



1	Equilibrium climate sensitivity increases with aerosol concentration
2	due to changes in rain efficiency
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# 34 Abstract

35 How Earth's climate reacts to anthropogenic forcing is one of the most burning questions faced by today's scientific community. A leading source of uncertainty in 36 37 estimating this sensitivity is related to the response of clouds. Under the canonical 38 climate-change perspective of forcings and feedbacks, the effect of anthropogenic 39 aerosols on clouds is categorized under the forcing component, while the modifications 40 of the radiative properties of clouds due to climate change are considered in the 41 feedback component. Each of these components contributes the largest portion of 42 uncertainty to its relevant category and is largely studied separately from the other. In 43 this paper, using idealized cloud resolving, radiative-convective-equilibrium simulations, with a slab ocean model, we show that aerosol-cloud interactions could 44 45 significantly affect cloud feedback. Specifically, we show that equilibrium climate 46 sensitivity increases under high aerosol concentration due to an increase in the 47 shortwave cloud feedback. The shortwave cloud feedback is enhanced under high 48 aerosol conditions due to a stronger increase in the precipitation efficiency with 49 warming, which can be explained by higher sensitivity of the droplet size and the cloud water content to the CO<sub>2</sub> concentration rise. These results indicate a strong connection 50 51 between cloud feedback and aerosol-cloud interactions.

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### 1. Introduction

54 Estimating Earth's equilibrium climate sensitivity (ECS), defined as the steady-state 55 global mean temperature increase for a doubling of CO<sub>2</sub>, is considered as a first-order, fundamental milestone on the way to understanding and predicting anthropogenic-56 57 driven climate change (Sherwood et al., 2020). Decades of research have tried to 58 accurately quantify ECS, with only limited success. The most probable current ECS 59 estimates are in the range of 2.3-4.5K (Sherwood et al., 2020). The largest source of 60 uncertainty in estimating ECS is related to the response of clouds to the externally forced warming and the feedback of these changes on the climate system (Sherwood et 61 62 al., 2020; Ceppi et al., 2017; Schneider et al., 2017). Clouds strongly modulate Earth's 63 radiation budget by reflecting the incoming shortwave radiation from the sun and by 64 absorbing and re-emitting the terrestrial longwave radiation (Loeb et al., 2018). Thus, changes in the cloud macro-physical properties (such as coverage and vertical extent) 65 66 and micro-physical properties (such as liquid/ice partition or hydrometeors size) due to 67 anthropogenic-driven climate change could significantly alter the climate system's





- response (Gettelman and Sherwood, 2016; Nuijens and Siebesma, 2019; Schneider et
- 69 al., 2017).

70 An important factor in determining cloud feedback magnitude is the sensitivity of the 71 Precipitation Efficiency ( $\epsilon$ ) (Lutsko et al., 2021; Li et al., 2022; Lutsko and Cronin, 72 2018).  $\epsilon$  quantifies the fraction of condensed water in a cloud to reach the surface as 73 precipitation. Using idealized cloud resolving simulations, it was shown that  $\epsilon$  is 74 expected to increase with temperature (Lutsko and Cronin, 2018). The increase in  $\epsilon$ 75 with warming was shown to be mostly driven by an increase in the efficiency with 76 which cloud condensate is converted into precipitation, while changes in the evaporation of falling precipitation was shown to play a smaller role (Lutsko and 77 78 Cronin, 2018).

79 An increase in  $\epsilon$  with warming more efficiently depletes the water from the clouds, thus affecting the radiation budget. On the one hand, increase in  $\epsilon$  with warming was 80 81 suggested to reduce the anvil cloud coverage and hence increase the outgoing longwave radiation (Lindzen et al., 2001; Mauritsen and Stevens, 2015), thus producing negative 82 83 feedback. On the other hand, however, it was recently shown that the longwave effect of an  $\epsilon$  increase is over-compensated for by changes in the shortwave flux (Li et al., 84 85 2019), i.e., a large reduction in the cloud optical depth, driving a reduction in the 86 shortwave cooling effect of clouds, dominates the response.

