Dansgaard-Oeschger events in climate models: Review and baseline MIS3 protocol

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Abstract.

Dansgaard-Oeschger (D-O) events, millennial-scale climate oscillations between stadial and interstadial conditions (of up to 10-15°C in amplitude at high northern latitudes), occurred throughout the Marine Isotope Stage 3 (MIS3; 27.8 – 59.4 ka) period. The climate modelling community up to now has not been able to answer the question: Are our climate models too stable to

- 5 simulate D-O events? To address this, this manuscript lays the ground-work for a MIS3 D-O protocol for general circulation models which are used in the International Panel for Climate Change (IPCC) assessments. We review: D-O terminology, community progress on simulating D-O events in these IPCC-class models (processes and published examples), and evidence about the boundary conditions under which D-O events occur. We find that no model exhibits D-O like behaviour under pre-industrial conditions. Some, but not all, models exhibit D-O like oscillations under MIS3 and/or full glacial conditions.
- 10 Greenhouse gases and ice-sheet configurations are crucial. However most models have not run simulations of long enough duration to be sure which models show D-O like behaviour, under either MIS3 or full glacial states. We propose a MIS3 baseline protocol at 34 ka, which features low obliquity values, medium-to-low MIS3 greenhouse gas values and the intermediate ice-sheet configuration which our review suggests are most conducive to D-O like behaviour in models. We also provide a protocol for a second freshwater (Heinrich-Event preconditioned) experiment, since previous work suggests that this variant
- 15 may be helpful in preconditioning a state in models which is conducive to D-O events. This review provides modelling groups investigating MIS3 D-O oscillations with a common framework, which is aimed at 1) maximising the chance of the occurrence of D-O like events in the simulations; 2) allowing more precise model-data evaluation and; 3) providing an adequate central point for modellers to explore model stability.

1 Introduction

20 During a Dansgaard-Oeschger (D-O) event, Greenland transitions between cold stadial (GS) and warmer Greenland Interstadial (GI) conditions. The warming can occur within a decade (Kindler et al., 2014; Huber et al., 2006), whilst cooling occurs over a much longer period that is typically several centuries in length. During a warming phase, surface air temperatures over Greenland increase by 10-15°C (Andersen et al., 2006; Kindler et al., 2014; Huber et al., 2006). D-O events are best documented during Marine Isotope Stage 3 (MIS3; between 27.8 – 59.4 thousand of years BP, hereafter ka Goni and Harrison,

- 25 2010), including being recorded in several ice cores from Greenland (Fig. 1 Johnsen et al., 2001). Whilst the D-O event recorded in these cores are renowned, the events are global in nature (Voelker et al., 2002; Sanchez Goñi and Harrison, 2010; Sánchez Goñi et al., 2017), with known climate signatures including imprints in surface temperature and the hydrological cycle at high northern latitudes (Andersen et al., 2004; Thomas et al., 2009; Seierstad et al., 2014), in the tropics (Deplazes et al., 2013; Baumgartner et al., 2014; Adolphi et al., 2018), in Eurasia (Genty et al., 2003; Wang et al., 2008; Jacobel et al.,
- 2017; Rousseau et al., 2017), and in North and South America (Wang et al., 2004; Wagner et al., 2010; Asmerom et al., 2010; Deplazes et al., 2013; Vanneste et al., 2015). While there are no Greenland ice core records of the previous glacial (MIS6 around 140-190 ka), speleothems and Antarctic ice cores indicate that it is extremely likely that D-O events also occurred during MIS6 and earlier glacial periods (Lang et al., 1999; Uriarte, 2019; Landais et al., 2004; Turner and Marshall, 2011; Barker et al., 2011; Lambert et al., 2012). This observational evidence shows that D-O events do not occur under interglacial or full Last Glacial Maximum conditions (Galaasen et al., 2014; Tzedakis et al., 2018; Galaasen et al., 2020).
- 35 or full Last Glacial Maximum conditions (Galaasen et al., 2014; Tzedakis et al., 2018; Galaasen et al., 2020). In 2011, Valdes (2011) argued that climate models used in the assessments of the Intergovernmental Panel on Climate Change (IPCC) have not proved their ability to simulate D-O events. This has several implications for the delivery of accurate projections of climate change, within the context of tipping points and abrupt climate change (Brovkin et al., 2021). Whilst in the intervening years a number of models have captured key features of D-O events through AMOC hysteresis behaviour
- 40 and/or produced D-O type millennial-scale variability under a range of forcings (Brown and Galbraith, 2016; Galbraith and de Lavergne, 2019; Klockmann et al., 2018; Peltier et al., 2020; Armstrong et al., 2021; Zhang et al., 2021; Vettoretti et al., 2022), we still do not know if climate models are too stable because too few models have run and published an appropriate simulation. This deficiency is related to both the computational expense which prevents models from being run for the longer time periods needed for investigating D-O events and to the lack of an agreed appropriate experimental set-up. The limited knowledge of pre-Last Glacial Maximum (LGM) boundary conditions, in particular in the case of the ice sheet height and
- 45 knowledge of pre-Last Glacial Maximum (LGM) boundary conditions, in particular in distribution, makes it challenging to generate an appropriate MIS3 experimental set-up.

An important question is if model stability is caused by the model parameters and MIS3 conditions are such that the models are in a mono-stable state, in an oscillatory state or if the models exhibits bi-modality where noise-induced transitions are not induced due to too low model variability (Ditlevsen and Johnsen, 2010). Previous studies have questioned the significance

50 of the periodic occurrence of DO events in MIS3 (~ 1470) (Ditlevsen et al., 2007). If the full glacial period is included, the distribution of waiting time between DO-events is consistent with a random process (Ditlevsen et al., 2005). Durations of stadials vs. interstadials indicate correlations with global ice volume and orbital parameters (Lohmann and Ditlevsen, 2018; Mitsui and Crucifix, 2017), thus underpinning the decision to focus on MIS3 boundary configurations.

Whether models can simulate abrupt changes is a crucial research question: if the current IPCC-class models are too stable
to simulate D-O events, their ability to predict future abrupt transitions, and their use in identifying tipping points is doubtful.
For example, a tipping point may have been recently reached in the Arctic's Barents Sea (Barton et al., 2018; Tesi et al., 2021);
sea ice loss in the area is linked with enhanced heat transport via an intensified throughflow, or "Atlantification" (Årthun et al., 2012; Polyakov et al., 2017). In addition, future enhanced precipitation, decline in Arctic sea ice and melting of glaciers and ice sheets could intensify the supply of freshwater to the North Atlantic and Arctic which could lead to the reorganization of the

60 Atlantic circulation and tip the energy distribution between South and North in a similar way as occurred during D-O events (Lenton et al., 2008). If climate models do not reliably simulate past tipping events, it suggests that simulations of the coming century may be giving us a false sense of security.

Coupled Model Intercomparison Project (CMIP) coordinates and designs climate model protocols for the past, present and future climates, and has become an indispensable tool to facilitate our understanding of climate change (IPCC, 2013; Eyring

et al., 2016). The Paleoclimate Model Intercomparison Project 4 (PMIP4) is one of the individual Model Intercomparison Projects which took part in CMIP6 (Kageyama et al., 2018). The design of a common MIS3 experimental protocol would allow the modelling community to address the questions posed above.

This manuscript compiles current information about unforced D-O like oscillations in CMIP5/CMIP6 models and discusses the boundary conditions and mechanisms responsible for these oscillations. Given the nomenclature on D-O events varies

- 70 throughout the literature. Firstly, Table 1 and Figure 1 provide a framework for a more consistent terminology for use within this proposed MIS3 DO protocol. Secondly, we review the literature to ascertain whether models reproduce D-O like events under MIS3, or other, climate conditions. We then use this information to develop a protocol for the simulations of D-O events. This protocol focuses on Marine Isotope Stage 3 (MIS3) partly because of the excellent records of D-O events during this period (Schulz et al., 1999) but also because, as our synthesis shows, MIS3 conditions are also conducive to promoting D-O
- 75 like events in some models. Given that D-O events did not occur under full glacial conditions in the last glacial period, the proposed modelling protocol is an important improvement on the use of an LGM PMIP protocol. It will undoubtedly help to shed light on the mechanism and processes involved in millennial-scale oscillations during MIS3. The common MIS3 climate modelling protocol is aimed at: 1) maximising the chance of the occurrence of D-O like events in the simulations; 2) improving model-data evaluation and; 3) providing an adequate central point for modellers to also explore model stability. In addition to
- 80 the protocol for a baseline simulation, we also outline a protocol for a Heinrich-Event preconditioned (freshwater) experiment. These protocols provide a common framework for model experiments to explore cold-period instabilities using commonly specified greenhouse gas (GHG), ice sheet, insolation, and freshwater-related forcings.

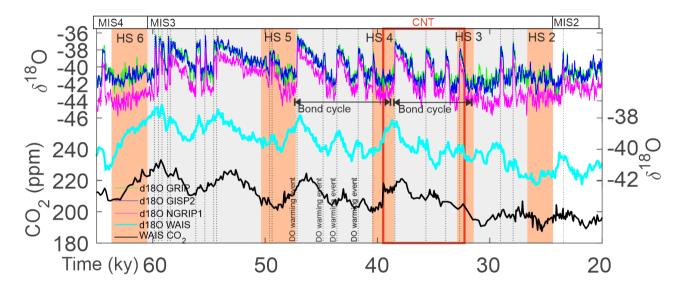


Figure 1. MIS3 ice core records and nomenclature. Stable water isotope and CO_2 measurements from Antarctic and Greenland ice cores (Bauska et al., 2021; NGRIP Project Members, 2004; Kindler et al., 2014). See also Table 1 for D-O nomenclature. The "cnt" red box indicates the 38 to 32 ka period proposed for the MIS3 baseline experiment.

Table 1: Terminology

| Term | Description |
|---------------------------|--|
| Abrupt change | We follow the IPCC Assessment Report 4 (IPCC AR4) definition of abrupt event/change (Meehl et al., 2007; Brovkin |
| | et al., 2021). This term refers to a large-scale change, which is much faster than the change in the pertinent forcing |
| | (e.g. rising atmospheric CO ₂ concentrations). |
| Tipping point | This term refers to a critical threshold at which a small perturbation can qualitatively modify the development or state |
| | of a system (Lenton et al., 2008). |
| Tipping element | This term describes large-scale components of the Earth system that could pass a tipping point (Lenton et al., 2008). |
| | Earth system components are the ocean, atmosphere, cryosphere, anthroposphere and biosphere, which have further |
| | important sub-components e.g. the meridional ocean circulation, the monsoon systems, sea ice, and various ecosystems |
| | (Brovkin et al., 2021). |
| D-O event | During the Last Glacial period, a series of dramatic climatic fluctuations occurred in the North Atlantic. These are |
| | known as D-O events, during which atmospheric and oceanic conditions alternated between relatively mild (intersta- |
| | dial) and full glacial (stadial) conditions (Dansgaard et al., 1982; Johnsen et al., 1992). Around 25 abrupt transitions |
| | (each completed within a decade) from stadial to interstadial conditions occurred during the Last Glacial period and |
| | their amplitude vary from 5 to 16°C (Landais et al., 2004; Huber et al., 2006; Kindler et al., 2014). The duration of |
| | interstadials varies from approximately a century to many millennia(Rasmussen et al., 2014). |
| D-O type oscillations | For the purpose of this MIS3 DO protocol, the term of D-O type oscillation refers to D-O scale climate variability |
| | reproduced by climate models, comparable to the D-O events observed in the Greenland ice core record. |
| Greenland Stadial and In- | We follow the INTIMATE (Integration of Ice core, Marine and Terrestrial records of the North Atlantic) definition |
| terstadial | of stadial/interstadial terms (Rasmussen et al., 2014). The Greenland Interstadials (GI) and Greenland Stadials (GS) |
| | periods terms are the Greenland expressions of the D-O events and represent warm and cold phases of the NA area, |
| | respectively. |
| Heinrich events | These are defined by the presence of layers of ice-rafted debris (IRD) of primarily (not exclusively) Laurentide origin |
| | in North Atlantic sediment cores (e.g. Heinrich, 1988; Hemming, 2004). Heinrich events have been observed during |
| | some of the longer stadials, but likely do not cover the entire period of these longer stadials (Roche et al., 2004; Marcott |
| | et al., 2011). |
| Heinrich Stadial (HS) | This term refers to a stadial containing a specific Heinrich Event. Rasmussen et al. (2014) indicates that the term of HS |
| | can refer to the complete stadial period, or to part of a stadial only, characterized by changes shown in proxies of IRD, |
| | AMOC or SST (Barker et al., 2009). |
| Bond cycle | D-O events tend to follow a pattern of diminishing amplitude (or a general cooling trend of the GSs) following each |
| | HE (Bond et al., 1992; Alley, 1998; Alley et al., 1999; Clark et al., 2007; Rousseau et al., 2022). These cycles of HE |
| | grouped D-O events were named Bond cycles by Broecker (1994) and Alley (1998). The average gap between HEs is |
| | around 7 ka, so this is the average length of a Bond Cycle (Clark et al., 2007). |

