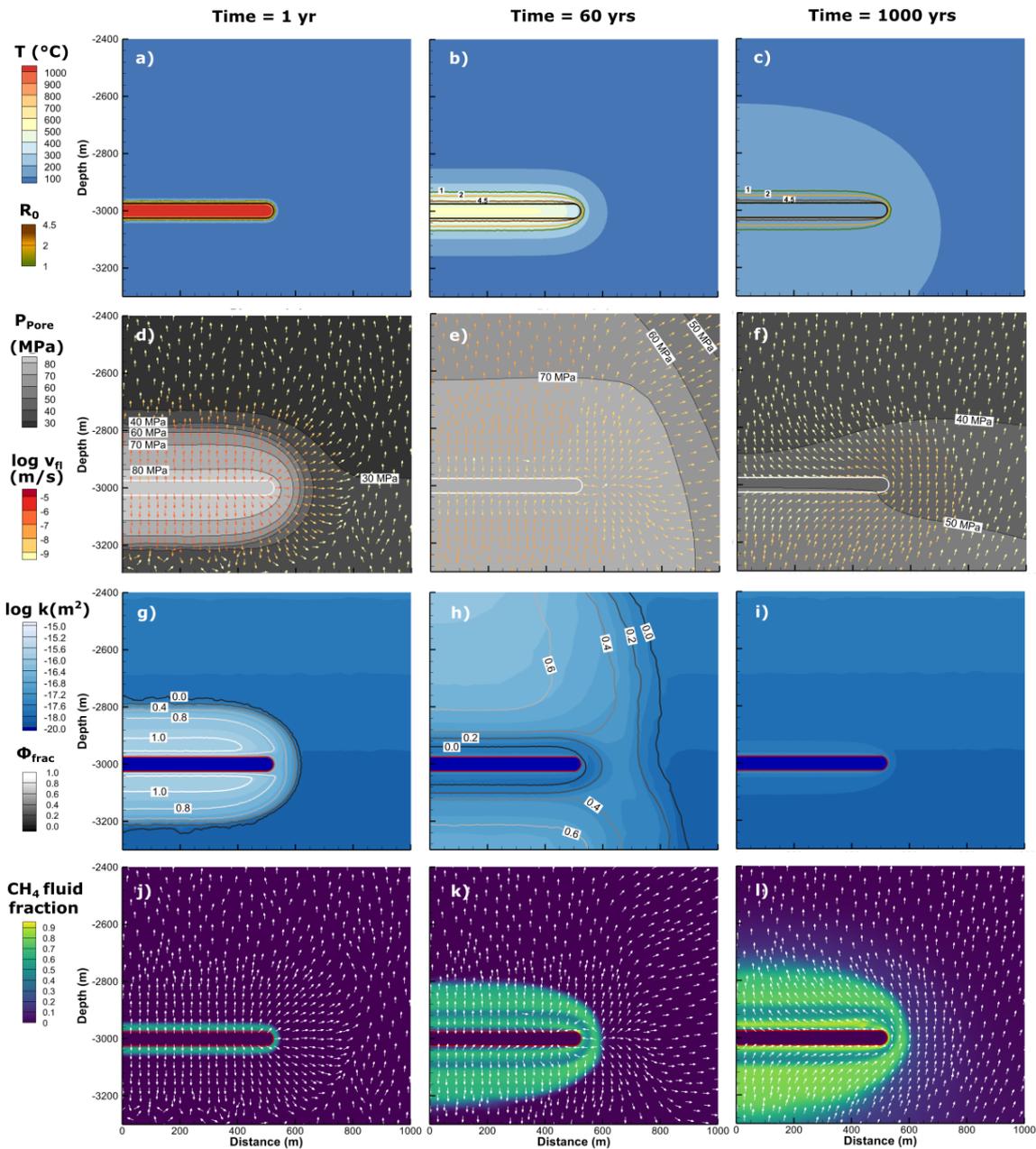
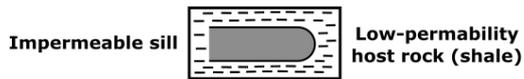
441 One year after emplacement (left column in Figure 6 and 7), the sill is still over 1000°C hot
442 and only the host rock within <10 m distance to the sill has been heated to temperatures of
443 >350°C. Within the thin thermal aureole, thermal maturity increases strongly, methane is
444 generated by thermal cracking and mineral dehydration takes place. Thermal expansion of the
445 heated fluid, mineral dehydration and methane generation close to the sill lead to strongly
446 elevated pore fluid pressures, which propagate away from the sill (Figure 6d and 7d). In both
447 low- and high-permeability scenarios we observe similar peak pore pressures of around 80 –
448 85 MPa at the sill contact during the first weeks after the sill emplacement. However, in the
449 high-permeability case (Figure 7d), the pressure front moves faster because of efficient fluid
450 flow even far away from the sill. In contrast, the pressure front moves much slower in the
451 low-permeability case (Figure 6d), where fast fluid flow is restricted to the region of active
452 hydrofracturing (Figure 6g vs. 7g). Flow direction and therefore methane transport is sill-
453 parallel towards the tips in the highly permeable contact aureole and radially outwards
454 outside of the aureole (Figure 6d,j and 7d,j).

455 After 60 years of sill cooling (central column in Figure 6 and 7), the high-temperature aureole
456 (>350°C) has expanded to ca. 25 m around the sill (Figure 6b and 7b), thermal maturity has
457 reached R_o values above 2, which indicates the gas window or overmaturity. At this stage,
458 pore pressure and permeability distribution in the two reference cases differ markedly. In the
459 low-permeability host, fluid overpressure is still sufficiently high to cause hydrofracturing in
460 a several 100 m wide halo around the sill, where permeability is elevated by 1-2 orders of
461 magnitude with respect to background values (Figure 6e, h). Note that hydrofracturing is
462 vanishing close to the sill where hydrofractures are closing again as pore pressure slowly
463 reduces. The high-permeability host allows for a more efficient dissipation of fluid
464 overpressures so that no more hydrofracturing occurs ca. 10 years after the sill emplacement.
465 After 60 years, the remaining fluid overpressure around the sill is only a few MPa above
466 hydrostatic (Figure 7e, h). Both models also develop porosity and permeability increase due
467 to clay mineral dehydration, but this is limited to 50 m distance from the sill contact. In both
468 models, the highest temperature in the inner aureole (ca. 10 m from contact) is reached ca. 10
469 years after the sill emplacement (around 670 °C at the sill contact), while the outer aureole
470 (up to ca. 50 m to contact) reaches its peak temperatures (300-400°C, depending on distance
471 to sill) after around 60 years. The combined action of thermal contraction of the cooling fluid
472 after reaching the peak-temperature and additional closure of hydrofractures (i.e., pore-space
473 reduction) in the low-permeability host cause an inversion of the flow direction. After 60
474 years, fluids carrying high methane concentrations migrate towards the sill within a ~100 m
475 thick region above and below the sill (Figure 6e,k, Figure 7e,k and supplementary movies).
476 Despite the differences in permeability structure and pressure regimes, flow patterns of both

477 reference simulations are relatively similar. The contact-parallel flow in the high permeability
478 host is stronger, and these higher flow velocities lead to a more pronounced plume of rising
479 methane on top of the sill near its tip (Figure 7k). Both cases also show a sizable methane
480 accumulation remaining below the sill.