The efficiency with which cloud condensate is converted into precipitation is closely 87 88 linked to the micro-physical properties of the clouds. The autoconversion of cloud droplets into rain becomes significant when liquid water amount and/or droplet radii 89 90 reach a critical threshold (Freud and Rosenfeld, 2012). An important factor influencing 91 the droplet radii (and also the liquid water amount, to some degree) is the amount of 92 available cloud condensation nuclei (CCN). Generally, an increase in aerosol 93 concentration drives an increase in CCN concentration, which results in more numerous and smaller droplets in the cloud (Twomey, 1974; Warner and Twomey, 1967). The 94 95 smaller droplets require longer time (or equivalently larger vertical distance) in the 96 clouds to grow by diffusion to the critical size enabling precipitation, thus delaying the initial warm rain formation (Rosenfeld, 2000; Dagan et al., 2015b). Therefore, aerosols 97 could affect  $\epsilon$  (Khain, 2009). 98

99 In addition to the effect on rain, aerosols could modify the radiative properties of clouds,

100 by modifying the droplet concentration and size distribution (Twomey, 1974) and by





101	affecting the clouds' macro-physical properties (Albrecht, 1989; Bellouin et al., 2019).
102	These changes to the radiative properties of clouds result in radiative forcing that could
103	affect the sea surface temperature [SST (Bellouin et al., 2019)]. Using cloud-resolving
104	radiative-convective-equilibrium simulations with interactive SST, Khairoutdinov and
105	Yang (2013) showed that the surface temperature decreases by 1.5K with each 10-fold
106	increase in aerosol concentration, an effect quite comparable to a 2.1-2.3K SST
107	warming obtained in a simulation with given (low) aerosol conditions but doubled $\mathrm{CO}_2$
108	concentration.
109	It has been suggested that cloud feedback and aerosol forcing are not independent of
110	each other (Mülmenstädt and Feingold, 2018; Igel and van den Heever, 2021). In

addition, the strong links between  $\epsilon$  and cloud feedback and between  $\epsilon$  and aerosol

concentration merit a dedicated study on the potential mutual CO<sub>2</sub> and aerosol effect on
clouds and thus also on ECS, which is the aim of the current study.

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# 116 **2.** Methods

# 117 Model description and experimental design

The model used herein is the System of Atmospheric Modeling [SAM - (Khairoutdinov 118 119 and Randall, 2003)] version 6.11.7. Subgrid-scale fluxes are parameterized using 120 Smagorinsky's eddy diffusivity model and gravity waves are damped at the top of the 121 domain. The microphysics scheme used is Morrison et al. (2005) 2-moment bulk microphysics. The cloud droplet number concentration source assumes that the number 122 of activated CCN depends on the super-saturation (S) according to a power-law: CDNC 123 =  $N_a S^k$ , where  $N_a$  is the prescribed concentration of CCN active at 1 % supersaturation, 124 and k is a constant (set in this study to 0.4 - a typical value for maritime conditions). 125 126 Changes in  $N_a$  serve as a proxy for the change in aerosol concentration. Three levels of  $N_a$  are considered here, covering an extreme range of conditions – 20, 200 and 2000 127 128 cm<sup>-3</sup>. While this wide range of conditions is unlikely to exist at any given geographical 129 location, they are used here in order to cover the range of possible conditions at different 130 locations and to maximize the effect for establishing better physical understanding. The 131 activation of CCN at the cloud base is parameterized following Twomey (1959), using the vertical velocity and CCN spectrum parameters. The model computes cloud water 132 133 and ice-crystal effective radius for the radiation; thus, the Twomey effect (Twomey,





1977) of both liquid and ice is considered. Direct interactions between aerosols and 134 135 radiation are not considered here. The simulations are conducted in a radiative-convective-equilibrium (RCE) mode and 136 137 generally follow the RCEMIP (RCE model inter-comparison project (Wing et al., 138 2018)) small-domain instructions (but with interactive SST and changes in the CO2 and aerosol concentration). The simulations were performed on a square, doubly periodic 139 140 domain. In this case, we want to avoid the effect of convective self-aggregation on  $\epsilon$ ; 141 thus, the domain size is set to 96x96 km<sup>2</sup>, which was shown to be small enough to 142 prevent convective self-aggregation (Muller and Held, 2012; Lutsko and Cronin, 2018; Yanase et al., 2020). The horizontal grid spacing is set to 1km and 68 vertical levels are 143 used, between 25m and 31km, with vertical grid spacing increasing from 50m near the 144 145 surface to roughly 1km at the domain top. A time step of 10s is used, and radiative fluxes are calculated every 5 min using the CAM radiation scheme (Collins et al., 2006). 146 The output resolution for all fields is 1h. The incoming solar radiation is fixed at 551.58 147 Wm<sup>-2</sup> with a zenith angle of 42.05° (Wing et al., 2018), producing a net insolation close 148 to the tropical-mean value. Convection is initialized with a small thermal noise added 149 150 near the surface at the beginning of the simulation. The initial conditions for the simulations are as in Wing et al. (2018). 151 152 Greenhouse gases are varied for three different levels: pre-industrial level (280 PPM, 153 1xCO<sub>2</sub>), 2 times pre-industrial level (2xCO<sub>2</sub>) and 4 times pre-industrial level (4xCO<sub>2</sub>). As in the case of the aerosol concentrations, the large range of  $CO_2$  conditions covered 154 155 here are used to examine the clouds' sensitivity to greenhouse gas concentrations under 156 a wide range of conditions. Nine different simulations, with all possible combinations 157 of  $N_a$  and CO<sub>2</sub> concentrations, were conducted.