2 Review of spontaneous D-O type quasi-oscillations in coupled climate models

85 We compile published evidence of long unforced quasi-oscillations (in the Atlantic Meridional Circulation; AMOC) in IPCCclass models under all climate states in Table A1, alongside glacial boundary condition simulations which do not show D-O type oscillations (Table A3). This permits us to explore the questions of: what proportion of models exhibit D-O like behaviour; which boundary conditions are most conducive to this; and what mechanisms are common to the modelled D-O behaviours. A number of PI/present-day model simulations exhibit spontaneous centennial-length cold events (Table A1), however, they do

90 not appear to be D-O like events. We deal with these first.

Under pre-industrial greenhouse gas (GHG) forcing and present-day ice sheets, spontaneous centennial-length cold events that last around 100-200 years occur in four IPCC-class models (Table A1). EC-Earth and Community Climate System Model version 4 (CCSM4) show high atmospheric blocking over the eastern subpolar gyre that causes a cold event under pre-industrial boundary conditions (Drijfhout et al., 2013; Kleppin et al., 2015, Table A1). ECHAM6-FESOM also produces cooling events

- 95 under pre-industrial conditions due to sudden reductions of deep water convection and increase of sea ice cover in the Labrador Sea (Sidorenko et al., 2015). Changes in convection also occur in the Kiel Climate Model (KCM; Martin et al., 2015), however here centennial-scale variability of the AMOC is linked to variability in Southern Ocean convection. Unlike the CCSM4 and the EC-Earth models, the KCM and ECHAM6-FESOM studies do not indicate an active role of the atmosphere. Although these four models all show abrupt spontaneous cooling events under pre-industrial boundary conditions, these events do not have the typical saw-tooth characteristics, or longer timescales, of D-O type events.
- Regular cycles of D-O type quasi-oscillations are found in UofT CCSM4 under LGM boundary conditions (Peltier and Vettoretti, 2014). The initiation of the abrupt D-O type warming events is associated with the opening of a large polynya over the Irminger Sea (Vettoretti and Peltier, 2016) (Table A1). The AMOC spontaneously exhibits D-O like quasi-oscillations (Peltier et al., 2020). The Peltier et al. (2020) salt oscillator is maintained by the salinity gradient between the subtropical gyre
- and the Northern North Atlantic. Although UofT CCSM4 is the only model to show long unforced quasi-oscillations in the AMOC under full glacial conditions, most of the other PMIP4 LGM simulations (Kageyama et al., 2021a) have not been run long enough to be sure that such oscillations would not arise if they were run for longer (see Table B1). Having said that, ideally models should not show oscillatory D-O type behaviour when configured under a full glacial climate state, given that in reality D-O events do not occur under full glacial conditions (Huber et al., 2006; Galaasen et al., 2014; Kindler et al., 2014;

D-O type quasi-oscillations are also found in MIROC4m under mid-glacial conditions (Kuniyoshi et al., 2022). Some aspects of the D-O warming mechanism observed in the UofT-CCSM4, in particular the spatial location of the opening of a big polynya in the Irminger Sea, determining the stadial-interstadial transition, is also identified in MIROC4m (Kuniyoshi et al., 2022) (Table A1).

- Under late glacial conditions, at 30 ka, a quasi-oscillating AMOC is produced by the HadCM3 model (Armstrong et al., 2021) and results from a North Atlantic salt oscillator mechanism similar to that in UofT CCSM4 (Peltier and Vettoretti, 2014; Vettoretti and Peltier, 2016; Peltier et al., 2020). The HadCM3 model also shows millennial-scale climate oscillations triggered by deglacial meltwater discharge in LGM simulations (Romé et al., 2022). Under intermediate glacial conditions (MIS3: 40-32 ka), the COSMOS model shows spontaneous millennial-scale climate oscillations triggered solely by orbitally driven insolation
- 120 changes (Zhang et al., 2021). Variations in either obliquity or eccentricity-modulated precession lead to climate variations over the tropical and subpolar North Atlantic which exert opposite effects on AMOC strength, and hence result in an oscillatory

¹¹⁰ Tzedakis et al., 2018).

climate regime (Zhang et al., 2021). The CM2Mc model also produces somewhat smoothed quasi-oscillating AMOC under intermediate MIS3-like boundary conditions, with a present-day ice sheet distribution in combination with a CO_2 concentration of 180 ppm and low obliquity (22°) (Brown and Galbraith, 2016; Galbraith and de Lavergne, 2019) (Table A1). The MPI-ESM

model exhibits more abrupt D-O like quasi-oscillations with a present-day ice sheet distribution in combination with CO_2

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concentrations ranging between 190-217 ppm (Table A1; Klockmann et al., 2018, 2020).

In contrast to the above, neither NorESM nor CCSM3 produce D-O type events or quasi-oscillations under MIS3 conditions (38 ka) (Table A3; Guo et al. (2019b); Zhang and Prange (2020)). The NorESM - with a reasonably simulated AMOC and Arctic sea ice distribution in the PI and historical simulations (as documented by Guo et al. (2019a)) - simulates a MIS3 climate that is in a stable regime with relatively strong convections in the Norwegian and Labrador seas. Indeed, NorESM sensitivity experiments including large reductions in atmospheric CO₂ levels and Laurentide Ice Sheet heights, aimed at perturbing the system into a cold stadial-like climate, indicate that the model state appears to be far from a possible threshold (Guo et al., 2019b). Zhang and Prange (2020) use the LGM ICE-5G ice sheet configuration (Peltier, 2004), with a high Laurentide Ice Sheet (at just over 4000 m) which may have contributed to a strong AMOC in the CCSM3 simulation, alongside its particular background climate

135 background climate.

In summary, IPCC-class models set up with pre-industrial or present-day conditions do not exhibit D-O type warming events, but can feature shorter centennial length cooling and warming events. This model behaviour is consistent with observations, since millennial timescale D-O events do not occur under interglacial conditions but periods of centennial-scale AMOC variability are present throughout several interglacials (Galaasen et al., 2014; Tzedakis et al., 2018; Galaasen et al., 2020). Some

140 models which are set up with more MIS3 like conditions exhibit D-O type warming events, but some do not. Under full LGM conditions only one model (UoT-CCSM4) out of ten (PMIP4 LGM simulations: Kageyama et al. (2021a)) show spontaneous D-O type oscillations (Tables A1 and B1; Kageyama et al., 2021a; Peltier and Vettoretti, 2014).

Since it can take some time for D-O type oscillations to evolve, it is unclear if some models would develop such oscillations if they were run for longer (at least for 2000 model years). Of the forty LGM/MIS3-like simulations (Table B1; Kageyama

- et al., 2021a; Armstrong et al., 2021; Klockmann et al., 2018), sixteen simulations have been run for less than 2000 years (Table B1), which makes it difficult to tell whether any of these simulations are capable of, or likely to, exhibit D-O like behaviour under specific boundary conditions. In addition the duration of LGM/MIS3 simulations is currently inadequate, we note that the majority of CMIP6 models appear not to have performed any form of glacial period simulation (Table C1). Thus, it is difficult to ascertain what proportion, or indeed which, models are capable of capturing D-O like behaviour, under any form of
- 150 glacial period state (Table C1).

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2.1 The role of ocean and sea-ice feedbacks

Changes in the AMOC are crucial to the correct simulation of D-O events (Broecker and Peteet, 1985). The AMOC features stabilising positive feedbacks: a strong AMOC transports warm and salty water into the subpolar North Atlantic, thus weakening the stratification and also keeping the sea ice cover reduced (*e.g.* Rahmstorf, 2002; Clark et al., 2002). As a consequence, there is a large transport of heat northward across the hemispheres (*e.g.* Feulner et al., 2013; Buckley and Marshall, 2016),

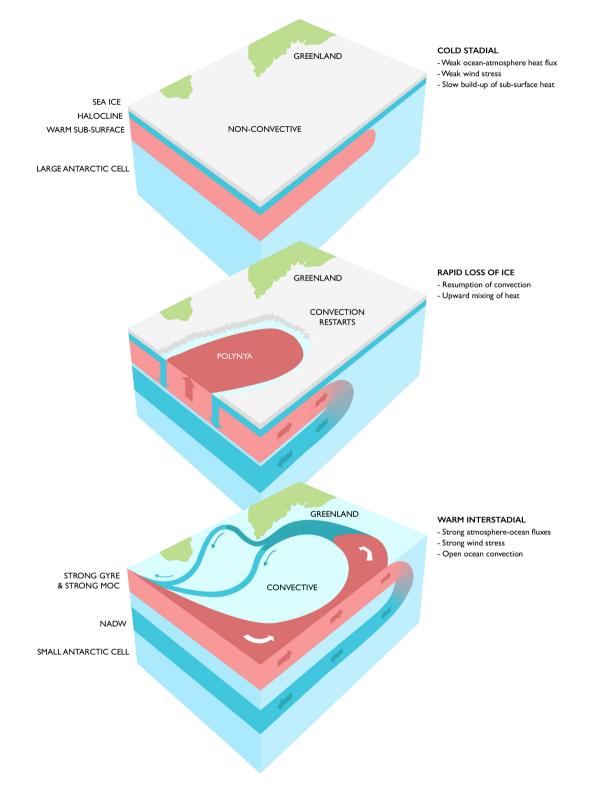


Figure 2. Schematic depicting the transition from GS to GI conditions *i.e.* a D-O warming event.

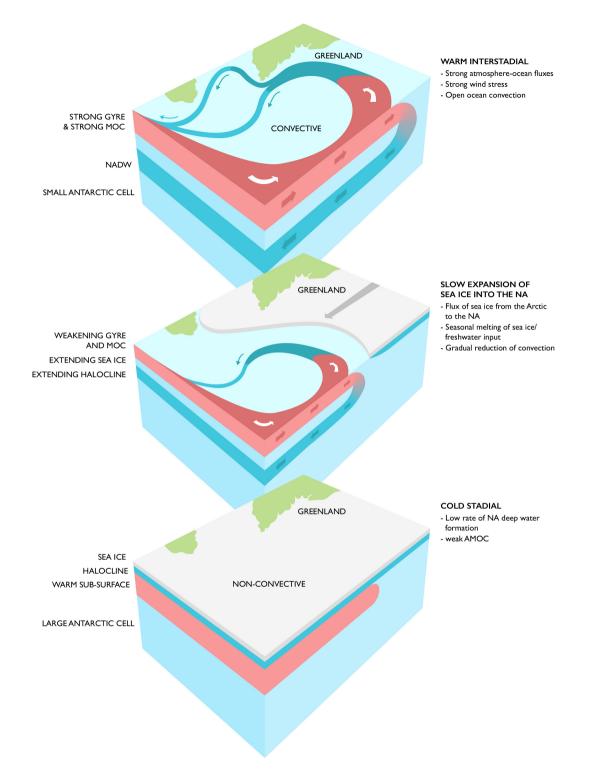


Figure 3. Schematic depicting the transition from GI to GS conditions *i.e.* a D-O cooling event.

strong heat loss in the North Atlantic and Arctic, and active deep convection that sustains the strong AMOC. A weak AMOC, on the other hand, is associated with a weaker northward transport of salt and heat. This increases the stratification in the subpolar North Atlantic and thus favors the expansion of sea ice. The weak northward heat transport and the insulating effect of the sea ice keep the density gain due to heat loss small and the AMOC in a weak state (*e.g.* Klockmann et al., 2018). This

weak AMOC state is stable when Antarctic Bottom Water becomes dense and salty enough to replace North Atlantic Deep

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Water (NADW) in the deep North Atlantic.