481 The right column in Figure 6 and 7 represents the end of the simulation after ca. 1000 years.
482 The temperatures throughout the model are still elevated with respect to the initial geotherm
483 but are now below 200°C everywhere. Fluid pressure in the low-permeability case is still up
484 to 20 MPa above hydrostatic but has dropped below the failure criterion and hydrofracturing
485 has stopped (Figure 6f,i). In the high-permeability host, pore pressure is reduced to values of
486 <1 MPa above hydrostatic (Figure 7f). The dehydration-related permeability increase has not
487 expanded significantly in either model. In the low-permeability scenario, some of the
488 methane rises to 250 m above the sill, but the highest concentrations (almost pure methane,
489 i.e., mass fraction close to 1) occur within 50–100 m to the sill contact (Figure 6l). In
490 contrast, the model with a high-permeability host shows the formation of a localized
491 secondary plume of very high methane concentrations (essentially pure methane) rising
492 above the sill (Figure 7l), and the initial methane plume has reached ca. 400 m above the sill.
493 The aureole below the sill has also accumulated high methane concentration in the fluids with
494 up >70% methane fraction within 30 m of the sill. These methane-rich fluids remain trapped
495 below the impermeable sill.

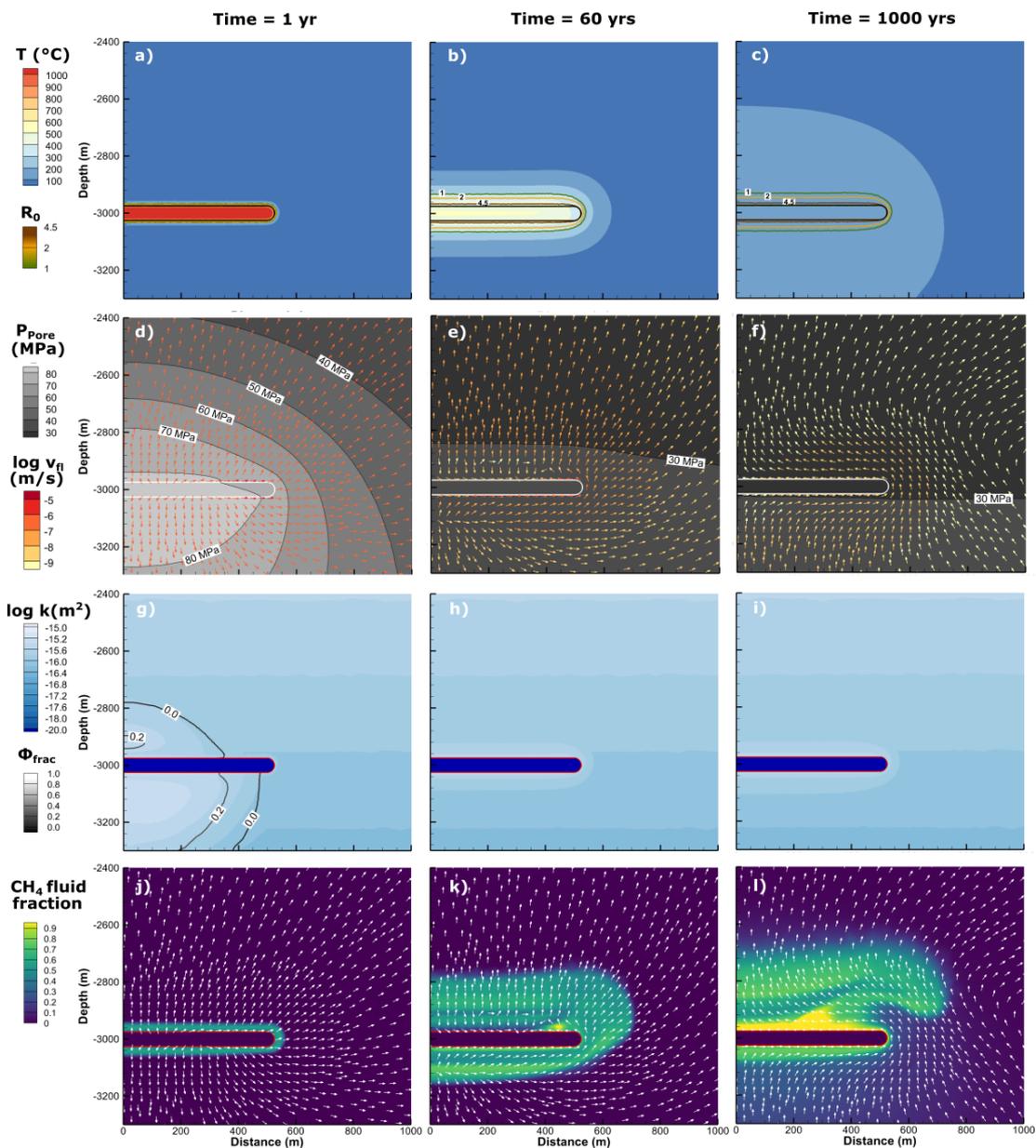
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497

498 Figure 6. Reference simulation results for an impermeable sill of 50 m thickness at 3 km
 499 depth in the low-permeability host case. The columns correspond to 1, 60 and 1000 years of
 500 simulated time after emplacement, respectively. The rows represent four parameters
 501 characterizing thermal state, contact metamorphism and hydrothermal transport of methane:
 502 (a-c) Temperature, thermal maturity as vitrinite reflectance R_0 contours, (d-f) pore fluid
 503 pressure with flow vectors coloured by pore velocities, (g-i) permeability with fracture
 504 porosity contours, (j-l) methane fraction in fluid.

505



506

507

508 Figure 7. Reference simulation results for an impermeable sill of 50 m thickness at 3 km
 509 depth in the high-permeability host case. The columns correspond to 1, 60 and 1000 years of
 510 simulated time after emplacement, respectively. The rows represent four parameters
 511 characterizing thermal state, contact metamorphism and hydrothermal transport of methane:
 512 (a-c) Temperature, thermal maturity as vitrinite reflectance R_0 contours, (d-f) pore fluid
 513 pressure with flow vectors coloured by pore velocities, (g-i) permeability with fracture
 514 porosity contours, (j-l) methane fraction in fluid.