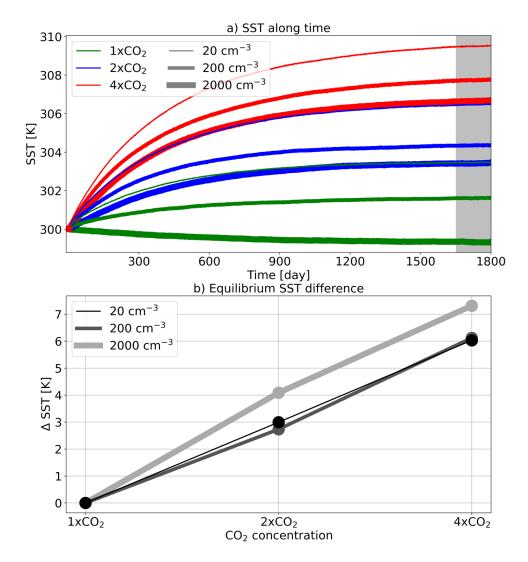
In all simulations, the SST is interactive and predicted by a slab ocean model (SOM). 158 159 The SOM's mixed layer depth is set to 5m, which represented a compromise between a 160 relatively deep layer ( $\geq 10m$ ), which reduces SST noise (Khairoutdinov and Yang, 161 2013), and a relatively shallow layer ( $\ll$ 1m), which requires a shorter computation time for equilibrium (Romps, 2020). As in Romps (2020), the SOM is cooled at a rate of 112 162 Wm<sup>-2</sup> in order to ensure that the simulations with 1xCO<sub>2</sub> are kept at around the initial 163 SST of 300K (Fig. 1). Each simulation was run for 1800 days, which is sufficient for 164 reaching close to equilibrium (the surface energy imbalance is  $\leq 0.1 \text{Wm}^{-2}$  in all 165





- 166 simulations during the last 150 days). The last 150 days of each run are used for
- 167 statistical sampling (gray shading in Fig. 1).

# 168



170Figure 1. a) the sea surface temperature (SST) along time for the different simulations171conducted under different aerosol and  $CO_2$  concentrations. The gray shaded area is172referred to as equilibrium conditions. b) Change in equilibrium SST due to a change in173 $CO_2$  concentration (compared to the  $1xCO_2$  case of each aerosol concentration), for the174different aerosol concentrations (the different curves).

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# 177 **3. Results**

178 Figure 1 presents the SST of the different simulations along time (panel a) and the 179 change in the equilibrium SST with the  $CO_2$  concentration for the different  $N_a$  cases 180 (panel b). As expected, the equilibrium SST (gray shading in Fig. 1a) increases with the 181  $CO_2$  concentration and decreases with  $N_a$  concentration. However, the rate of increase in equilibrium SST with  $CO_2$  concentration increases under extremely high  $N_a$ 182 183 concentrations (2000 cm<sup>-3</sup>), compared with the low and medium  $N_a$  concentrations (20 and 200 cm<sup>-3</sup>, respectively - Fig. 1b). Calculating the average ECS based on the three 184 combinations available for each  $N_a$  condition [2xCO<sub>2</sub>-1xCO<sub>2</sub>, 4xCO<sub>2</sub>-2xCO<sub>2</sub> and 185 186  $(4xCO_2-1xCO_2)/2$ ], demonstrates that it increases with  $N_a$  from 3.0K at the lowest  $N_a$ to 3.7K at the highest  $N_a$  (i.e., a 23% increase – Table 1). 187