Sea ice can act as both a slow and fast positive feedback on AMOC-induced changes in climate. Extensive stadial sea ice cover during a weak AMOC state cools Greenland and suppresses atmosphere-ocean exchange of heat and oceanic convection in the North Atlantic (Li et al., 2005, 2010). This also leads to a slow build up of heat in the North Atlantic subsurface.

165 Foraminifera from marine sediment cores offer evidence to back-up that this sub-surface warming occurred before the onset of fast D-O warming events (Rasmussen and Thomsen, 2004; Singh et al., 2014; Dokken et al., 2013). This heat build-up sets up the conditions for subsequent fast losses of GS sea ice.

Wind-driven, AMOC, and sea-ice linked salinity changes also play a crucial role in D-O positive and negative feedbacks. Indeed the net freshwater transport in the Atlantic basin by the AMOC can be used to assess the stability regime of the AMOC

- 170 (Rahmstorf, 1995; Huisman et al., 2010). The interaction of subpolar and tropical salinity anomalies at the surface and in the subsurface (Jackson and Vellinga, 2013), and possible roles of the intertropical convergence zone and freshwater export through the Fram Strait, are also important in D-O related salinity feedbacks. Klockmann et al. (2018) note that if the subtropical gyre shifts northward and the sub polar gyre contracts, an inflow of salty subtropical water extends over the entire Atlantic basin east of the Mid-Atlantic Ridge. This inflow can supply salty water to the deep-convection sites in the Iceland Basin and Irminger
- 175 Sea, and help maintain continuous deep convection and a strong AMOC even at low CO₂ concentrations (Brown and Galbraith, 2016; Klockmann et al., 2018; Guo et al., 2019b; Muglia and Schmittner, 2015; Sherriff-Tadano et al., 2018), thus preventing the initiation of GS-like conditions. Where the AMOC does enter a weak state for a prolonged period, and the climate enters a GS, a build-up of heat in subsurface waters and salt in the tropical Atlantic can enable the very rapid resumption of the AMOC (Lynch-Stieglitz, 2017), with the upward mixing of heat from the subsurface and importation of salt from the tropic Atlantic
- 180 via gyre mechanisms (Peltier and Vettoretti, 2014).

The importance of vertical (diapycnal) mixing in the ocean for these long timescale, D-O type, instabilities has long been recognised (Welander, 1982). However, we note that the different ocean- and climate-models (Table A1) parameterise diapycnal mixing in very different ways (*e.g.* Nilsson et al., 2003; de Lavergne et al., 2019). The lack of a single consistent paramaterisation and differences in the strength of diapycnal mixing across climate models means it is to be expected that

- 185 some models will produce D-O like oscillations "out of the box" under MIS3 boundary conditions, but others may require changes, or tuning, to their diapycnal mixing parameters. Even within the same model, a large range of diapycnal diffusivities may yield steady states that satisfy common plausibility constraints such as AMOC transport and sea ice distribution (see e.g. Holden et al., 2010). This is partly because wind-driven Southern Ocean upwelling plays a complementary role to diapycnal mixing in setting the steady state overturning (Samelson, 2004), and partly because surface buoyancy forcing controls the rel-
- 190 ative strength of the upper (Atlantic) and lower (Antarctic) overturning cells (Oliver and Edwards, 2008). A realistic AMOC

transport may be obtained due to compensating biases in these processes, which has serious implications for whether AMOC feedbacks (necessary for capturing D-O behaviour) are represented in an adequate manner within these models.

Figure 2 and 3 show some of the key states, processes, and ocean sea-ice feedbacks that enable D-O events. Following Lohmann and Ditlevsen (2019), D-O events can be broken down into four periods: (1) cold stadial state (Fig. 2a), (2) rapid

- 195 warming phase governed by very fast-time-scale mechanisms (Fig. 2b), (3) warm interstadial state (Fig. 2c and Fig. 3a) and, (4) gradual cooling phase (Fig. 3b) followed by a faster abrupt transition into a cold stadial phase (Fig. 3c). For some of the D-O events, the magnitude of the warming transitions are on the order of ten degrees in a decade, while the slow cooling in the interstadials is on the order of a few degrees in a millennium (the sawtooth shape) (Lohmann and Ditlevsen, 2019). This picture of rapid retreat of North Atlantic sea ice (Spolaor et al., 2016; Dokken et al., 2013) associated with the resumption of
 - 200 convection and the AMOC, alongside an upwards mixing of salt and heat, followed by a slower cooling phase back into stadial conditions matches accumulation, temperature, and water isotopes retrieved from Greenland ice core records of D-O warming events (Li et al., 2005, 2010; Sime et al., 2019).

2.2 The role of Northern Hemisphere Ice Sheets

Section 2 and Table A1 suggest that large Northern Hemisphere Ice Sheets and the wind regime associated with these can
contribute to a strong AMOC which stabilises the North Atlantic and prevents D-O events. Thus ice sheets have a critical role
to play in setting up the conditions for D-O events (Zhang et al., 2014; Klockmann et al., 2018; Brown and Galbraith, 2016;
Muglia and Schmittner, 2015; Sherriff-Tadano et al., 2018). Figure 4 and 5 show some of the key mechanisms and feedbacks
that are behind a state of reduced likelihood for D-O events and a potentially D-O type oscillating state, respectively.

210 The Northern Hemisphere Eurasian ice sheet was most probably limited to mountainous areas during mid-MIS3 (Helmens, 2014; Hughes et al., 2016), and its impact on D-O dynamics was probably relatively small. However, the size and presence (or absence) of the Laurentide ice sheet (LIS), which has elevations reaching a maximum of approximately 3000 m (Abe-Ouchi et al., 2015) at the LGM, does appear to cause important and robust (across multiple models) changes to Northern Hemisphere atmospheric circulation and resultant wind forcing of the ocean. LIS-dependent wind changes influence the subpolar gyre and the stability of the atmosphere-ice-ocean coupled system (Li and Born, 2019; Zhang et al., 2014).

A larger LIS (especially its height) causes stronger Northern Hemisphere winds (Li and Battisti, 2008; Pausata et al., 2011; Hofer et al., 2012; Ullman et al., 2014; Löfverström et al., 2014; Merz et al., 2015); an amplified stationary wave over North America (Manabe and Broccoli, 1985; Cook and Held, 1988); the North Atlantic glacial jet to be more stable due to differences in wave-mean flow feedbacks (Riviere et al., 2010); and alters variability of the large-scale atmospheric circulation, especially in the North Atlantic (Justino and Peltier, 2005; Pausata et al., 2009; Riviere et al., 2010). In addition, LIS height could control

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the sea-ice coverage and gyre circulation by shifting the westerlies over the North Atlantic region Zhang et al. (2014). LIS altered winds that have wide implications for D-O relevant tipping elements (Seager and Battisti, 2007; Wunsch, 2006).

Lis and Born (2019) note that, first, the presence of a large LGM-type LIS is linked to a strong, more zonal and equatorwardshifted North Atlantic jet which weakens atmospheric heat transport into the North Atlantic (van der Schrier et al., 2010)

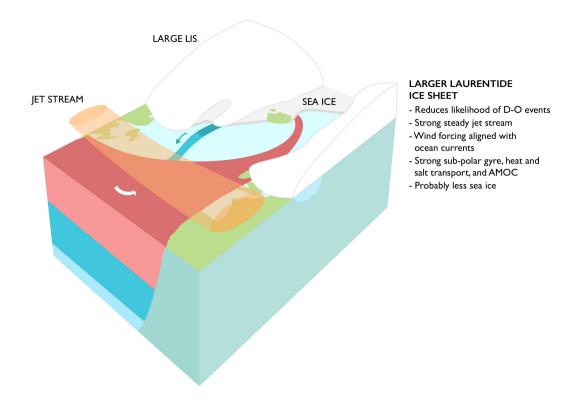


Figure 4. Schematic showing a state of reduced likelihood for D-O events.

- and favours episodes of Greenland blocking (Madonna et al., 2017). Both could trigger the atmosphere-ice-ocean feedbacks that cause abrupt climate change in this area. Second, a steadier and stronger North Atlantic jet strengthens the wind-driven component of the subpolar gyre (Li and Born, 2019). Given that at latitudes north of about 45N, the subpolar gyre, which is essentially wind-driven, plays a crucial role in the northward transport of heat and salt, and is strongly linked to the AMOC (*e.g.* Jungclaus et al., 2013), wind-driven changes in this gyre have a strong impact on the density gain in the North Atlantic.
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In many simulations with a large LIS (LGM-like ice sheets), the subtropical gyre can shift northward and cause an inflow of salty subtropical water over deep-convection sites, contributing to continuous deep convection and a strong AMOC even at low CO₂ concentrations (Brown and Galbraith, 2016; Klockmann et al., 2018; Guo et al., 2019b; Muglia and Schmittner, 2015; Sherriff-Tadano et al., 2018; Zhang et al., 2014). Similarly, Zhang et al. (2014) note that a higher LIS can promote less South Labrador Sea sea ice export to northeastern North Atlantic (which reduces sea ice concentration) to permit deep convection and shift the core of westerlies northwards, strengthening subtropical gyre for heat and salt transport (Zhang et al., 2014).

For these reasons, large LGM-type ice sheets, particular a large LIS, tend to lead to a density gain over the North Atlantic and the northward salt transport is enhanced with respect to the PI ice sheet case. For many, but not all models, this tends to lead to more active convection in the North Atlantic and a strong AMOC (across a wide range of CO_2 concentrations). That said, the AMOC in many LGM simulations is likely too strong (Klockmann et al., 2018; Kageyama et al., 2021b). Thus the

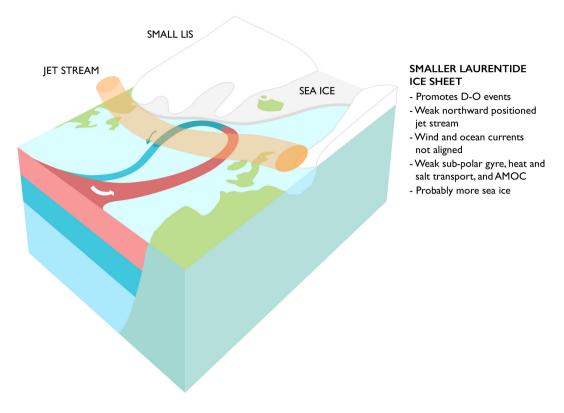


Figure 5. Schematic of a potentially D-O type oscillating state.

AMOC is far away from a tipping point with LGM-size ice sheets for many models (Zhang et al., 2014; Klockmann et al., 2018; Guo et al., 2019b).