515

516

517 3.4.2 Permeable sill: Opened upward flow path and flow reversal

518 The introduction of cooling-related permeability generation in the sill profoundly changes the
519 development of hydrothermal flow and methane transport patterns for both host rock types.
520 We show identical parameters and time steps as for the reference simulations for low- and
521 high-permeability host in Figure 8 and 9, respectively. To describe the details of the evolving
522 flow patterns and hydrocarbon transport in the first 60 years, we add close-up figures for both
523 cases displaying fluid pressure with flow vectors and permeability plots with (transient)
524 fracture porosity (Figure 10). Finally, we quantify the total generated methane mass, the
525 methane mass exposed to temperatures $>400^{\circ}\text{C}$ (graphitization conditions), the accumulation
526 of methane in the permeable and porous sill, and the average sill temperature over time
527 (Figure 11). Again, it is instructive to also view the supplementary movies for the respective
528 simulations.

529 One year after emplacement, initial cooling of the sill leads to a progressing
530 porosity/permeability front where the temperatures approach the solidus defined as 900°C
531 (Figure 8 and 9, left column). At this stage, which continues until the sill becomes fully
532 permeable, the simulations closely resemble the reference runs (see also supplementary
533 animations of the simulations). Temperatures in the sediments are elevated only close to the
534 intrusion, where the sediments almost instantly produce gas or become thermally overmature,
535 i.e., vitrinite reflectance is larger than 1.5 (Figure 8a, 9a). Again, pore pressures are strongly
536 elevated, and porosity and permeability in the sediments increase due to hydrofracturing
537 around the sill (Figure 8d, g and 9d, g). Although large-scale methane distribution appears
538 nearly identical to the reference scenarios, the detailed view shows that the onset of porosity
539 and permeability generation in the sill's outermost regions allows flow within the sill (Figure
540 10, first and second column). Nevertheless, most fluid flow and methane transport is directed
541 away from the sill contact with some sideways flow towards the sill tip along the bottom
542 contact.

543 The entire sill has become fully porous and permeable after about 12 years (just after Figure
544 10, second column). While the same processes as described for the reference cases take place
545 here as well (vanishing hydrofracturing, cooling of the contact aureole, fluid contraction and
546 beginning inversion of the flow field towards the sill), the opening of the sill leads to stronger
547 and more focused flow towards it. The flow direction below the sill changes from downward
548 or contact-parallel to near-vertical upwards (Figure 8e, Figure 9e), "flushing" high methane
549 concentrations directly into the sill (Figure 8k, 9k, 10c, i). However, at this point there are
550 differences between the simulations considering a low- vs. high-permeability host rock.

551 In the low permeability case, a convection cell evolves within the sill, whose permeability is
552 2 – 3 orders of magnitude higher than that of the surrounding shale (Figure 10c). Methane
553 transport into the sill is further enhanced by the much slower overpressure dissipation
554 through this host rock and the fading hydrofracturing. After 60 years, hydrofracturing and the
555 associated fracture porosity are progressively reduced and eventually stopped near the sill
556 (Figure 8h, 10f), where fluid pressures have dropped below the failure criterium. The
557 pressure drop is primarily caused by thermal contraction of the cooling fluids in the aureole
558 and within the cooling sill. While the front of fading hydrofracturing and associated fracture
559 porosity propagates away from the sill, methane-rich fluid is "squeezed" out of the host rock
560 (Figure 8l, 10c, f) and contributes to a sustained flow and methane transport towards the sill.
561 In this way, the sill is charged with up to 11,000 tons of methane (7.5% of total generated
562 methane) within the first 100 years (Figure 11a). Methane mass in the cooling sill rises to ca.
563 16,000 tons (8.8% of total) after 250 years.

564 In the high permeability host rock, the porosity and permeability structure are much simpler,
565 because hydrofracturing is absent due to generally lower fluid pressure (Figure 9e). The
566 permeable sill and its dehydrated aureole, with slightly higher permeability than the
567 background (Figure 9h, k, 10l), now represent an upward pathway and storage layer for
568 methane-rich fluids. In addition to upward flow from below the sill, the flow directions at the
569 intrusion tip also change and a circular flow pattern develops centered around a vortex
570 located at the top of the intrusion tip (Figure 9e, 10i). This vortex initiates a sideward and
571 downward directed flow that transports some of the methane from above and next to the sill
572 tip back towards and into the sill tip (Figure 9k). During this phase of “methane flushing”,
573 more methane enters the sill from below and through the tip than is lost through the top
574 contact, and thus methane stored in the intrusion rises to 10,000 tons (6.3% of total) until ca.
575 90 years of simulation time (Figure 11b).