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189	Table 1. Average equilibrium climate sensitivity (ECS), cloud-feedback parameter ( $\lambda_{cloud}$ ),			
190	hydrological sensitivity ( $\eta$ ), and change in precipitation efficiency ( $\Delta\epsilon$ ) of the three			
191	combinations available for each N <sub>a</sub> condition [2xCO <sub>2</sub> -1xCO <sub>2</sub> , 4xCO <sub>2</sub> -2xCO <sub>2</sub> and 4xCO <sub>2</sub> -			
192	1xCO <sub>2</sub> ]. For the calculation of the average ECS, the difference between 4xCO <sub>2</sub> and 1xCO <sub>2</sub>			
193	is divided by 2. The rest of the quantities are normalized by the SST change between the			
194	relevant simulations.			
	$N_{\rm c}$ [cm <sup>-3</sup> ] ECS [K] $\lambda_{\rm clust}$ [W m <sup>-2</sup> K <sup>-1</sup> ] $n$ [% K <sup>-1</sup> ] $\Lambda_{\rm c}$ [% K <sup>-1</sup> ]			

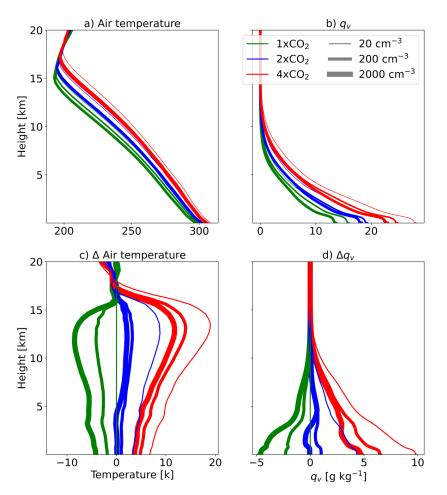
$N_a$ [cm <sup>-3</sup> ]	ECS [K]	$\lambda_{cloud} [W m^{-2} K^{-1}]$	$\eta \ [\% \ \mathrm{K}^{-1}]$	$\Delta \epsilon \ [\% K^{-1}]$
20	3.0	-0.45	3.8	1.2
200	3.1	-0.38	4.3	1.3
2000	3.7	-0.08	4.6	2.7

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Figure 2 presents the time and domain mean vertical profiles of temperature and water 196 197 vapor mixing ratio  $(q_y)$  in the different simulations (panels a and b) and their difference from the simulation with the lowest  $N_a$  and  $CO_2$  concentrations (panels c and d). It 198 199 demonstrates, as expected, that the vertical profile of air temperature is set by the surface temperature (increases with  $CO_2$  concentrations and decreases with  $N_a$ ) with an 200 201 amplification of the change at the upper troposphere, as the profiles follow the moist adiabatic lapse-rate. It also shows that  $q_c$  increases with the temperature, as expected 202 203 (Held and Soden, 2006).







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Figure 2. Time and domain mean vertical profiles of air temperature and water vapor mixing ratio  $(q_v)$  in the different simulations (a and b) and how they differ from the simulation with the lowest  $N_a$  and CO<sub>2</sub> concentrations (panels c and d).

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In order to understand the increase in ECS with  $N_a$ , we next examine the top-ofatmosphere (TOA) energy budget. Figure 3 presents the change in the net shortwave and longwave TOA energy gain ( $R^{SW}$  and  $R^{LW}$ , respectively) with the CO<sub>2</sub> concentration for the different  $N_a$  conditions. In addition, Fig. 3 presents the change in the cloud radiative effect (CRE) with increasing the CO<sub>2</sub> concentration, where CRE is computed by subtracting the clear-sky from the all-sky TOA radiative fluxes ( $R-R_{clear-sky}$ ), again for the shortwave and longwave separately (CRE<sup>SW</sup> and CRE<sup>LW</sup>,