In some simulations with reduced ice sheets, the jet stream shifts northwards, leading to regional cooling and a rise in seasonal sea ice concentration over the subpolar gyre region (Armstrong et al., 2021). This freshens the area and lowers deep-water formation, which weakens the subpolar gyre and as a result the simulations are more prone to enter a weak convection, weak AMOC mode which is conducive to D-O type oscillations (Klockmann et al., 2018; Armstrong et al., 2021). Thus, with intermediate MIS3 LIS, *i.e.* reduced in its height compared to the LGM, multiple AMOC states are more likely (Zhang et al., 2014; Kawamura et al., 2017; Zhang and Prange, 2020; Armstrong et al., 2021; Klockmann et al., 2018).

3 Contours of a baseline MIS3 experiment protocol

Although the choice of a time within MIS3 for a D-O baseline experiment should be unimportant, given that in reality D-O events occurred during the whole of the MIS3, our analysis of existing simulations, boundary conditions and mechanisms above suggests that there are periods which may be particularly conducive to D-O events occurring in models. Oscillatory D-O type behaviour appears to be more likely, but not guaranteed (Guo et al., 2019b; Zhang and Prange, 2020), when models are run with intermediate or low MIS3 CO_2 values and ice-sheets, *i.e.* reduced in size compared to the LGM (Brown and Galbraith, 2016; Kawamura et al., 2017; Klockmann et al., 2018; Zhang and Prange, 2020; Galbraith and de Lavergne, 2019; Zhang et al., 2014; Vettoretti et al., 2022), and particularly without a high LIS.

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The impact of orbital parameters has been investigated in less detail than the role of GHGs and ice sheets. Using the model COSMOS, Zhang et al. (2021) demonstrated that under intermediate glacial conditions, obliquity appears to play a significant role in the occurrence of D-O type behaviour. In particular, the orbital parameters at 40 ka do not produce D-O type behaviour, whilst at 34 ka lower obliquity (22.6°) leads to D-O type behaviour (see Figure 2 from Zhang et al. (2021)). Additionally,

the MIROC4m model produces D-O-type oscillations (under mid-glacial conditions) and low obliquity (22.9°) (Kuniyoshi et al., 2022). From these COSMOS and MIROC4m results, we deduce that low obliquity seems conducive to D-O behaviour in models.

These considerations suggest that the interval starting at 38 ka to 32 ka is a good choice for the proposed baseline experiment: it is characterised by (1) a rather regular sequence of D-O events (Fig. 1), and (2) has the ideal intermediate MIS3 ice-sheet configuration conducive to generating D-O-type quasi-oscillations (Section 2).

A baseline simulation needs to be run for a sufficient duration to allow the strong positive feedbacks, together with long time-scale negative feedbacks, that enable D-O type oscillations. The analysis of existing simulations (Section 2) suggests this should be a minimum of 5000 years (Peltier and Vettoretti, 2014; Kleppin et al., 2015; Sidorenko et al., 2015; Brown and Galbraith, 2016; Klockmann et al., 2018, 2020). However, given computational constraints, a minimum duration of around 2000 years, with a spin-up period of 1000 years, may be a more practical minimum requirement for most modeling groups.

It would, however, be important to examine and document key metrics for model drift (such as top-of-atmosphere radiation imbalance, deep ocean or global mean ocean temperature) during the initial spin-up. The exact length of spin-up is thus subject to discretion of each modelling group based on these key metrics.

There are two obvious possibilities for spinning up the MIS3 control experiment (MIS3-cnt). The baseline experiment could

- 275 be initialised from either the end of a well spun-up LGM or PI experiment. Other possibilities could be to spin up from a linear combination of LGM and PI states (as done in Klockmann et al., 2016, 2018) or spinning up from present day's observations (as done in Guo et al., 2019b). Modelling groups are encouraged to choose whichever option is more feasible/convenient for them. In case that several spin-up options are available, short spin-ups with diagnosed top-of-atmosphere (TOA) imbalance or global mean ocean temperature could help distinguish the faster spin-up option. It is worth noting that initial ocean state (i.e.
- 280 Atlantic salinity stratification) does play a role in abrupt AMOC change and associated feedbacks (Zhang et al., 2013; Knorr et al., 2021), of which impacts shall be considered and evaluated in the future.

We suggest performing a MIS3-cnt experiment centered at 34 ka, using GHG and orbital conditions for 34 ka (Fig. 6); and ice sheet configuration as outlined below (sections 3.1 and 3.2).

We acknowledge that some models might not oscillate under the proposed 34 ka baseline scenario. Indeed, this is expected for NorESM, which under 38 ka conditions, is in a stable regime and the model state seems to be far from a possible tipping point. In spite of that, standardised MIS3 simulations which do not show D-O like behaviour are still highly valuable, for exactly the same reasons that LGM simulations are relevant to the wider modelling community. These standardised MIS3

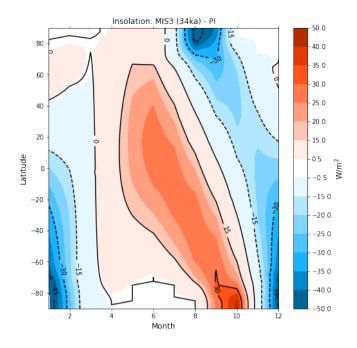


Figure 6. Monthly zonal-mean MIS3 (34ka) - PI anomalies of the top-of-atmosphere short-wave incoming radiation (W m⁻²).

simulations could contribute to progress on the overarching CMIP6 questions 1 and 2 (Eyring et al., 2016): "How does the Earth System respond to forcing?", and "What are the origins and consequences of systematic model biases?" With a larger number of standardised MIS3 simulations, we would be able to answer questions such as:

- Are state-of-the-art climate models capable of representing D-O events under more realistic MIS3 conditions? Benchmarking these simulations will deliver a measure of how well models simulate abrupt changes, and tipping events.
- Standardized MIS3 simulations can help explore the existence of a theoretical sweet spot for millennial activity in current climate models (Barker and Knorr, 2021). As close to or within the sweet spot, the AMOC is characterized by high sensitivity to transient and/or noisy climatic forcing (Zhang et al., 2014; Lohmann and Ditlevsen, 2018) or by self-oscillating behaviors (Zhang et al., 2021).
- If models are too stable to simulate abrupt transitions, what are the processes that contribute to relative levels of model stability?

 In addition, a larger number of standardised MIS3 simulations could encourage the creation of new data sets, improving model-data evaluation.

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3.1 Atmospheric trace gases

MIS3 atmospheric CO₂ values varied between a maximum of ~ 233 ppm to a minimum of ~ 187.5 ppm (Table 2; Figure 1; Bauska et al., 2021). Interestingly, increases of around 5 ppm happened during the abrupt warming of most D-O events and increases of up to 10 ppm happened within some Heinrich stadials (Bauska et al., 2021). GHG forcing is critical to model stability. Low (LGM-like) to intermediate (MIS3) CO₂ concentrations tend to be associated with abrupt D-O type AMOC transitions in models (Section 2 and Klockmann et al., 2018; Zhang et al., 2017, 2014; Brown and Galbraith, 2016; Vettoretti et al., 2022). We thus suggest to perform the MIS3-cnt experiment using the GHGs values specified in Table 2 and keep these

3.2 **Northern Hemisphere Ice Sheets**

values fixed for the whole duration of the simulation including the spin-up.

- 310 Constraining MIS3 ice-sheet boundary conditions is a challenge. Scarcity and fragmentation of evidence (Kleman et al., 2010; Batchelor et al., 2019) is an issue. In particular, it is difficult to determine the size and shape of the ice sheets during MIS3 because subsequent larger LGM configurations have overridden and destroyed evidence of the position of the margins of these smaller ice sheets.
- Global sea level fluctuations during the mid-MIS3 were driven nearly exclusively by the LIS (Gowan et al., 2021). Global 315 average sea level remained above -55 m for the period between 30-55 ka. From glacial isostatic modelling and geological constraints, a global mean sea level between -30 m and -50 m is inferred (Dalton et al., 2022). For much of MIS3, since the Eurasian ice sheets and the Cordilleran Ice Sheet were likely restricted to mountain-based caps (Helmens, 2014; Hughes et al., 2016; Clague and Ward, 2011), the primary control on ice volume is assumed to be from the LIS. Recent work in the area of the Hudson Bay (Dalton et al., 2016, 2019; McMartin et al., 2019; Dalton et al., 2022) suggests ice-free conditions may have
- 320 occurred during mid-MIS3. This implies climatic conditions in this region similar to present (Dalton et al., 2017), and a LIS margin removed from the southern Hudson Bay. Similarly Tarasov et al. (2012) show a considerably lower and less extensive LIS compared to ICE-5G and ICE-6G (Peltier, 2004; Peltier et al., 2015) LGM ice sheet reconstructions. Pico et al. (2017) sea-level curves are consistent with the estimated MIS3 ice-sheet volumes from Batchelor et al. (2019). Using Glacio Isostatic Adjustment modeling, Pico et al. (2017) also show that a small LIS can explain high MIS3 sea-level estimates alongside the
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eastern coast of the United States. Batchelor et al. (2019)'s synthesis of numerical modelling results and empirical data provides additional support for a considerable reduction in the MIS3 LIS extent and very minimal European ice sheet.

The recent MIS3 ice sheet reconstruction, PaleoMIST 1.0 (Paleo Margins, Ice Sheets, and Topography), was developed independently of far-field sea-level records and indirect proxy records by Gowan et al. (2021). This reconstruction is based on trying to fit the evolution of ice flow indicators, as well as chronological constraints of ice-free conditions.

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Gowan et al. (2021) provide a maximum and minimal MIS3 reconstruction, specifically for the Laurentide Ice Sheet. The maximum scenario is more consistent with recently discovered eastward oriented, pre-LGM ice flow direction indicators found in southeastern Manitoba (Gauthier and Hodder, 2020), so we currently consider it to be more likely. However, at 35 ka, the difference between the two scenarios is minor. The difference is primarily with the thickness (and therefore also topography)

of the ice sheet, rather than extent, but it amounts to less than 1 m of sea level equivalent. The 35 ka time slice represents conditions after Heinrich Event 4 (Andrews and Voelker, 2018), and the ice margin in Hudson Strait is retreated about 350 km from the edge of the continental shelf. The ice margin elsewhere for the Laurentide Ice Sheet is based on chronological constraints, most that are documented in the compilation by Dalton et al. (2019). The Cordillera Ice Sheet extent is based on evidence of relatively restricted ice cover during MIS 3 (Clague and Ward, 2011). The Greenland Ice Sheet margin is set to be

intermediate of the LGM and present day extent. The Eurasian ice extent at 35 ka includes an advance of ice into the Baltic

340 Sea, which happened after Heinrich Event 4 (Hughes et al., 2016). For East Antarctica, the margin is set to be the same as present. In West Antarctica, the margin at 35 ka is close to the shelf edge, as the maximum extent may have been achieved by 30 ka (Larter et al., 2014).

Given its strong evidence basis, we thus suggest the use of the maximum 35 ka Gowan et al. (2021) PaleoMIST ice sheet configuration. We note the LIS is considerably reduced in size, compared to the ICE-6G LGM reconstruction in the southeastern margin (Fig. 7a,d); the EIS is also significantly smaller (Fig. 7a,d).

Whilst the implementation of the ice sheet will differ between models, the steps of Kageyama et al. (2017) describe how to implement a glacial state ice sheet in the IPSL climate model. For consistency, we likewise recommend the same steps should be followed as far as possible. Since a reduced sea-level can modify river courses, Kageyama et al. (2017) recommend that as a minimum, rivers should reach the oceans. Also, the ocean should be initialized with a salinity 0.6 psu higher than the PI experiment, to account for the sea-level difference between MIS3 and PI experiment (freshwater stored as ice on land) (Guo et al., 2019b).

The single ice sheet reconstruction MIS3 set-up summarized above contrasts with the PMIP4 LGM protocol, which provides three different possible ice sheet configurations (PMIP3, ICE6G-C and GLAC-1D) for the tier 1 LGM experiment (Kageyama et al., 2017). This partially reflects the more limited knowledge of ice sheet in pre-LGM periods. Exploration of the effect of MIS3 ice sheet reconstructions uncertainties on climate models, particularly on model stability, would be valuable. For this purpose, further additional 34 ky / MIS3 ice sheet reconstructions would be very valuable.