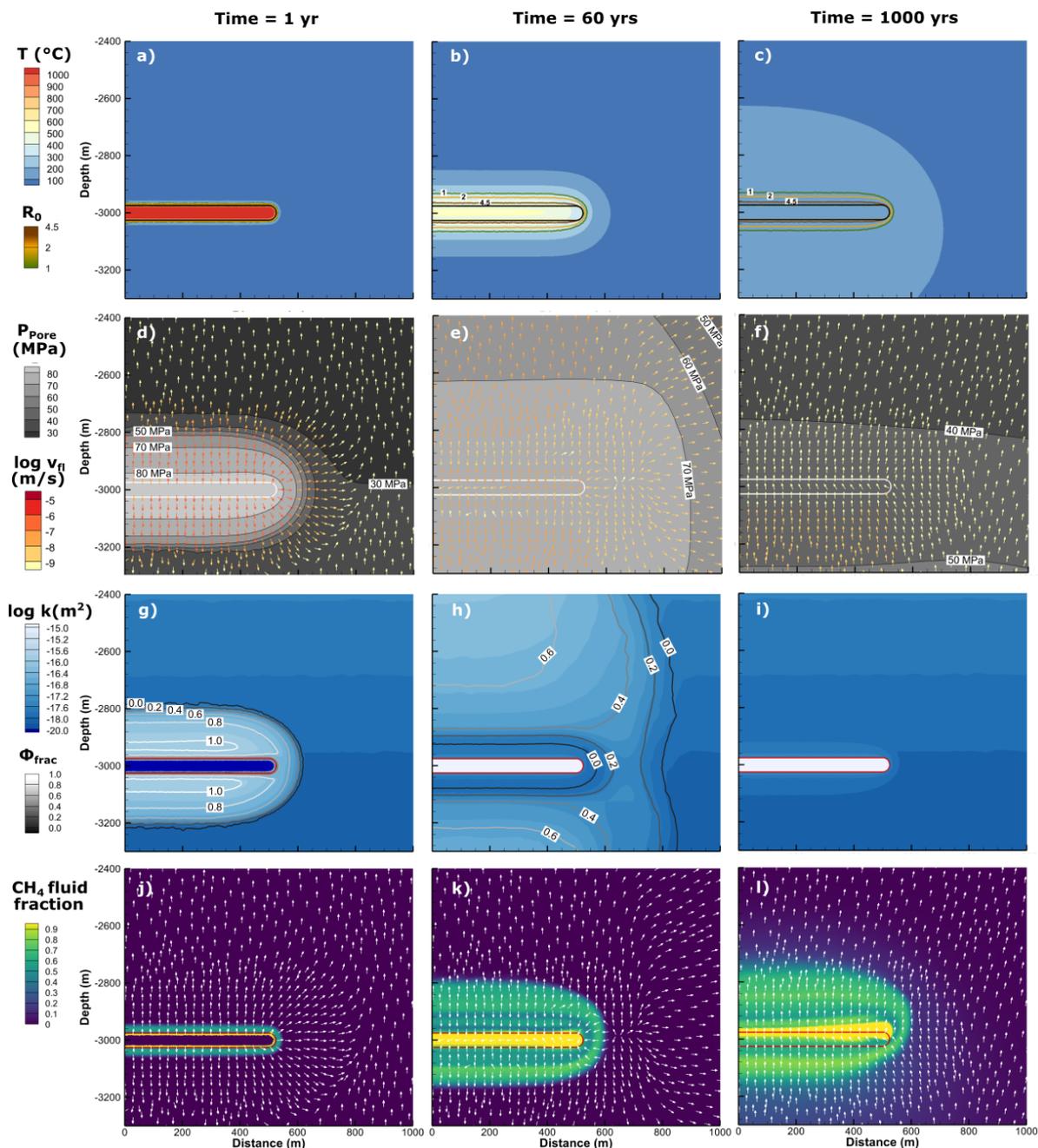
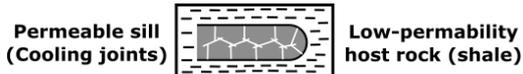
576 Interestingly, in both cases the average temperature in the sill during this stage of “flushing”
577 is still between 400-800 °C (Figure 11). Thus, up to 7.5% of the overall generated methane
578 that reaches the sill and the innermost 10 – 20 m of the aureole in the first ca. 100 years is
579 exposed to these temperatures.

580 In the following phase until the end of both simulations, the sill and sediments cool down to
581 below 200°C, thermal maturity increases only marginally, and fluid pressures dissipate to a
582 level below the hydrofracturing point (right columns in Figures 8 and 9). In the simulation
583 comprising a low-permeability host rock, highest methane concentrations accumulate within
584 the sill or within 50 m of the upper contact (Figure 8l). In the high-permeability case, the
585 release of methane previously trapped under the sill creates a slowly rising band of very high
586 methane concentrations above the sill (Figure 9l). Flow velocities have further reduced by
587 another 2 orders of magnitude, indicating that hydrothermal flow is stalling. In the last few
588 hundred years of each simulation, the methane amount in the sill reduces to about 10,000 tons
589 (5.5% of total) and 4000 tons (2.5% of total), respectively (Figure 11).

590 In summary, we identify three hydrothermal flow phases in the case of a sill with porosity
591 and permeability evolution due to cooling joint formation. These include
592 (i) fluid flow and methane transport away from the sill and contact-parallel as long as the sill
593 core is impermeable,
594 (ii) partial backflow and “methane-flushing” into, or through the sill once it is completely
595 fractured; in low-permeability environments, this phase is accompanied by the closure of
596 hydrofractures around the sill, thereby “squeezing out” methane-rich fluids that enter the sill
597 (iii) stabilization of hydrocarbon in and around the sill during fading hydrothermal
598 circulation.

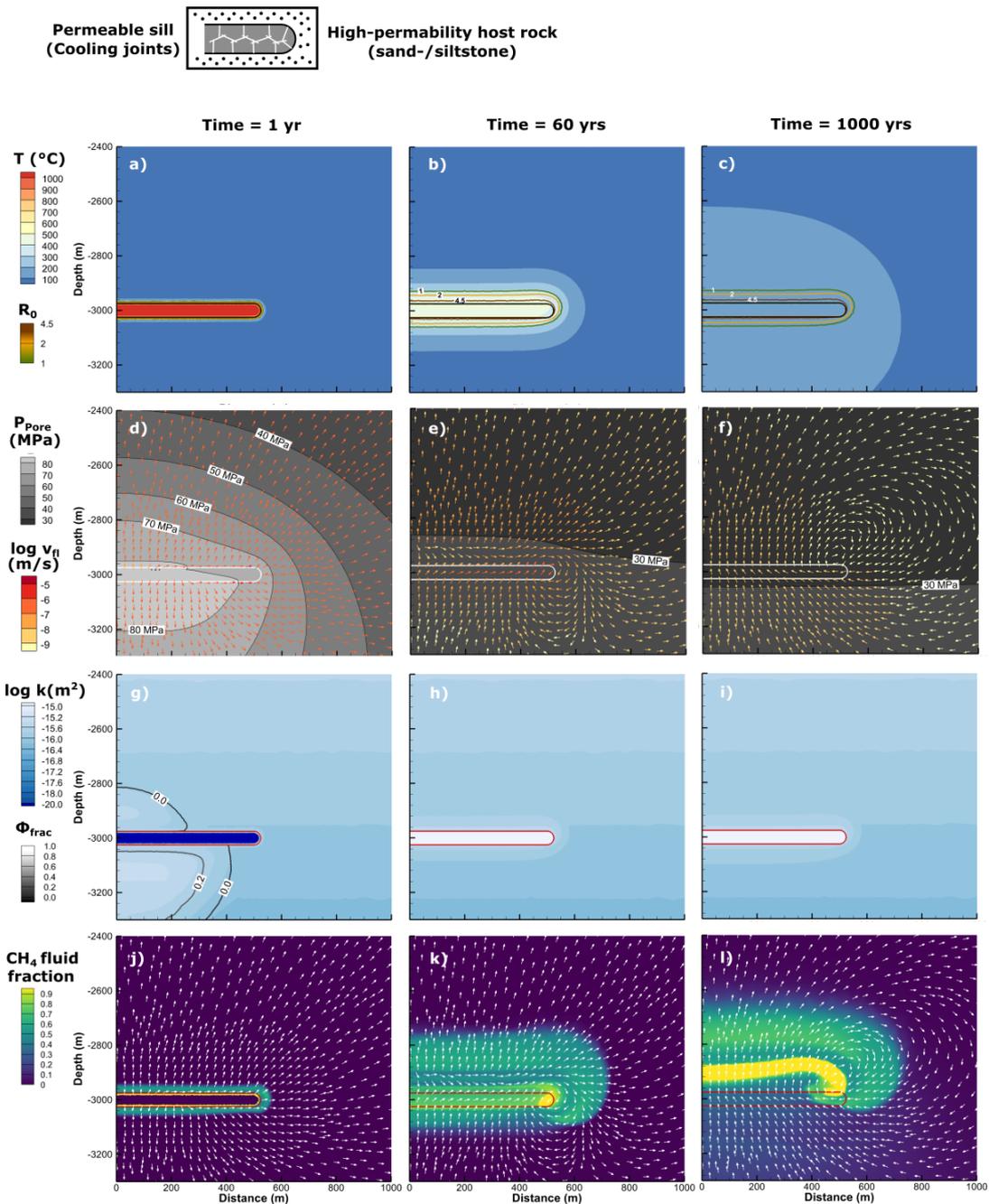
599 Despite the differences in the physical behavior between low- and high-permeability host
600 rocks in phase (ii), i.e. fracture-facilitated fluid flow vs. matrix flow, substantial amounts of
601 methane-rich fluids from below enter the porous and permeable sill shortly after its
602 solidification at 900 °C. Compared to the simulations with impermeable sills, fluid
603 overpressures dissipate slightly faster in the same host rock type, because the opening of pore
604 space in the sill compensates for a small fraction of the overpressure.