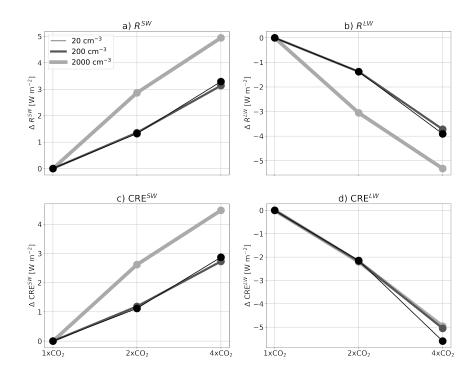


218	respectively). Figure 3a and b demonstrates that under equilibrium conditions $R^{SW}$
219	increases, while $R^{LW}$ decreases with the CO <sub>2</sub> concentration. However, the rate of change
220	in both $R^{SW}$ and $R^{LW}$ is much faster under the high $N_a$ conditions than under the low
221	and medium $N_a$ conditions. The trend in CRE <sup>SW</sup> under the different $N_a$ conditions (Fig.
222	3c) resembles the trend in $R^{SW}$ , suggesting that the clouds' response dominates the
223	changes in the TOA shortwave fluxes. $\mbox{CRE}^{\mbox{\tiny LW}},$ on the other hand, decreases at a similar
224	rate with CO <sub>2</sub> concentration for the different $N_a$ conditions (Fig. 3d). Thus, the different
225	decrease rates in $R^{LW}$ with CO <sub>2</sub> concentration for the different $N_a$ conditions (Fig. 3b)
226	must be driven by clear-sky changes (specifically, the plank and the lapse-rate
227	feedbacks – see Fig. 2 above).
228	In Table 1 above, we estimate the average cloud radiative feedback ( $\lambda_{cloud})$ as the change
229	in CRE with increasing surface temperature, i.e., $\lambda_{cloud} = dCRE/dT$ , for the different $N_a$
230	conditions. The table shows that $\lambda_{cloud}$ becomes less negative with the increase in $N_a$ ,
231	leading to higher climate sensitivity. The differences in the values of $\lambda_{cloud}$ between the
232	different $N_a$ conditions is mostly derived from the shortwave part of the spectrum (Fig.
233	3).
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Figure 3. The change in the net top-of-atmosphere energy gain (R) in the shortwave (a) and in the longwave (b), and the change in the cloud radiative effect (CRE) in the shortwave (c) and in the longwave (d), due to a change in the CO<sub>2</sub> concentration (compared to the 1xCO<sub>2</sub> case of each aerosol concentration), for the different aerosol concentrations (the different curves).

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Thus far, we have seen that the ECS increases with  $N_a$  (Fig. 1 and Table 1) and that this 243 increase can be explained by changes in  $\lambda_{cloud}$  (Table 1) and specifically in CRE<sup>SW</sup> (Fig. 244 3). To understand the changes in the cloud properties driving the changes in  $\lambda_{cloud}$ , and 245 246 hence also in ECS, under the different  $N_a$  conditions, in Fig. 4 we present the change in cloud liquid water path (CWP), ice water path (IWP), rain water path (RWP) and cloud 247 fraction (CF) with increasing CO<sub>2</sub> concentrations for the different  $N_a$  conditions. The 248 figure shows that the CWP decreases with the CO2 concentrations at a much faster rate 249 (about 3 times faster) under the highest  $N_a$  conditions compared to the low and medium 250 251  $N_a$  conditions (Fig. 4a). The changes in the IWP, on the other hand, are about an order 252 of magnitude smaller than the changes in CWP and are not consistent in sign for the 253 different  $N_a$  conditions (Fig. 4b). The RWP increases with the CO<sub>2</sub> concentrations at a 254 slightly faster rate (about 20% faster) under the highest  $N_a$  conditions compared to the





- low and medium  $N_a$  conditions (however the response is non-monotonic with  $N_a$  Fig. 4c). The CF decreases with the CO<sub>2</sub> concentrations, at a similar rate for the different  $N_a$
- $257 \qquad \text{conditions (about 1.5\% decrease in CF for each doubling of the CO_2 concentrations -}\\$
- 258 Fig. 4d).
- 259 The faster decrease in CWP with  $CO_2$  concentrations under high  $N_a$  conditions drives the faster increase in CRE<sup>SW</sup> as the clouds become less opaque in the shortwave. We 260 note that the difference in  $CRE^{SW}$  trend under different  $N_a$  conditions could not be 261 explained by the minor differences in the CF trends. In addition, the small differences 262 263 in the IWP between the different  $N_a$  conditions are consistent with the small differences in the CRE<sup>LW</sup> seen above. The general increase in RWP with CO<sub>2</sub> concentrations is 264 consistent with an increase in rain efficiency with warming (Lutsko and Cronin, 2018), 265 266 as elaborated below.
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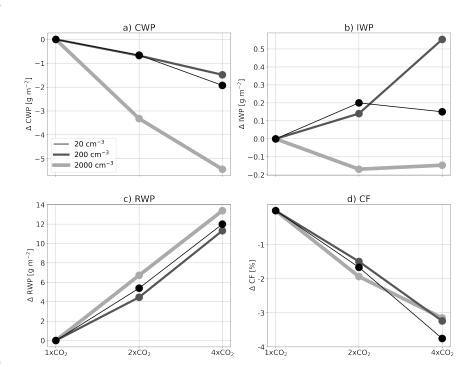