3.3 Heinrich-Event preconditioned option

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The term "kick" Heinrich event-like was initially introduced by Peltier and Vettoretti (2014) to invoke a 'kicked' salt oscillator hypothesis (pseudo-Heinrich type behaviour), to induce D-O type oscillations in an LGM simulation performed with the UofT

360 CCSM4. During the first thousand years of the simulation as the model is spun up and the ocean cools to reach a state consistent with glacial boundary conditions, there are two thermal thresholds during which the strength of the AMOC rapidly reduces (see Figure 2 in Peltier et al., 2020). These abrupt transitions in the AMOC coincide with abrupt reductions in surface temperatures in the North Atlantic and abrupt expansions of sea ice coverage. During the second of these events, the AMOC is reduced to approximately 12 Sv, about half its strength in the pre-industrial control (Peltier et al., 2020). This event may resemble the impact of a Heinrich event-like "kick" to the AMOC though no freshwater perturbation was imposed (Peltier et al., 2020).

In a more recent study, Pedro et al. (2022) examine the CCSM4 simulations that shows unforced D-O type oscillations (Peltier and Vettoretti, 2014; Vettoretti and Peltier, 2018), but with the addition of a (freshwater) H-like event. (Pedro et al.,

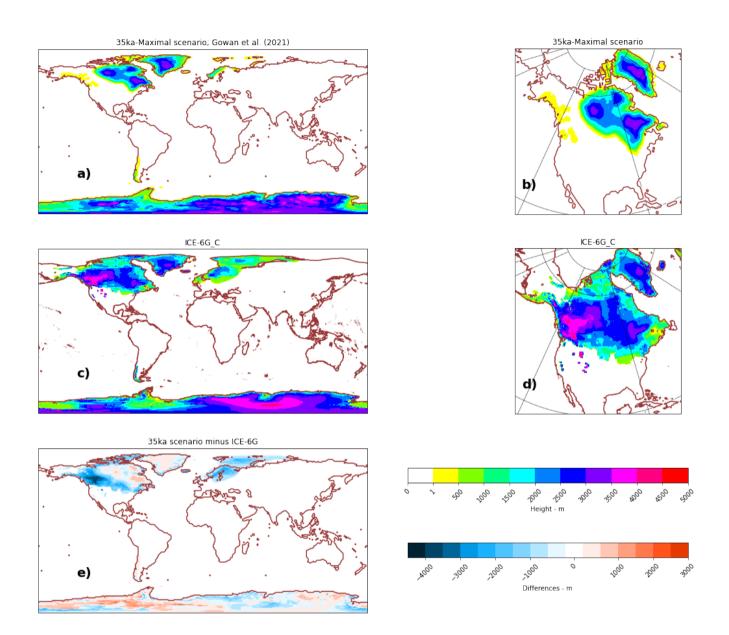


Figure 7. MIS3 ice sheet reconstruction from (a,b) Gowan et al. (2021). Also shown for comparison is the (c,d) LGM ICE-6G ice sheet reconstruction from Peltier et al. (2015). Third row shows the differences between the MIS3 ice sheet reconstruction and the LGM ICE-6G reconstruction.

2022). The freshwater flux (0.05 Sv injected into the NA for 500 years) leads to (1) a 5 Sv weaker AMOC compared to the one seen in the unforced model stadial and, (2) a stronger D-O warming transition into the interstadial phase (Pedro et al., 2022).

370 Thus this type of H-E preconditioning can trigger abrupt reductions in the AMOC strength and in NA surface temperatures and sea ice coverage – and it may also help induce a stadial state in other models which is more conducive to unforced (D-O type) oscillations (Pedro et al., 2022).

Given the importance of HEs to starting Bond Cycles of D-O events, an additional experiment to investigate how HE meltwater preconditioning impacts the simulation of D-O like oscillations under MIS3 boundary conditions would be valuable.

375 HE freshwater preconditioning may, as in reality, be more conducive to a (Bon Cycle-like) sequence of spontaneous D-O type oscillations (see Table 1).

The freshwater delivered during Heinrich event iceberg discharge suppresses the AMOC, leading to accumulation of heat in the Southern Hemisphere, and in the North Atlantic subsurface waters (e.g. Stocker and Johnsen, 2003). Estimates of the meltwater input into the North Atlantic during Heinrich events range between 2 m and 15 m of sea level equivalent ice volume

- 380 (Hemming, 2004; Chappell, 2002; Rohling et al., 2004; Roche et al., 2004; Roberts et al., 2014b; Siddall et al., 2008; Grant et al., 2014). It is logical to presume that these freshwater events are important in preconditioning the climate system with respect to D-O behaviour (Peltier and Vettoretti, 2014; Peltier et al., 2020; Pedro et al., 2022). Freshwater perturbations can trigger changes between AMOC states (e.g. Ganopolski and Rahmstorf, 2001; Timmermann et al., 2003; Stouffer et al., 2006; Hu et al., 2008; Kageyama et al., 2012; Roberts et al., 2014a; Sime et al., 2019) and a relatively small freshwater flux applied
- 385 over convection areas can lead to a shutdown of the AMOC (e.g. Roche et al., 2010). Studies have shown similarities between observed global features of abrupt D-O changes and the behaviour seen in freshwater forcing experiments (Liu et al., 2009; Menviel et al., 2014). However, the sensitivity of the AMOC to a wide range of freshwater inputs varies according to model, where the meltwater is added, and the background climate state (Ganopolski and Rahmstorf, 2001; Timmermann et al., 2003; Stouffer et al., 2006; Hu et al., 2008; Kageyama et al., 2012; Roberts et al., 2014a; Zhang et al., 2014). Given these various uncertainties, we suggest that it would be useful to run an additional experiment to investigate how preconditioning through a

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H-like (freshwater) event impacts the simulation of D-O like oscillations under MIS3 boundary conditions.

The range of H-event volumes calculated using ice sheet models varies from 24.2 to 125×10^4 km³ (MacAyeal, 1993; Dowdeswell et al., 1995; Marshall and Clarke, 1997; Hulbe, 1997), whilst isotope based estimates and a precipitation balance approach have yielded 86, 649, and 946 x 10^4 km³ of ice volume (Hemming, 2004; Roche et al., 2004; Levine and Bigg, 2008).

- Roberts et al. (2014b), and using a sediment modelling approach estimated a discharge of 30 to $120 \times 10^4 \text{km}^3$ of ice volume. 395 Some of the spread in these estimates could be because the relationship between the oxygen isotope record, sea level, and meltwater volume is not constant when ice is lost from marine basins, such that the use of oxygen isotopes for calculating H-event volumes may produce unrealistically high values (Gasson et al., 2016; Hemming, 2004; Roberts et al., 2014b). There is also some uncertainty about the duration of the H-events, with some previous studies suggesting they could be as short as 250
- 400 years (Hemming, 2004) and others suggesting a duration of 500 yr is more typical (Roberts et al., 2014b). These considerations suggest that it is possible to justify the use of anywhere between 0.02 - 0.6 Sy freshwater flux over 500 years; or 0.04 - 1.2 Sy over 250 years. More recent estimates of H-event magnitudes tend to favour the lower end of this range. If all forcings are set

to MIS3-cnt values and the H-event freshwater flux is distributed across the North Atlantic this could yield a range of stadial climates (Sime et al., 2019; Zhang and Prange, 2020). After 250-500 years this freshwater forcing should be switched off.

Table 2. Summary of the boundary conditions (BC) and forcings for the MIS3-cnt experiment.

| CO_2 : 208 ppm (Bauska et al., 2021) CH_4 : 420 ppb (Loulergue et al., 2008) N_2O : 204 ppb (Schilt et al., 2010) Eccentricity: 0.01567Berger et al. (1998) | | |
|--|--|--|
| N_2O : 204 ppb (Schilt et al., 2010) | | |
| | | |
| Eccentricity: 0.01567Berger et al. (1998) | | |
| | | |
| Obliquity: 22.6° Berger et al. (1998) | | |
| Precession: -0.016 Berger et al. (1998) | | |
| Same as PI control | | |
| 35ka ice sheet reconstruction (Gowan et al., 2021); | | |
| mean global salinity increased by 0.6PSU to account for ice volume | | |
| Closed to avoid drifts; Snow should not accumulate | | |
| over ice sheets and rivers should flow into the ocean. | | |
| Models need to consider lakes when closing | | |
| the global freshwater budget | | |
| Dynamic or fixed as in PI. | | |
| If fixed vegetation: tundra in new land points | | |
| As in PI control | | |
| initial 0.04 - 1 Sv over 250-500 years | | |
| followed by standard MIS3-cnt simulation | | |
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405 4 Conclusions

D-O events are abrupt, large climate changes that punctuated the last glacial period. There is uncertainty whether current IPCC-class models can effectively represent the processes that cause D-O events. We have shown that reduced ice sheets relative to LGM, low obliquity values and low-to-medium MIS3 CO₂ values are more likely to lead to unforced quasi-oscillatory D-O type behaviour. However, the simulations need to be run long enough to allow the strong positive AMOC feedbacks, along with negative feedbacks on long time-scales, which can then lead to D-O type oscillations. Around 40% of the simulations set-up with full LGM or more MIS3-like conditions, have a run length of less than 2000 model years, which makes it difficult to tell whether any of these simulations are capable of, or likely to, exhibit D-O like behaviour. In addition, the vast majority of PMIP4/CMIP6 models have not run LGM or MIS3-like simulations long enough to be sure which models have the capability to oscillate.

- We have provided boundary conditions for a baseline MIS3-cnt simulation, and a H-event preconditioned variant (freshwater forced experiment). The MIS3-cnt experiment is centered at 34 ka because it yields the ideal combination of intermediate ice sheets (smaller in size compared to LGM), low obliquity values and medium-to-low GHG values conducive to oscillatory D-O type behaviour in models. Ideally, the MIS3 baseline experiment should be run for 5000 years, however, given computational constraints a minimum duration of 2000 years together with a spin-up of at least 1000 years is a more practical minimum re-
- 420 quirement. This baseline MIS3-cnt protocol provides a common framework to explore cold-period instabilities using particular GHG-, insolation-, freshwater-, and NH ice sheet-related forcings, together with diapycnal mixing. More model simulations run under the here proposed MIS3 DO protocol together with analyses across models, could provide better insights, along the lines of atmospheric-ice-ocean feedbacks behind DO events. These simulations will allow us to answer questions such as: are current climate models able to reproduce DO-type behaviour under more realistic MIS3 conditions? How well models simu-
- 425 late tipping events, abrupt changes? What are the mechanisms that lead to relative levels of model stability? Moreover, a large number of standardised MIS3 simulations could encourage the creation of new data sets, improving model-data evaluation.