605



606

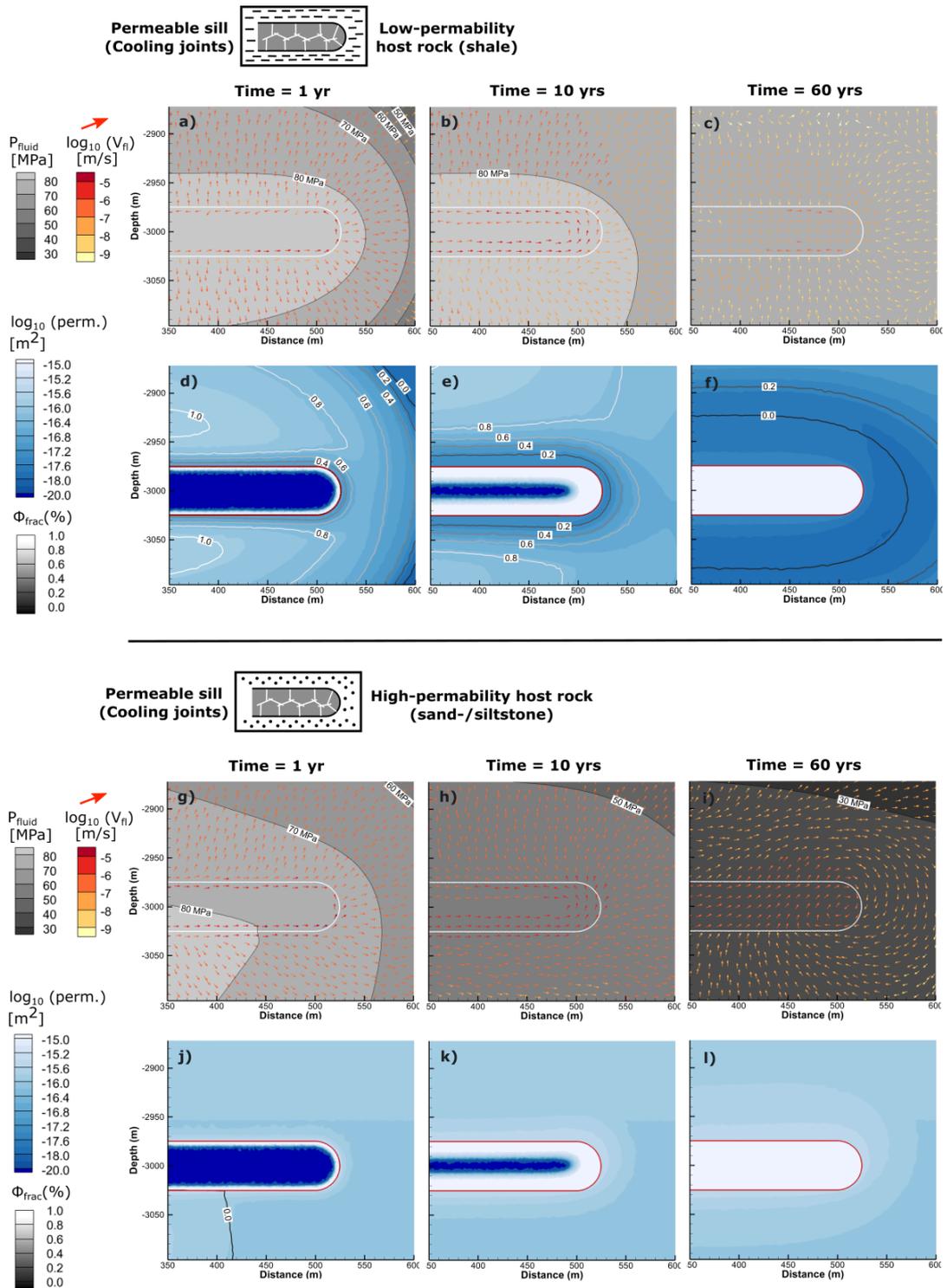
607 Figure 8. Simulation results for identical conditions as the reference case in Figure 6 (50 m
 608 thick, 3 km depth, low-permeability host), but the sill develops porosity and permeability
 609 with cooling. The rows represent four parameters characterizing thermal state, contact
 610 metamorphism and hydrothermal transport of methane: (a-c) Temperature, thermal maturity
 611 as vitrinite reflectance R_0 contours, (d-f) pore fluid pressure with flow vectors coloured by
 612 pore velocities, (g-i) permeability with fracture porosity contours, (j-l) methane fraction in
 613 fluid.



615

616 Figure 9. Simulation results for identical conditions as the reference case in Figure 7 (50 m
 617 thick, 3 km depth, high-permeability host), but the sill develops porosity and permeability
 618 with cooling. The columns correspond to 1, 60 and 1000 years of simulated time after
 619 emplacement, respectively. The rows represent four parameters characterizing thermal state,
 620 contact metamorphism and hydrothermal transport of methane: (a-c) Temperature, thermal
 621 maturity as vitrinite reflectance R_0 contours, (d-f) pore fluid pressure with flow vectors
 622 coloured by pore velocities, (g-i) permeability with fracture porosity contours, (j-l) methane
 623 fraction in fluid.