Figure 4. The change in: a) cloud liquid water path (CWP), b) ice water path (IWP, c) rain water path (RWP), and d) cloud fraction (CF) due to a change in the CO<sub>2</sub> concentration (compared to the 1xCO<sub>2</sub> case of each aerosol concentration), for the different aerosol concentrations (the different curves).



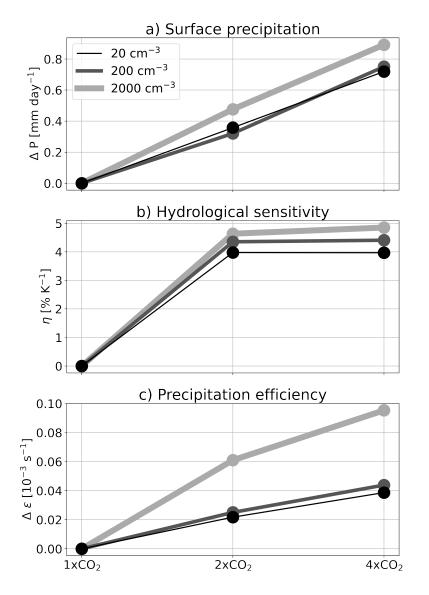


274	Figure 4 suggests that the largest difference in the cloud response to $\mathrm{CO}_2$ under different
275	$N_a$ conditions is due to changes in CWP. The higher sensitivity of CWP to CO <sub>2</sub>
276	concentration under higher $N_a$ conditions can explain the higher $\lambda_{cloud}$ and thus also the
277	larger ECS. Hence, the question arises: What causes the faster reduction in CWP with
278	$CO_2$ concentration under high $N_a$ conditions? A major sink for CWP is via precipitation.
279	Hence, in Fig. 5 we present the change in the mean surface precipitation rate, the
280	hydrological sensitivity ( $\eta$ - the rate of change in the surface precipitation per 1K
281	increase in surface temperature) and the precipitation efficiency ( $\epsilon$ - calculated
282	following Li et al. (2022) as the ratio of surface precipitation-to-condensed water path,
283	i.e., CWP+IWP+RWP). As expected, the surface precipitation increases with CO <sub>2</sub> (i.e.,
284	$\eta$ is positive) and so does $\epsilon$ (Lutsko and Cronin, 2018). This is true for all $N_a$ conditions.
285	However, the rates of increase in surface precipitation and $\epsilon$ with CO <sub>2</sub> concentration
286	are higher under the highest $N_a$ conditions (see also Table 1). We note that the larger
287	rate of increase in surface precipitation under the highest $N_a$ conditions is not solely due
288	to the higher surface temperature increase, as $\eta$ also increases with $N_a$ .
289	The much larger (more than double- Table 1) rate of increase in $\epsilon$ with the CO <sub>2</sub>
290	concentration under the highest $N_a$ conditions depletes the cloud water more efficiently
291	from the atmosphere, leading to a faster reduction in CWP with $\mathrm{CO}_2$ concentration (Fig.
292	4), which in turn leads to higher $\lambda_{cloud}$ and ECS. The faster increase in RWP with CO <sub>2</sub>
293	concentration under the highest $N_a$ conditions presented in Fig. 4c is consistent with
294	this explanation.
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Figure 5. The change in: a) surface precipitation, b) hydrological sensitivity ( $\eta$ ), and c) precipitation efficiency ( $\epsilon$ ) due to a change in the CO<sub>2</sub> concentration (compared to the 1xCO<sub>2</sub> case of each aerosol concentration), for the different aerosol concentrations (the different curves).