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845 Appendix A

| oscillations. |
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| Models |
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| Table |

| Study | Model | Doriod | CHC | Ire | Incolation | EWF | Bun | Main findinge | Machanisms for D-O cooling acout | Mechanisms for D-O warming avant | Similar Machaniem to |
|--|-------------------|----------------------------------|-----------------------------------|----------------------------------|------------|------|--------------------|--|--|--|--|
| | | | | sheet | | | Length | 0 | 0 | D | |
| Drijfhout et al. (2013) | BC-Earth model | Ы | Ы | Id | Ы | None | 1125 yrs | Spontaneous cold event that lasts around 100 years | An anomalous high atmospheric block- ing over the castern subpolar gyre causes the cold event. Ocean currents transport east ice southwards and there is a shut- sen ice southwards and there is a shut- down of deep water convection in the Labrador Sea. | No warning event reported. | No. Atmospheric blocking-sea-ice-ocean feedback identified as a main cause behind the cold event. |
| Kleppin et al. (2015) | CCSM4 | ы | ı L | E. | Ы | None | 1000 yrs | Spontaneous cold event that lasts around 200 years | The cooling event has a duration of 200 years and is linked to a weakened state of the SPG and deep water convection in the Labrador Sea. | The warming event is triggered by a stronger leelandic low and therefore deep water convection recovers and SPG circulation resumes. | No. Stochastic atmospheric forcing identified as a poten- tial cause for sea ice varia- tions. |
| Sidorenko et al. (2015) | ECHAM6- FESOM | Present day (PD) | 0691 - CI4 | G | PD - 1990 | None | 350 yrs | Events of sudden reduc- tion of deep water con- vection and increase of sea ice cover in the Labrador Sea. | An anomalous inflow of warm and suline water into the deep Labrador Sea causes a weakening of the subpolar gyre and modifies the upper freshwater budget over the Labrador Sea. | No warming event reported. | No. The strong surface winds over the subtropical NA al- ter the Ginzhata Strait outflow path into the NA and are iden- pited as the main cause behind the stilne and warm anomalies in the deep NA. |
| Martin et al. (2015) | КСМ | Q | G | Q | Q | None | 1300 yrs | Centennial-scale vari- ability of the AMOC as well as variations in the NA heat content and subpolar gyre strength. | As the NA Current accelerates, deep convection in the Weddell Sca enables a positive heat content anomaly to prop- agate northwards in the upper Athunic Ocean. Eventually, the heat anomaly creaches the northern NA resulting in a reduction in deep water formation there. | The retreat of Antarctic Bottom Wa- ter (AABW) leads to enhanced merid- ional density gradient that results in an increased North Atlantic Deep Water (NADW) cell. | No. Interhentispheric telecon- nection; Variability in the Southern Ocean deep water convection identified as the main cause behind AMOC os- cillations. |
| Peltier and Vet- toretti (2014); Vettoretti and Peltier (2016, 2018) | UofT CCSN4 | ILGM | LGM – 21ka | LGM - ICE- 6G (VM5a) | LGM – 21ka | None | 5000 yrs | Spontaneous millermial- scale D-O type oscilla- tions | The continuous flux of sea ice from the Arctic basin into the NA subpolar gyre area across the East Greenhand Current, favours melting of sea ice as it moves over the warm ocean surface. This fresh- water input restartifies the high-latitude NA and results in a considerable de- crease in the rate of NA Deep Ware for- mation. | The initiation of the abrupt warming events is associated with the opening of a large polynya over the Irminger Sea. The stability of the water column is key and depends on transport of saft to the suppolar gyre along the Irminger Cur- rent and Denmark Strait in the decades preceding the warming event. | Yes |
| Vettoretti et al. (2022) | ссѕин | Glacial sim- ula- tions | CO2 levels from 190 to 225 ppm | LGM 6G (VM5a) | rgM | None | 8000- 10000 yrs | Spontaneous milternial- scale D-O type oscilla- tions within a window of CO_2 levels from 190 to 225 ppm to 225 ppm | Old sea ice from the Arctic is ex- ported to the NA, sea;ce growth is favoured through ice-albedo feed- back, high-latitude convection is re- duced through sea;ce melt, and con- sequently Anturcis, and Greenhand cool. The interstatial-to-statial transi- tion happens with fast NA sea-ice expan- sion and NADW production collapses. | During a studial, sea ice thins in the Southern Ocean and Antarctica warms. There are increases in salt convergence in the NA, NADW fluctuations are am- plified via sult advection feedback, and the volume of NADW increases, allowed by late-studial decreases in AABW for- mation. Late-studial decreases in AABW for- mation. Late-studial decreases in AABW for- ition experiences thermohaline instabil- ity and the Nordic and Irminger Seas are detabilised, triggering repid sea-ice loss in the NA and the transition from stadial to interstadial states. | Yes |
| Armstrong et al. (2021) | HadCM3B | MIS3 - 30ka | 30ka | 30ka | 30ka | None | 6000 yrs | Spontaneous miltennial- scate D-O type oscilla- tions | Ocean forcing initiates the stadial phase; The collapse of the stalinity gradient be- tween the Northern NA and STO feads to a reduced advection in the Nordic Sats and decretased deep-water forma- tion | The initiation of the interstalial phase is associated with a wind-driven atmo- spheric forcing in the North Seas due to increased regional temperatures, re- duced sea ice cover and increased sea level pressure, which enhances wind stress and convection | Yes, partially. The D-0 type oscillations reflect a salt oscil- lator mechanism in the NA. |
| | | | | | | | | | | Continued on next page | |

| | | | | | Ia | ble A1 – co | Table A1 – continued from previous page | revious page | | | |
|------------------|---------|---------|--------------------------|-------|--------------------------------|-------------|---|--------------------------|--|--|----------------------------------|
| Study | Model | Period | GHG | Ice | Insolation | FWF | Run | Main findings | Mechanisms for D-O cooling event | Mechanisms for D-O warming event | Similar Mechanism to |
| | | | | sheet | | | Length | | | | schematic in Fig. 2 or 3 |
| Zhang et al. | COSMOS | MIS3: | 40ka | 40ka | One transient | None | +5000 yrs | Spontaneous millennial- | Transitions from warm interstadial | While the AMOC is in its weak phase, a | Some similarities: unforced |
| (2021) | | 40- | | | simulation of | | | scale D-O type oscilla- | to cold stadial are linked to (1) a | gradual increase in subsurface tempera- | AMOC oscillations are trig- |
| | | 32ka | | | 40-32ka and | | | tions; Orbitally induced | precession-controlled rise in low- | ture in the subpolar ocean together with | gered by either the tropi- |
| | | | | | one 40ka | | | AMOC changes | latitude boreal summer insolation by | enhanced northward transport of salt in | cal salt impact (linked to |
| | | | | | snapshot | | | | modifying the NA low-latitude hydrocli- | the NA, drive the AMOC back to its | preccesion-controlled summer |
| | | | | | simulation | | | | mate and/or (2) an obliquity-controlled | strong phase. | insolation) and/or the subpo- |
| | | | | | (with 34 ka | | | | reduction in high-latitude annual in- | | lar thermal impact (linked to |
| | | | | | orbital | | | | solation by altering high-latitude sea | | obliquity-controlled mean an- |
| | | | | | parameters) | | | | ice-ocean-atmosphere interactions. | | nual insolation). |
| Brown and Gal- | CM2Mc | Mixed | $CO_2 = 180 \text{ ppm}$ | Η | Obliquity: 22°; | None | more than | Spontaneous millennial- | During a weak AMOC phase, NA deep | During a strong AMOC phase, NA deep | Yes, partially. Salt advection |
| braith (2016) | | forc- | | | Precession: | | 8000 yrs | scale D-O type oscilla- | convection is largely reduced and there | convection is intense and there is a re- | is a key driver of the oscilla- |
| | | ing | | | °06 | | | tions | is an expansion of sea ice in the north- | treat of sea ice in the northeast Atlantic | tions, specifically the salt ex- |
| | | | | | | | | | east Atlantic. Heat accumulates at depth | | change between subpolar and |
| | | | | | | | | | in the NA linked to the weak advection | | subtropical NA. |
| | | | | | | | | | of warm waters from the tropics. | | |
| Klockmann | MPI-ESM | Mixed | $CO_2 = 195-217$ | Η | LGM – 21 ka | None | 8000- | Spontaneous millennial- | Stadial phases correspond to weak | During interstadial phases, the AMOC | Yes, partially. The proposed |
| et al. (2018, | | forc- | ppm; $CH_4 =$ | | | | 12350 yrs | scale D-O type oscilla- | AMOC and strong SPG phases. The ex- | is strong and the SPG is contracted and | mechanism behind the spon- |
| 2020) | | ing | 396-494 ppb; | | | | | tions | tensive SPG results in low northward | weak. There is a broad inflow of salty | taneous AMOC oscillations |
| | | | $N_2 O = 209-227$ | | | | | | salt transport and deep convection only | subtropical water to the subpolar NA. | compromises three com- |
| | | | qdd | | | | | | occurs sporadically in the Iceland basin. | Changes in the SPG are driven by vari- | ponents: (1) oscillations in |
| | | | | | | | | | The Nordic Seas are entirely ice covered | ations in the cross-gyre density differ- | salinity comparable to Peltier |
| | | | | | | | | | which results in a weak Icelandic Low | ence. The eastern NA is fully ice-free. | and Vettoretti (2014), (2) a |
| | | | | | | | | | and therefore in a weak wind stress curl. | Deep convection occurs continuously in | density-driven feedback loop |
| | | | | | | | | | Subsurface waters in the Nordic Seas are | the Iceland basin, Irminger Sea and the | comparable to Montoya et al. |
| | | | | | | | | | around 3 K warmer than during intersta- | Nordic Seas. | (2011), and (3) a wind-driven |
| | | | | | | | | | dial phases. | | feedback loop comparable to |
| | | | | | | | | | | | Drijfhout et al. (2013) and |
| | | | | | | | | | | | Kleppin et al. (2015). |
| Kuniyoshi et al. | MIROC4m | -Mid- | $CO_2 = 220 \text{ ppm}$ | LGM | Obliquity: | None | 6000 yrs | Spontaneous millennial- | Changes in subsurface ocean tempera- | The SPG remains weak (strong) when | Yes, partially. The opening of |
| (2022) | | glacial | | , | 22.949°; | | | scale D-O type oscilla- | ture in the NA plays an important role in | the AMOC is weak (strong), as well | a big polynya determines the |
| | | con- | | ICE- | Eccentricity: | | | tions | modifying the stratification of the verti- | as during the transitions between the | stadial-interstadial transition. |
| | | -ib | | 5G | 0.04; | | | | cal water column and then reversing the | two AMOC modes. There is a positive | Abrupt changes in AMOC |
| | | tions | | | Perihelion: | | | | AMOC mode (thermohaline oscillator) | feedback between AMOC and SPG, in | lead to changes in salt ad- |
| | | | | | 270° and 90° | | | | | agreement with (Li and Born, 2019). | vection with the NA subpolar |
| | | | | | | | | | | | gyre and works as a positive |
| | | | | | | | | | | | feedback. |

| Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run Lenoth | Main findings | |
|---------------|-----------|--------|--------------------------------|-------------------------|------------|------------------------|---------------|-------------------------------------|---|
| Guo et al. | NorESM1-F | MIS3 – | $CO_2 = 215 \text{ ppm};$ | Data-constrained 38 ka | 38 ka | None | 2500 | The equilibrium MIS3 simula- | 1 |
| (2019b) | | 38 ka | $CH_4 = 550 \text{ ppb}; N_2O$ | | | | yrs (re- | tion does not show spontaneous | |
| | | | = 260 ppb | | | | cently | D-O type oscillations. Attempts | |
| | | | | | | | ex- | at perturbing the system into a | |
| | | | | | | | tended | cold stadial state, by modify- | |
| | | | | | | | to 6000 | ing the height of the LIS and | |
| | | | | | | | model | atmospheric CO2 levels, show | |
| | | | | | | | years) | that the modelled MIS3 inter- | |
| | | | | | | | | stadial state is rather stable, and | |
| | | | | | | | | thus questioning the occurrence | |
| | | | | | | | | of spontaneous D-O type oscil- | |
| | | | | | | | | lations in the lack of interactive | |
| | | | | | _ | | | ice sheet-meltwater dynamics. | |
| Zhang and | CCSM3 | MIS3 - | $CO_2 = 215 \text{ ppm};$ | ICE-5G ice sheet | 38 ka | 12 hos- | 2170 | AMOC is more sensitive to | |
| Prange (2020) | | 38 ka | $CH_4 = 501 \text{ ppb}; N_2O$ | configuration (Peltier, | | ing/extraction | yrs | meltwater fluxes under MIS3 | |
| | | | = 234 ppb | 2004). | | experiments | | conditions than under LGM | |
| | | | | | | with freshwater | | conditions. The lower AMOC | |
| | | | | | | fluxes from | | stability under MIS3 conditions | |
| | | | | | | ± 0.005 Sv to | | proposes that D-O type oscilla- | |
| | | | | | | \pm 0.2 Sv, | | tions could have been triggered | |
| | | | | | | injected in the | | by small perturbations in the | |
| | | | | | | Nordic Seas for | | ocean surface meltwater forc- | |
| | | | | | | 500 years. | | ing e.g. linked to ice-sheet pro- | |
| | | | | | _ | | | cesses. | |
| | | | | | | Continued on next page | next page | | |