624



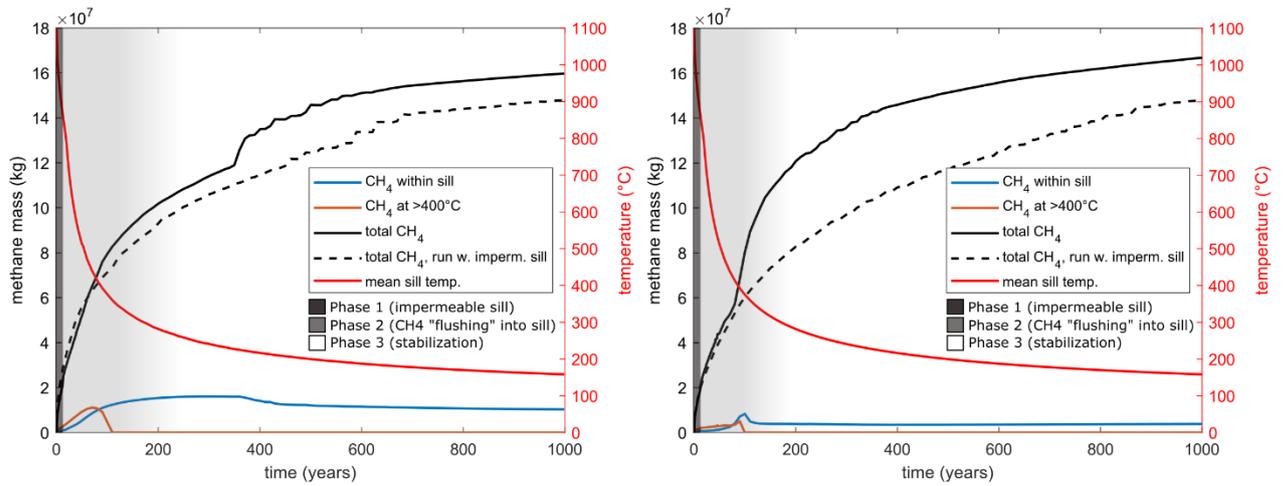
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626 Figure 10. Close-up view of fluid pressure with flow vectors and permeability with fracture
 627 porosity contours at 1, 10 and 60 years of the permeable sill simulations in low-permeability
 628 (a-f) and high-permeability (g-l) host rocks.

629

a) Low-permeability host rock

b) High-permeability host rock



630

631 Figure 11. Cumulative methane mass for the simulations with permeable sills in low-
 632 permeability host (a) and high-permeability host (b) grouped by different criteria: within rock
 633 at >400 °C, i.e., sufficient for graphitization (orange line), within sill (blue line), total mass
 634 generated in the model (black line). Dotted black line gives the total generated methane mass
 635 for the respective reference simulation with an impermeable sill. Red line represents average
 636 sill temperature.

637

638 **4 Interpretation and discussion**

639 **4.1 Hydrocarbon transport through hot sills**

640 Outcrop data strongly suggests that transport of hydrocarbons generated around the sills in
641 the northern Neuquén Basin occurs both vertically through the igneous intrusions as well as
642 around their tip. Importantly, geochemical data from the study area suggests that the sills are
643 responsible for most, if not all, organic matter transformation in the host rock, because
644 background maturity between the sills is essentially zero (Spacapan et al., 2018; Rabbel et al.,
645 2021). We thus interpret the observed hydrocarbons in the field to result from magmatic
646 heating and hydrothermal activity rather than burial-related maturation.

647 Cooling joints and veins or dykes filled with black shale and bitumen are pervasive
648 throughout the fracture networks of sills in our study area (Figure 2 and 3). The example
649 shown in Figure 2 is particularly clear in showing the relationship between upwelling
650 fluidized structures and the fill of the cooling joints. Note that similar observations have been
651 documented in other outcrops at larger scales (Rabbel et al., 2021). We therefore propose that
652 the flow of hot fluids and fluidized sediments through sills is a common occurrence in the
653 northern Neuquén Basin. We thus provide solid outcrop evidence that sills can become
654 preferred hydrocarbon transport pathways upon cooling. In contrast, this finding contradicts
655 the widespread assumption of impermeable intrusions in modeling studies of hydrothermal
656 fluid flow in volcanic basins (e.g., Iyer et al., 2013; Iyer et al., 2017; Galerne and
657 Hasenclever, 2019).

658 Our observations complement numerous growing evidence that (carbon-rich) fluids or
659 fluidized sediments entered sills of widely different sizes during their cooling stage, for
660 instance in the Faroe-Shetland Basin (Rateau et al., 2013; Schofield et al., 2020), Karoo
661 Basin (Svensen et al., 2010; Lenhard et al., 2023) and Guyamas Basin (Teske et al., 2021).
662 Lenhard et al. (2023) concluded that carbon-rich fluids must have entered the sill during
663 magma solidification, and Svensen et al. (2010) found metamorphosed sandstone dykes
664 within dolerite sills in the Karoo Basin with mineral assemblage indicative of temperatures
665 $>300^{\circ}\text{C}$. They proposed that strong pressure gradients between the overpressured host rocks
666 and the solidifying and contracting sill intrusions are likely responsible for liquefied
667 sediments entering the intrusions shortly after cooling. Similar to this study, Svensen et al.
668 (2010) based this interpretation on coupled thermo-hydraulic models, but explicitly
669 considered thermal contraction of the magma. Although our approach of adding porosity to
670 the sill is a simplification of this process, the general mechanism of hydrocarbon-rich fluids
671 and sediments entering the cooling joint network in the intrusions in the Neuquén Basin is
672 similar. We thus infer that our observations of hydrocarbon flow through hot sills in the
673 Neuquén Basin are widespread in volcanic basins worldwide.

674

675 **4.2 Impact of permeable sills on hydrothermal flow**

676 Our numerical simulations allow us to assess the effects of porosity and permeability increase
677 due to cooling joint formation when the sill has reached the solidus temperature. In
678 simulations with an impermeable (i.e., unfractured) sill, the intrusion acts as a constant
679 barrier for the hydrothermal flow and methane transport, while the aureole shows porosity
680 and permeability evolution due to hydrofracturing and dehydration (Figure 6 and 7). In this
681 configuration, an upward-rising methane plume initiates from the top contact of the sill. This
682 plume is much more pronounced and methane is transported further upwards if the host rock
683 is relatively permeable, such as silt- or sandstone (Figure 7). In addition, large amounts of

684 methane are trapped below the sill, since no vertical pathways are available through the sill
685 (Figure 6l and Figure 7l).