The last open question is why  $\epsilon$  increases faster with CO<sub>2</sub> concentration under the highest  $N_a$  conditions. The increase in  $\epsilon$  with warming was shown to be mostly driven by an increase in the efficiency with which cloud condensate is converted into





307 precipitation (Lutsko and Cronin, 2018). As was mentioned in the introduction, the 308 conversion of cloud condensate into precipitation (or autoconversion of cloud droplets) 309 becomes significant only when liquid water amount and/or droplet radii reach a critical 310 threshold (Freud and Rosenfeld, 2012). To understand the faster  $\epsilon$  increases with CO<sub>2</sub> concentration under the highest  $N_a$  conditions, we present the histograms over the 311 312 domain and time (during the last 150 days of the simulations based on 3D output in 1hour resolution) of liquid cloud droplets mixing ratio ( $q_c$  – Fig. 6) and mean cloud 313 droplet radii ( $\overline{r_c}$  – Fig. 7) around the height of the maximum in cloud droplet effective 314 315 radii (1950m) and its mean sensitivity to doubling of  $CO_2$  concentration for each  $N_a$ condition. 316

Figure 6 demonstrates that the cut-off of the  $q_c$  distribution (the mixing ratio for which the probability density function starts to decrease sharply) increases with the CO<sub>2</sub> concentration and decreases with the aerosol concentration. However, the sensitivity of the relatively large  $q_c$  with CO<sub>2</sub> concentration is significantly larger under high aerosol concentrations compared to the lower aerosol concentrations (Fig. 6b). The larger relative increase in high  $q_c$  promotes the autoconversion process and hence enhances  $\epsilon$ , more under high aerosol concentrations than under low aerosol concentrations.

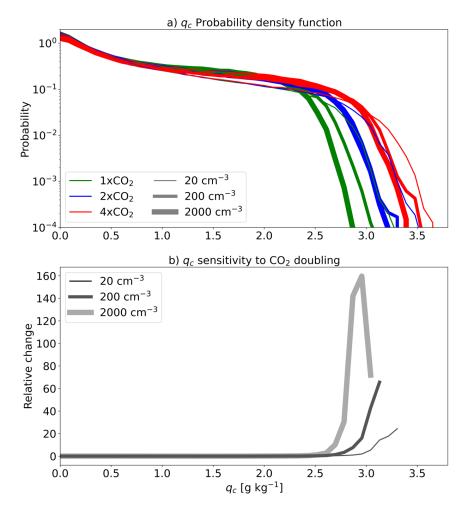
324 Figure 7 demonstrates, in line with expectations, that  $N_a$  has a strong effect on  $\overline{r_c}$ . In addition, it shows that under all  $N_a$  conditions,  $\overline{r_c}$  increases with the CO<sub>2</sub> concentration. 325 This could be explained by the increase in the availability of water vapor (Fig. 2), 326 327 which, for a given  $N_a$  conditions, enable larger diffusional growth of the droplets. Here 328 again, the highest  $N_a$  conditions demonstrate the largest sensitivity of  $\overline{r_c}$  to CO<sub>2</sub> concentration, especially at the right-hand side of the distribution (Fig. 7b). This could 329 330 be explained by the fact that under these high  $N_a$  conditions, the cloud droplet growth 331 is primarily limited by the availability of water vapor, as large number of droplets 332 compete for the available water vapor (Koren et al., 2014; Dagan et al., 2015a; Reutter et al., 2009). Similarly to the  $q_c$  case, the larger relative increase in the relatively large 333 334 droplets promotes the autoconversion process and hence enhances  $\epsilon$ , more under high aerosol concentrations than under lower aerosol concentrations. 335

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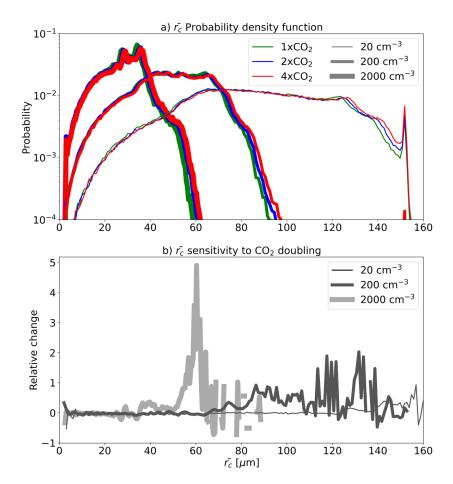