Table A2: List of simulations run under MIS3/mid-glacial conditions.

| | | | | Table 74 - continued from previous page | | | | |
|-------------------|----------|---------|--------------------------------|---|------------|---------------------------|-----------|-----------------------------------|
| Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run | Main findings |
| | | | | | | | Length | |
| Kawamura | MIROC 4m | Mid- | $CO_2 = 215 \text{ ppm};$ | Intermediate-size | 15 ka | Hosing | More | The climate response to fresh- |
| et al. (2017) | | glacial | $CH_4 = 350 \text{ ppb}; N_2O$ | ice-sheet configuration | | experiments | than | water perturbations is much |
| | | state | = 200 ppb | (at 15 ka) | | with freshwater | 2000 | lower under LGM conditions |
| | | | | | | fluxes of 0.05 | yrs | than under MIS3 conditions. |
| | | | | | | Sv and 0.1 Sv, | | The unperturbed LGM AMOC |
| | | | | | | injected in the | | is unusually weak (around 6 |
| | | | | | | North Atlantic | | Sv) and thus could barely be |
| | | | | | | Ocean (50°N to | | further lessened, such that melt- |
| | | | | | | 70° N) for 500 | | water hosing does not largely |
| | | | | | | years. | | affect the large-scale climate. |
| Vettoretti et al. | CCSM4 | Glacial | $CO_2 = 210 \text{ ppm}$ | TGM - ICE-6G | LGM | Hosing | 0008 | The Heinrich simulation has |
| (2022) | | run | | (VM5a) | | experiment (H | yrs | a large Northern Hemisphere |
| | | | | | | event-like | | temperature and AMOC over- |
| | | | | | | pulse) with | | shoot after the Heinrich sta- |
| | | | | | | freshwater | | dial ends. Nevertheless, this |
| | | | | | | fluxes of 0.05 | | fast AMOC rise above regular |
| | | | | | | Sv for 500 | | interstadial levels is in agree- |
| | | | | | | years in two | | ment with observations only for |
| | | | | | | separate stadial | | a few H-stadial periods (H4 and |
| | | | | | | periods in a | | H5). |
| | | | | | | glacial | | |
| | | | | | | simualtion run | | |
| | | | | | | with CO_2 of | | |
| | | | | | | 210 ppm. The | | |
| | | | | | | freshwater flux | | |
| | | | | | | is injected in | | |
| | | | | | | the North | | |
| | | | | | | Atlantic (50°N | | |
| | | | | | | to 70°N). | | |
| | | | | | | Continued on next page | next page | |

Table A2 – continued from previous page

| Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run | Main findings |
|--------------|----------|---------|-----|-------------------------|------------|------------------------|-----------|----------------------------------|
| • | | | | | | | Length | D |
| Zhang et al. | COSMOS | Mixed | LGM | Sensitivity experiments | LGM | Freshwater | Snapshot | An AMOC bi-stability regime |
| (2014, 2017) | (ECHAM5- | forcing | | applying different | | fluxes from | and | is found under intermediate |
| | JSBACH- | | | heights of the NH ice | | 0.00 Sv to | tran- | CO_2 and ice sheet conditions |
| | (MOIdM) | | | sheets | | ±0.02 Sv, | sient | roughly resembling that of the |
| | | | | | | injected in the | simula- | MIS3 climate. In the bi-stable |
| | | | | | | NA for | tions | MIS3 regime, transitions from |
| | | | | | | 100-300 years | 250- | weak to strong AMOC state |
| | | | | | | | 700 yrs | and vice versa could be initi- |
| | | | | | | | | ated by not only gradual vari- |
| | | | | | | | | ations in LIS height and at- |
| | | | | | | | | mospheric CO2 but also fresh- |
| | | | | | | | | water perturbations. Changes in |
| | | | | | | | | the LIS height can initiate a |
| | | | | | | | | positive atmosphere-ocean-sea |
| | | | | | | | | ice feedback leading to D-O |
| | | | | | | | | type climate shifts. A gradual |
| | | | | | | | | increase in the ice sheet height |
| | | | | | | | | results in a northward shift of |
| | | | | | | | | the winds, and favours a more |
| | | | | | | | | saline Labrador Sea both by re- |
| | | | | | | | | ducing the sea ice/freshwater |
| | | | | | | | | import from the Arctic and in- |
| | | | | | | | | creasing the advection of salt |
| | | | | | | | | into the area. |
| | | | | | | Continued on next page | next page | |

Table A2 – continued from previous page

| StudyModelPeriodCHGIce sheetIce sheetIce sheetIce sheetIce sheetIce sheetIce sheetIce sheetIce sheetIce sheetMain findingsZhamg et al.GYMOSMixelCO2 levels increased40% of LOM ice-sheetLGMNome500 ysSrow an eison and instanceZ0014.017)ECHAMS-forming 500 years.MPOMOmoderateice onlige0.05 ymstate A variation of 15 pmMPOMOMPOMOParaforming 500 years.paraparaparastate A variation of 15 pmMPOMOParaParaParaforming 500 years.paraparastate A variation of 15 pmMPOMOParaParaParaParaparaparaparaMPOMOParaParaParaParaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCO2ParaparaparaparaMPOMOParaCOSParaParaparaparaParaCOSONSMixelCO2ParaparaparaParaCOSONSMixelCO2Paraparapar | | | | | aged shows of more manufactor and a state | | | | |
|--|--------------|----------|---------|-------------------------|---|------------|-----------------|----------|-------------------------------------|
| Image: NormeImage: Norme </th <th>Study</th> <th>Model</th> <th>Period</th> <th>GHG</th> <th>Ice sheet</th> <th>Insolation</th> <th>FWF</th> <th>Run</th> <th>Main findings</th> | Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run | Main findings |
| COSMOS Mixed CO2 levels increased 40% of LGM ice-sheet LGM None 500 yrs JSBACH- forcing from 185 to 205 ppm configuration 500 yrs 500 yrs JSBACH- during 500 years. during 500 years. configuration 600 yrs 500 yrs MPIOM) MPIOM) MPIOM More CO2 levels increased Intermediate ice sheet LGM None 500 yrs MPIOM) Mixed CO2 levels increased Intermediate ice sheet LGM None 600 yrs GECHAMS- from 185 to 239 ppm at volume equivalent to None 600 yrs 600 yrs MPIOM) MPIOM more configuration - ice LGM None 600 yrs SBACH- MPIOM more configuration - ice LGM None 600 yrs SBACH- MPIOM more configuration - ice sea-level drop of 40 m None 600 yrs GECHAMS- forcing from 185 ppm and 245 sea-level drop of 40 m More 600 yrs GECHAMS- forcing from 185 ppm and 245 gen-level drop of | | | | | | | | Length | |
| (ECHAMS- ISBACH- MPIOM)forcing houring 500 years.from 185 to 205 ppm during 500 years.configurationmISBACH- ISBACH- MPIOM)during 500 years.during 500 years.doring 500 years.doring 500 years.MPIOM)mLGMNonedoring moredoring moredoring moredoring moredoring moreMPIOM)MixedCO2 levels increased a rate of 0.02 ppm per year.Internediate ice sheet approximately a sea-level drop of 40mLGMNone600 yrsMPIOM)MixedCO2 levels increased a rate of 0.02 ppm per year.Internediate ice sheet approximately a sea-level drop of 40mLGMPersistent frexwater flux of 0.15 SVMore of 0.15 SV | Zhang et al. | COSMOS | Mixed | CO_2 levels increased | 40% of LGM ice-sheet | LGM | None | 500 yrs | For a moderate ice volume, |
| JSBACH- during 500 years. MPIOM) during 500 years. MPIOM) note during 500 years. during 500 years. MPIOM) note during 500 years. during 500 years. MPIOM) Mixed COSMOS Mixed Mixed CO2 levels increased Intermediate ice sheet LGM None Mone JSBACH- arate of 0.02 ppm per MPIOM) arate of 0.02 ppm per MPIOM) year. sea-level drop of 40m LGM MPIOM) Mixed MPIOM) Mixed MPIOM) CO2 levels increased MPIOM) Mixed MPIOM) CO2 levels increased MPIOM) Mixed MPIOM) CO2 levels increased MPIOM) Mixed MPIOM) Mixed MPIOM) Mixed MPIOM) Mixed MPIOM) Persistent flux MPIOM) MPI a rate of 0.05 | (2014, 2017) | (ECHAM5- | forcing | from 185 to 205 ppm | configuration | | | | 0.25-0.45 times the LGM, two |
| MPIOM)MPIOM | | JSBACH- | | during 500 years. | | | | | stable AMOC modes are iden- |
| Image: NoticeImage: NoticeImage: Notice <th< th=""><th></th><th>(MOIdM)</th><th></th><th></th><th></th><th></th><th></th><th></th><th>tified. A variation of 15 ppm</th></th<> | | (MOIdM) | | | | | | | tified. A variation of 15 ppm |
| Image: Cost of the cost of | | | | | | | | | in atmospheric CO2 concen- |
| COSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrsCOSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrsCOSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrsSBACH-a rate of 0.02 ppm per year.volume equivalent to approximately a sea-level drop of 40 mNone600 yrsMPIOM)MixedCO2 levels increasedLGM ice volumeLGMPersistent600 yrsCOSMOSMixedCO2 levels increasedLGM ice volumeLGMPersistent600 yrsSBACH-ppm at a rate of 0.05ppm at a rate of 0.05mPIOM)Persistent flux600 yrs | | | | | | | | | tration - equivalent to changes |
| COSMOSMixedCO2Intermediate ice sheetLGMNone600 yrsCOSMOSMixedCO2levels increasedIntermediate ice sheetLGMNone600 yrs(ECHAM5-froringfrom 185 to 239 ppm atconfiguration - iceNone600 yrs600 yrsJSBACH-a rate of 0.02 ppm pervolume equivalent toapproximately asea-level drop of 40 m600 yrsMPIOM)MixedCO2levels increasedLGM ice volumeLGMPersistent600 yrsCOSMOSMixedCO2levels increasedLGM ice volumeLGMPersistent600 yrsSBACH-forcingfrom 185 ppm and 245ppm at arate of 0.05mpm or year.of 0.15 Svof 0.15 Sv | | | | | | | | | during D-O cycles containing |
| COSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrsCOSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrs(ECHAM5-forcingfrom 185 to 239 ppm atconfiguration - iceLGMNone600 yrsJSBACH-a rate of 0.02 ppm pervolume equivalent tovolume equivalent tosea-level drop of 40 mSea-level drop of 40 mSea-level drop of 40 mMPIOM)MixedCO2 levels increasedLGM ice volumeLGMPersistent600 yrs(ECHAM5-forcingfrom 185 ppm and 245from 185 ppm and 245of 0.15 Svof 0.15 SvJSBACH-ppm at a rate of 0.05ppm at a rate of 0.05ppm at a rate of 0.05of 0.15 Svof 0.15 Sv | | | | | | | | | HE - is enough to trigger oscil- |
| COSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone600 yrs(ECHAM5-forcingfrom 185 to 239 ppm atconfiguration - iceLGMNone600 yrsJSBACH-a rate of 0.02 ppm pervolume equivalent toarate of 0.02 ppm pervolume equivalent toJSBACH-proximately asea-level drop of 40 mEGMForeing600 yrsMPIOM)MixedCO2 levels increasedLGM ice volumeLGMPersistent600 yrsCOSMOSMixedCO2 levels increasedLGM ice volumeLGMrestrent flux600 yrsJSBACH-forcingfrom 185 ppm and 245ppm at a rate of 0.05ycer.of 0.15 Sv90 yrsMPIOM)Ppm at a rate of 0.05ppm per year.ppm per year.of 0.15 Sv90 yrs90 yrs | | | | | | | | | lations between a weak stadial |
| COSMOSMixed CO_2 levels increasedIntermediate ice sheet LGM None600 yrs(ECHAM5-forcingfrom 185 to 239 ppm atconfiguration - ice LGM None600 yrsJSBACH-a rate of 0.02 ppm pervolume equivalent tovolume equivalent to $PearcPearcJSBACH-mPIOM)year.a rate of 0.02 ppm persea-level drop of 40 mPercPercMPIOM)MixedCO_2 levels increasedLGM ice volumeLGMPersistent600 yrs(ECHAM5-forcingfrom 185 ppm and 245tom 185 ppm and 245Persistent flux600 yrsJSBACH-ppm at a rate of 0.05ppm at a rate of 0.05ppm at a rate of 0.05of 0.15 SvPersistent flux$ | | | | | | | | | state to a strong interstadial cir- |
| COSMOSMixedCO2 levels increasedIntermediate ice sheetLGMNone(ECHAM5-forcingfrom 185 to 239 ppm atconfiguration - iceNoneNoneJSBACH-a rate of 0.02 ppm pervolume equivalent toapproximately asealevel drop of 40 mMPIOM)year.sea-level drop of 40 msea-level drop of 40 mPersistentCOSMOSMixedCO2 levels increasedLGM ice volumeLGM ice volume(ECHAM5-forcingfrom 185 ppm and 245Sea-level drop of 40 mPersistentJSBACH-ppm at a rate of 0.05ppm at a rate of 0.05MPIOM)ppm at a rate of 0.05 | | | | | | | | | culation state. |
| (ECHAM5-forcingfrom 185 to 239 ppm at a rate of 0.02 ppm perconfiguration - ice volume equivalent toJSBACH-a rate of 0.02 ppm pervolume equivalent to approximately aMPIOM)volume equivalent to sea-level drop of 40 mPersistent freshwater fluxCOSMOSMixedCO2 levels increased from 185 ppm and 245LGM ice volumeISBACH-forcingfrom 185 ppm and 245Persistent freshwater fluxMPIOM)MPIOM)ppm at a rate of 0.05of 0.15 Sv | Zhang et al. | COSMOS | Mixed | CO_2 levels increased | Intermediate ice sheet | LGM | None | 600 yrs | |
| JSBACH-a rate of 0.02 ppm pervolume equivalent toMPIOM)a rate of 0.02 ppm perapproximately aMPIOM)year.approximately aMPIOM)sea-level drop of 40 mCOSMOSMixedCO2 levels increasedMixedCO2 levels increasedLGM ice volumeCOSMOSforcingfrom 185 ppm and 245JSBACH-ppm at a rate of 0.05of 0.15 SvMPIOM)mpromer year.ppm per year. | (2014, 2017) | (ECHAM5- | forcing | from 185 to 239 ppm at | configuration - ice | | | | |
| MPIOM)year.approximately a sea-level drop of 40 mperiodCOSMOSMixedCO2 levels increasedLGM ice volumePersistentCOSMOSfroringfrom 185 ppm and 245LGM ice volumefreshwater fluxJSBACH-ppm at a rate of 0.05ppm at a rate of 0.05of 0.15 SvMPIOM)ppm per year.ppm per year.ppm | | JSBACH- | | a rate of 0.02 ppm per | volume equivalent to | | | | |
| Image: Mode of Mode Section Image: Section Image: Section Image: Section COSMOS Mixed CO2 Evels increased LGM ice volume Persistent COSMOS Mixed CO2 Evels increased LGM ice volume Persistent (ECHAM5- forcing from 185 ppm and 245 LGM ice volume LGM Persistent JSBACH- ppm at a rate of 0.05 ppm per year. of 0.15 Sv | | MPIOM) | | year. | approximately a | | | | |
| COSMOSMixedCO2 levels increasedLGM ice volumeLGMPersistent(ECHAM5-forcingfrom 185 ppm and 245freshwater fluxfreshwater fluxJSBACH-ppm at a rate of 0.05ppm at a rate of 0.05of 0.15 SvMPIOM)ppm per year.ppm per year.of 0.15 Sv | | | | | sea-level drop of 40 m | | | | |
| (ECHAM5-forcingfrom 185 ppm and 245JSBACH-ppm at a rate of 0.05MPIOM)ppm per year. | Zhang et al. | COSMOS | Mixed | CO_2 levels increased | LGM ice volume | LGM | Persistent | 600 yrs | |
| ppm at a rate of 0.05 ppm per year. | (2014, 2017) | (ECHAM5- | forcing | from 185 ppm and 245 | | | freshwater flux | | |
| | | JSBACH- | | ppm at a rate of 0.05 | | | of 0.15 Sv | | |
| | | MPIOM) | | ppm per year. | | | | | |