686 In contrast, a sill becoming permeable during cooling introduces drastic changes in the flow
687 and methane transport patterns, and results in three phases with very different characteristics.
688 The details of these phases differ depending on the type of host rock (low- vs. high-
689 permeability), but share many similarities as described below.

690 *Phase 1: Diverging and contact-parallel flow around the impermeable sill.*

691 Prior to complete solidification, the sill acts as a flow boundary and the flow patterns are
692 essentially the same as for impermeable sills. Hydraulic fractures in the aureole and around
693 the sill edge initiate in this early phase (Figure 10d, j), which may lead to the formation of the
694 large bitumen dykes and calcite veins observed in the field (Figure 3). In addition, porosity
695 generation inside the sill creates a suction effect that drives fluids into the intrusion (cf.
696 Svensen et al., 2010).

697 *Phase 2: Reversed flow and hydrothermal “flushing” of the solidified, permeable sill.*

698 The generation of cooling-related porosity and permeability inside the sill creates a hydraulic
699 connection between the lower and upper aureole and initiates the abrupt change to vertical
700 upflow into and through the sill (central column of Figure 8 and 9, Figure 10c, i). This rush of
701 hydrocarbons (here: methane-rich fluids) into and through the sill occurs at average sill
702 temperatures of 400–800 °C (Figure 11). Given the flow velocities and the timespan of
703 methane accumulation (Figure 8k, Figure 9k, Figure 11), the model shows that fluids flowing
704 through a 50 m thick sill are exposed to this high-temperature environment for tens of years,
705 which could be sufficient for graphite generation and therefore seems to fit well with outcrop
706 data and models for hydrothermal graphite generation (Figure 4; Rumble, 2014). We stress
707 that this is also valid for low-permeability host rocks such as shale, because sustained
708 hydrofracturing around the sill facilitates fluid flow by temporarily increasing porosity and
709 permeability in the host rock (Figure 8, 10d, e). “Squeezing” of hydrocarbons towards the
710 porous/permeable sill as the transiently opened fractures close corresponds well with the
711 observation of multi-generation bitumen dykes/dyke arrays in the study area (e.g., Figure 3a,
712 see also Rabbel et al. (2021)).

713 *Phase 3: Stabilization of hydrocarbon and fading hydrothermal flow.*

714 In low-permeability host rocks, reversed flow towards the sill continues for a few hundred
715 years and seems to be driven by closing of hydrofractures and thermal contraction of fluids
716 within the still cooling sill (Figure 8l, 10c, supplementary movies). In high-permeability
717 rocks, however, the sudden change in the pore pressure distribution due to porosity and
718 permeability creation within the sill also initiates a vortex at the sill tip. This leads to
719 transport of methane-rich fluids towards the sill tip from the surrounding host rock (Figure
720 9l). Eventually the amount of methane stabilizes within the sill (Figure 11). Hydrocarbons
721 entering the sill in this last phase do not experience the extreme temperatures and are unlikely
722 to experience graphitization.

723 **4.3 Implications for flow and methane transport in volcanic basins**

724 Our results have two main implications for fluid flow and methane transport in volcanic
725 basins. First, they demonstrate that part of the generated methane may transform to graphite
726 and thus be permanently stored in fractured sills. Second, the opening of vertical flow paths
727 through fractured sills changes hydrocarbon migration routes compared to impermeable sills,
728 and may thereby affect atmospheric degassing.

729 Despite observable differences in the physical processes at work during sill emplacement in
730 low- vs. high-permeability host rocks, both environments allow for a significant amount of
731 hydrocarbon flushing into a fractured sill at temperatures that are sufficient for graphitization.
732 The occurrence of graphitic bitumen in dykes and as filling of cooling joints indicates that
733 part of the carbon gases rising through permeable sills could be reduced to graphite and be
734 stored as solids (Figure 2, 3, 11). Our results show a plausible mechanism for the creation of
735 these features, as they demonstrate that methane-rich fluids can exploit hydrofractures to flow
736 into the porous and permeable while temperatures in the sill are $>400^{\circ}\text{C}$, i.e., high enough for
737 graphite precipitation (Buseck and Beyssac, 2014). This fraction of the mobilized carbon is
738 trapped permanently underground and is not available for degassing into the atmosphere. The
739 widespread observations of graphite in the field suggests that a significant fraction of
740 hydrocarbons matured in the metamorphic aureoles of sills may not be transported away
741 from the sill. Currently, the fraction of methane transformed to immobile graphite is not
742 known. Our simulations with permeable sills provide a first estimate, showing that up to 10-
743 11,000 tons (6.3-7.5%) of the generated methane experiences temperature $>400^{\circ}\text{C}$ (Figure
744 11a).

745 Sequestration of significant amounts of generated thermogenic carbon into cooling sills by
746 secondary mineral formation is suggested by borehole data in the Karoo Basin, South Africa
747 (Lenhardt et al., 2023). Despite this carbon mass not being available for degassing, a recent
748 investigation on the Karoo Large Igneous Province and the related Toarcian crisis (ca 183
749 Ma) indicates that a total of ca. 20,500 Gt C is needed to replicate the Toarcian $p\text{CO}_2$ and
750 $\delta^{13}\text{C}$ proxy data (Heimdal et al., 2021). Based on existing quantitative model outcomes on the
751 Karoo LIP (Galerie and Hasenclever, 2019), Lenhardt et al., (2023) pointed out that
752 sequestration and degassing are not contradictory but rather point at synchronous processes
753 during sill emplacement and cooling.

754 Additionally, permeable sills favor upward vertical flow and can thus contribute to fluid
755 pressure release below the sill both in low-permeability and high-permeability host rocks.
756 Conversely, impermeable sills favor sustained fluid pressure build-up below the sill and force
757 contact-parallel fluid flow and methane transport towards sill tips (Iyer et al., 2017). These
758 latter mechanisms, combined with a saucer-shaped sill geometry, are the main responsible for
759 the formation of hydrothermal vent complexes at sill tips (Iyer et al., 2017; Galerie and
760 Hasenclever, 2019). In line with field observations in the Neuquén Basin, the simulations
761 considered here use relatively small, flat sills and therefore do not develop sufficient
762 overpressure for hydrothermal vent formation (Galerie and Hasenclever, 2019; Spacapan et
763 al., 2019). They are therefore not suited to quantify the effect of permeable and porous sills
764 on venting. Nevertheless, the results allow us to speculate that for settings in which venting is
765 generally more likely, permeable sills could reduce vent formation potential, or at least
766 reduce atmospheric degassing. This is because the opening of vertical pathways through the
767 sills offers a more direct fluid escape route and helps to dissipate overpressure below the sill.