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Figure 6. Probability density functions (PDF) of the cloud droplet mixing ratio  $(q_c)$  for the different simulations (a), and the mean sensitivity of the  $q_c$  PDF to a doubling of the CO<sub>2</sub> concentration based on the three combinations available for each  $N_a$  condition [2xCO<sub>2</sub>-1xCO<sub>2</sub>, 4xCO<sub>2</sub>-2xCO<sub>2</sub> and (4xCO<sub>2</sub>-1xCO<sub>2</sub>)/2] (b), calculated for the heights around which the cloud droplet effective radii reach a maximum (1950m) and using 3-D files output every hour of the last 150 days of the simulations. Note the logarithmic scales for the yaxes of a.







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Figure 7. Probability density functions (PDF) of cloud droplet mean radii ( $\overline{r_c}$ ) for the different simulations (a), and the mean sensitivity of the  $\overline{r_c}$  PDF to a doubling of the CO<sub>2</sub> concentration based on the three combinations available for each  $N_a$  condition [2xCO<sub>2</sub>-1xCO<sub>2</sub>, 4xCO<sub>2</sub>-2xCO<sub>2</sub> and (4xCO<sub>2</sub>-1xCO<sub>2</sub>)/2] (b), calculated for the heights around which the cloud droplet effective radii reach a maximum (1950m) and using 3-D files output every hour of the last 150 days of the simulations. Note the logarithmic scales for the yaxes of a.

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#### 357 4. Summary and conclusions

The role of clouds in a climate-change is manifested by two pathways: (1) effects of anthropogenic aerosol on clouds, and (2) feedback that clouds exert on the changing climate. These two pathways are usually studied separately, and even by different





- 361 scientific communities. In this paper, we demonstrate that the two pathways are closely
- 362 linked to each other and should be examined concurrently.
- Using long, idealized RCE simulations over a small domain with a slab ocean model,
  we demonstrate that the ECS, i.e., the increase in surface temperature under equilibrium
- conditions due to doubling of the CO<sub>2</sub> concentration, increases with the aerosol concentration. The ECS increase is explained by a faster increase in precipitation efficiency with warming under high aerosol concentrations, which more efficiently depletes the water from the cloud and thus is manifested as an increase in the cloud feedback parameter. The precipitation efficiency increases faster under high aerosol concentration due to a higher sensitivity of the relatively high liquid water mixing ratios and the relatively large mean droplet sizes to a CO<sub>2</sub> concentration increase.

The results presented here are based on idealized simulations over a small domain.
Under more realistic conditions, other processes, not included here, that could affect
the precipitation efficiency and hence the general trend will be introduced. In particular,

convective self-aggregation could be of interest as, while it is inhibited in the small
domain used here, it was shown to affect precipitation efficiency (Lutsko et al., 2021)
and to be affected by aerosols (Nishant et al., 2019). Other processes that should be
accounted for in future research include the presence of large-scale circulation and
direct aerosol radiative effects (Dagan et al., 2019; Dingley et al., 2021).

The results presented here suggest a strong connection between cloud feedback and aerosol-cloud interactions. The regulation of aerosol emissions is known to be more effective than the effort to reduce greenhouse gas emissions. This, together with the short lifetime of aerosols in the atmosphere, has resulted in a reduction in the value of the global mean aerosol effective radiative forcing in recent years (Quaas et al., 2022). If the conclusions of this paper hold under higher levels of complexity (e.g., large-scale circulation, convective self-aggregation, etc.) this might mean that the reduction in

soo enculation, convective sen aggregation, etc.) and might mean that the reduction in

387 global aerosol emissions could lead to a reduction in ECS, which could compensate, at

least partially, for the reduction in the negative forcing induced by aerosols (Quaas et

al., 2022; Bellouin et al., 2019), thus providing yet additional motivation for reducing

aerosol emissions globally.

391

# 392 Code availability

393 SAM is publicly available at: http://rossby.msrc.sunysb.edu/~marat/SAM.html





# 395 Data availability

- 396 The data presented in this study will become publicly available via zenodo prior to
- 397 publication.

398

# **399** Author contributions

400 GD carried out the simulations and analyses presented and prepared the article.

401

### 402 Competing interests

- 403 The authors declare that they have no conflict of interest.
- 404

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