Table A2 – continued from previous page

| Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run Length | Main findings |
|----------------|-----------|------------|--------------------------------|-----------|-----------------|------------------------|---------------|----------------------------------|
| Brugger et al. | CLIMBER | 415 Ma | $CO_2 = 1500 \text{ ppm. Few}$ | None | S0=1315 | None | 5000 | Decadal to centennial tempera- |
| (2019) | 3α | and | runs with 500 ppm, | | W/m^2 (415 | | yrs | ture fluctuations at high north- |
| | | 380 Ma | 800 ppm, 2000 ppm | | Ma), 1319 | | | ern latitudes. |
| | | | | | W/m^2 (380 | | | |
| | | | | | Ma), various | | | |
| | | | | | obliquity | | | |
| | | | | | values, | | | |
| | | | | | eccentricities, | | | |
| | | | | | precession | | | |
| | | | | | angles | | | |
| Rind et al. | GISS E2-R | Future | $4 \text{ x PI } CO_2$ | PD | PD | None | 2500 | Multi-centennial cessation with |
| (2018) | (TCADI) | warm- | | | | | yrs | restoration and rapid over- |
| | | ing | | | | | | shooting in NADW formation |
| | | simula- | | | | | | |
| | | tion | | | | | | |
| | | (exper- | | | | | | |
| | | iment | | | | | | |
| | | "abrupt | | | | | | |
| | | 4 x | | | | | | |
| | | C02" | | | | | | |
| | | in Rind | | | | | | |
| | | et al. | | | | | | |
| | | (2018) | | | | | | |
| | | | | | | Continued on next page | n next page | |

Table A3: Possible oscillatory behaviour for other time periods

| | | | lable A3 – continut | 1able A3 - continued from previous page | | | | |
|-------------|-----------|---------|------------------------|---|------------|------|--------|---------------------------------|
| Study | Model | Period | GHG | Ice sheet | Insolation | FWF | Run | Main findings |
| | | | | | | | Length | |
| Rind et al. | GISS E2-R | Future | CO_2 rises at 1% per | PD | PD | None | 2500 | Multi-centennial cessation with |
| (2018) | (TCADI) | warm- | year until 4 x CO_2 | | | | yrs | restoration and rapid over- |
| | | ing | (after 140 years) and | | | | | shooting in NADW formation |
| | | simula- | then held constant. | | | | | |
| | | tion | | | | | | |
| | | (exper- | | | | | | |
| | | iment | | | | | | |
| | | "1 pct | | | | | | |
| | | C02" | | | | | | |
| | | in Rind | | | | | | |
| | | et al. | | | | | | |
| | | (2018) | | | | | | |
| Rind et al. | GISS E2-R | Future | rcp85 emissions | PD | Δd | None | 4300 | Multi-centennial cessation with |
| (2018) | (TCADI) | warm- | | | | | yrs | restoration and rapid over- |
| | | ing | | | | | | shooting in NADW formation |
| | | simula- | | | | | | |
| | | tion | | | | | | |
| | | (exper- | | | | | | |
| | | iment | | | | | | |
| | | "rcp85" | | | | | | |
| | | in Rind | | | | | | |
| | | et al. | | | | | | |
| | | (2018) | | | | | | |

Table A3 – continued from previous page

Appendix B

Table B1. Summary of LGM/MIS3-like simulations discussed in the text. Highlighted in red the models that reproduce D-O type oscillations.

| Study | Model | Period | N° of simulations | Run length |
|-------------------------------|----------------|------------------------|--------------------------|-----------------------|
| Peltier and Vettoretti (2014) | UofT CCSM4 | PMIP4 LGM | 1 | 5000 |
| Lohmann et al. (2020) | AWI-ESM1-1-LR | PMIP4 LGM | 1 | 1300 |
| Sidorenko et al. (2019) | AWI-ESM-2-1-LR | PMIP4 LGM | 1 | 600 |
| Tierney et al. (2020) | CESM1.2 | PMIP4 LGM | 1 | 1800 |
| Valdes et al. (2017) | HadCM3B-M2.1aD | PMIP4 LGM | 3 | 400-2900 ^a |
| Lhardy et al. (2021) | iLOVECLIM1.1.4 | PMIP4 LGM | 2 | 5000 |
| Volodin et al. (2018) | INM-CM4-8 | PMIP4 LGM | 1 | 50 |
| Sepulchre et al. (2020) | IPSLCM5A2 | PMIP4 LGM | 1 | 1200 |
| Ohgaito et al. (2021) | MIROC-ES2L | PMIP4 LGM | 1 | 8960 |
| Mauritsen et al. (2019) | MPI-ESM1.2 | PMIP4 LGM | 1 | 3850 |
| Armstrong et al. (2021) | HadCM3B-M2.1aD | MIS3 (30 ka) | 1 | 6000 |
| Zhang et al. (2021) | COSMOS | MIS3 (40-32 ka) | 2 | 5000 |
| Guo et al. (2019b) | NorESM | MIS3 (38 ka) | 1 | +6000 |
| Zhang and Prange (2020) | CCSM3 | MIS3 (38 ka) | 1 | 2170 |
| Kawamura et al. (2017) | MIROC4m | Mid-glacial conditions | 1 | +2000 |
| Kuniyoshi et al. (2022) | MIROC4m | Mid-glacial conditions | 2 | 6000 |
| Vettoretti et al. (2022) | CCSM4 | Glacial conditions | 4^b | 8000 |
| Brown and Galbraith (2016) | CM2Mc | Mixed forcing | 1 | +8000 |
| Klockmann et al. (2018) | MPI-ESM | Mixed forcing | 3 | +8000 |
| Zhang et al. (2014) | COSMOS | Mixed forcing | 11^c | 300-4000 ^d |

^{*a*} Only one simulation run longer than 2000 model years

 b Four simulations run with CO_{2} levels: 200, 210, 220, 225 ppm

 c We do not consider FWF runs nor transient simulations forced with varying CO_{2} and/or NH ice sheet height

^d Only two simulations with a duration of +2000 model years

Appendix C

| Model | Period | Run length |
|-----------------|--------|------------|
| ACCESS-ESM1-5 | - | - |
| AWI-ESM-1-1-LR | LGM | 1300 |
| CESM2 | - | - |
| CNRM-CM6-1 | - | - |
| EC-Earth3-LR | - | - |
| FGOALS-f3-L | - | - |
| FGOALS-g3 | - | - |
| GISS-E2-1-G | - | - |
| HadGEM3-GC31-LL | - | - |
| INM-CM4-8 | LGM | 50 |
| IPSL-CM6A-LR | - | - |
| MIROC-ES2L | LGM | +8000 |
| MPI-ESM1-2 | LGM | 3850 |
| MRI-ESM2-0 | - | - |
| NESM3 | - | - |
| NorESM1-F | MIS3 | 6000 |
| NorESM2-LM | - | _ |