768 **4.4 Implications for igneous petroleum systems in the Río Grande Valley**

769 Our study also provides further insight into the evolution of the igneous petroleum systems in
770 the Río Grande Valley with respect to the timing of the charging of the igneous hydrocarbon
771 reservoirs. The current conceptual models for these petroleum systems state that a first
772 migration pulse of hydrocarbons into the reservoirs happens when cooling joints open (Witte
773 et al., 2012; Spacapan et al., 2018; Spacapan et al., 2020). Our study demonstrates that
774 hydrocarbons indeed migrate into the intrusions when the cooling joint network forms, but
775 likely experience far too high temperatures (at least 400°C) for them to survive as producible
776 liquids or gases (Figure 8, Figure 11a). Therefore, a survival of hydrocarbons entering the

777 cooling joint network shortly after its creation seems highly unlikely. This result suggests that
778 the igneous reservoirs are charged during the late-stage cooling of the sills, and after
779 significant amounts of the hydrocarbons have flowed through the sills at high temperature.
780 The migration model proposed earlier thus needs to be revised and split into two sub-stages:
781 (1) a first influx of hydrocarbons through still hot sills, with no or very limited survival of
782 hydrocarbons, and (2) a later migration of hydrocarbons within the cooled sill, where the
783 hydrocarbons can survive and be trapped to form producible reservoirs.

784 **4.5 Study limitations and future recommendations**

785 Finally, we present selected recommendations for future work arising from the limitations of
786 our study. Despite the complexity of the current numerical model, some known processes in
787 the host rock are not yet considered. First, mineral precipitation at high temperatures can
788 occur at non-negligible rates and lead to fast porosity-decrease, causing pore pressure
789 increase and possibly fracturing (Townsend, 2018). In addition, buoyancy effects of methane
790 or even two-phase flow are not considered, which may be an important parameter especially
791 for venting (Iyer et al., 2017). Finally, we believe that due to the strong impact on
792 hydrothermal flow patterns, future hydrothermal modelling studies of intrusion in
793 sedimentary basins should consider the possibility of early porosity-permeability generation
794 within the sills, especially in light of growing evidence for gas sequestration in sills
795 worldwide. One important point would be to better constrain the permeability of sill fracture
796 networks. Optimally, we should seek to reach beyond ad-hoc models and develop a physical
797 model that quantifies porosity-permeability evolution under given conditions, e.g., depth,
798 thickness, cooling rate, composition of magmatic intrusions.

799

800 **5 Conclusions**

801 We integrate geological field data with numerical models to investigate hydrothermal
802 transport of hydrocarbon-rich fluids around fractured igneous sills. We use outcrops of
803 fractured sills emplaced in organic-rich shales in the northern Neuquén Basin to establish that
804 sills can become permeable fluid pathways upon solidification while still hot, which affects
805 the fate of locally generated hydrocarbons. The numerical modelling study allows us to
806 understand in detail the hydrothermal flow patterns in response to porosity and permeability
807 generation inside a sill intrusion. This provides new insights into hydrothermal flow in
808 volcanic sedimentary basins, because previous studies commonly assume that sills represent
809 prominently impermeable bodies. The main conclusions of this study are as follows:

810 1. Widespread occurrence of veins with solid, strongly graphitized bitumen as well as
811 cooling joints filled with solid bitumen or organic-rich shales evidence transport of
812 hydrocarbon-rich fluids and liquefied sediments into the sill in a high-temperature
813 (probably $>350^{\circ}\text{C}$), high-pressure environment. This happens within years to several
814 decades after solidification of the sill.

815

816 2. Numerical modelling indicates three flow phases around sills that become porous and
817 permeable upon solidification, which differ markedly from flow around impermeable
818 sills:

819 (1) Contact-parallel flow toward the tip prior to solidification, creating an early
820 plume above the sill tip.

821 (2) Sudden change to vertical flow upon complete solidification. This leads to
822 flow of hydrocarbons from the lower contact aureole upwards into and/or through
823 the sill ("flushing"). This effect is present in both low- and high-permeability host
824 rocks, because hydrofracturing around the sill increases permeability and thus
825 facilitates flow. In addition, hydrocarbons stored in closing hydrofractures are
826 expelled, which, directly around the sill, pushes hydrocarbon-rich fluid back
827 towards the sill.

828 (3) Stabilization of the flow regime with slow rise of hydrocarbon-rich fluids
829 above the sill center, and backward-downward flow towards the sill due to either
830 closing of hydrofractures (low-permeability host) or a vortex driven by the
831 permeable sill (high-permeability host).

832 3. Simulations indicate that flow of methane through the sill occurs at temperatures $>400^{\circ}\text{C}$,
833 which meets the conditions for hydrothermal graphitization. This may explain field
834 observations of graphitic bitumen dykes and could lead to permanent storage of part of
835 the mobilized carbon (estimated up to 7.5%).

836
837 4. Thus, in contrast to proposed conceptual models, flow of hydrocarbons into newly formed
838 cooling joints is likely not a viable migration/charge mechanism for sill reservoirs in the
839 northern Neuquén Basin, as the intrusions are too hot for survival of liquid hydrocarbons.

840
841 5. Permeability creation with the cooling sills does not significantly reduce pressure build-
842 up below the sill, but creates efficient upward pathways for fluid and reduces focusing of
843 flow around the sill tip. With growing evidence for permeable sills in volcanic basins
844 globally, the permeability evolution of sills should be addressed in future modelling
845 studies focused on sill-related venting.

846

847

848 **Author contributions**

849 O. Rabbel: study concept design, field work, evaluation of field data, preparation and
850 evaluation of numerical simulations, manuscript preparation and revision

851 J. Hasenclever: development of numerical code, preparation and evaluation of numerical
852 simulations, preparation and revision of model description and results in manuscript

853 C. Y. Galerne: study concept design, preparation and evaluation of numerical simulations,
854 manuscript preparation and revision

855 O. Galland: study concept design, field work and evaluation of field data, manuscript
856 preparation and revision

857 K. Mair: study concept design, field work and evaluation of field data, manuscript
858 preparation and revision

859 O. Palma: field work, geochemical analysis, evaluation of field data, manuscript preparation
860 and revision

861

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867

868 **Supplementary material**

869 Additional figures and animations of all described simulation runs, plus some additional ones
870 using a thinner sill, can be found in the following FAIR data repository:

871 https://osf.io/28whp/?view_only=5ad1d2cc52844f1cb715516d076a5dd0

872

873 **Competing interests**

874 The authors declare that they have no conflict of interest.

875

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