



1 Hydrogeological controls on the spatio-temporal variability of surge- 2 induced hydraulic gradients along coastlines: implications for beach 3 surface stability

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13 **Abstract.** Ocean surges pose a global threat for coastal stability. These hazardous events alter flow conditions and pore
14 pressures in flooded beach areas during both inundation and subsequent retreat stages, which can mobilize beach material,
15 potentially enhancing erosion significantly. In this study, the evolution of surge-induced pore-pressure gradients is studied
16 through numerical hydrologic simulations of storm surges. The spatiotemporal variability of critically high gradients is
17 analyzed in 3D. The analysis is based on a threshold value obtained for momentary liquefaction of beach materials under
18 groundwater seepage. Simulations of surge events show that during the run-up stage, head gradients can rise to the calculated
19 critical level landward of the advancing inundation line. During the receding stage, critical gradients were simulated seaward
20 of the retreating inundation line. These gradients reach maximum magnitudes just as sea level returns to pre-surge level, and
21 are most accentuated beneath the still-water shoreline, where the model surface changes slope. The gradients vary along the
22 shore owing to variable beach morphology, with the largest gradients seaward of intermediate-scale (1-3m elevation)
23 topographic elements (dunes) in the flood zone. These findings suggest that the common practices in monitoring and mitigating
24 surge-induced failures and erosion, which typically focus on the flattest areas of beaches, might need to be revised.



25 1 Introduction

26 Groundwater seepage can destabilize land areas, especially at the interface between terrestrial and submerged systems
27 (Iverson, 1995; Iverson & Major, 1986; Iverson & Reid, 1992; Schorghofer et al., 2004; Stegmann et al., 2011). Recent studies
28 have examined the characteristics of pore pressure behavior, the associated groundwater seepage, and its effect on the stability
29 of geomaterials (soils, rocks, etc.), including field observations (Mory et al., 2007; Sous et al., 2016), physical experiments
30 (Schorghofer et al., 2004; Sous et al., 2013), numerical simulations (Orange et al., 1994; Rozhko et al., 2007; Schorghofer et
31 al., 2004), and analytical models (Sakai et al., 1992; Yeh & Mason, 2014). There are several examples of seepage-induced
32 failure of the surface (i.e. the mobilization of the soil skeleton) from around the world, including Japan (Yeh & Mason, 2014),
33 California (Orange et al., 2002), and France (Sous et al., 2016; Stegmann et al., 2011).

34 Soil liquefaction occurs when pore pressures in the geomaterial rise to a point where its effective stress drops to zero and the
35 material is fluidized, and thus acts as a liquid. At the coast, ocean (waves, surge, tides, inundation) and terrestrial (groundwater
36 heads, precipitation, and overland flows) processes concurrently contribute to changing pore pressures in beach and nearshore
37 sediments, and changes in pore pressure distributions and gradients could induce failure of the surface. Ocean effects on pore
38 pressures, groundwater flow, and seepage occur due to wind waves, storm surges, and tsunamis. For example, a 1D analytical
39 model suggests that during a tsunami, vertical hydraulic gradients can destabilize sediments and increase the potential for
40 sediment momentary liquefaction, consistent with laboratory experiments (Abdollahi & Mason, 2020; Yeh & Mason, 2014).
41 Laboratory experiments (Sous et al., 2013) suggest that the magnitude of hydraulic gradients in the beach due to infiltration
42 from sea-swell and infragravity waves depend on the wave frequency, cross-shore position, water table overheight, and the
43 presence of standing waves. A large-scale (250 m) flume study of a barrier island showed that waves can alter the coastal
44 groundwater head distribution significantly, and can change cross-island and local (under the ocean beach) hydraulic gradient
45 directions (Turner et al., 2016). Field observations of pore-pressure over several tidal cycles in a microtidal beach (Sous et
46 al., 2016) suggest that breaking-wave-driven onshore increases in the water surface (setup) over the 10 m nearest the shoreline
47 induced groundwater head changes of O(0.1 m) (Sous et al., 2016). Furthermore, density-driven flow at the subsurface
48 transition zone between fresh terrestrial groundwater and saline groundwater can produce intense, localized seepage (Burnett
49 et al., 2006). Rapid changes in seepage characteristics (locations, magnitudes, direction) during extreme events may lead to
50 sediment liquefaction (i.e., loss of particle-to-particle contacts and sediment effective stresses) and mobilization, resulting in
51 erosion and structure destabilization.

52 Observations, theories, and simulations have shown that the pore-pressure changes owing to energetic ocean waves can reduce
53 effective stresses and may cause liquefaction (Chini & Stansby, 2012; Mory et al., 2007; Sakai et al., 1992; Sous et al., 2013;
54 Yeh & Mason, 2014). Measured pore-pressure changes in beach sediments during intense waves suggest that momentary
55 liquefaction may occur at shallow depths (<1 m) below the surface (Mory et al., 2007), consistent with theory (Sakai et al.,
56 1992). Analytical solutions for the effective stress in an idealized seabed suggest that waves can alter the stresses in the upper
57 meters of the seafloor significantly (Mei & Foda, 1981; Sakai et al., 1992). Simulations of a theoretical 2D porous medium,



58 where an increase in pore pressure is applied at the bottom of the layer from a point source, revealed that different spatial
59 failure patterns (i.e. the geometry of the slip surface) can occur under various stress regimes (i.e. distribution of stresses in the
60 soil) (Rozhko et al., 2007), although the process that leads to the simulated change in the pore-pressure distribution was
61 unexplored.

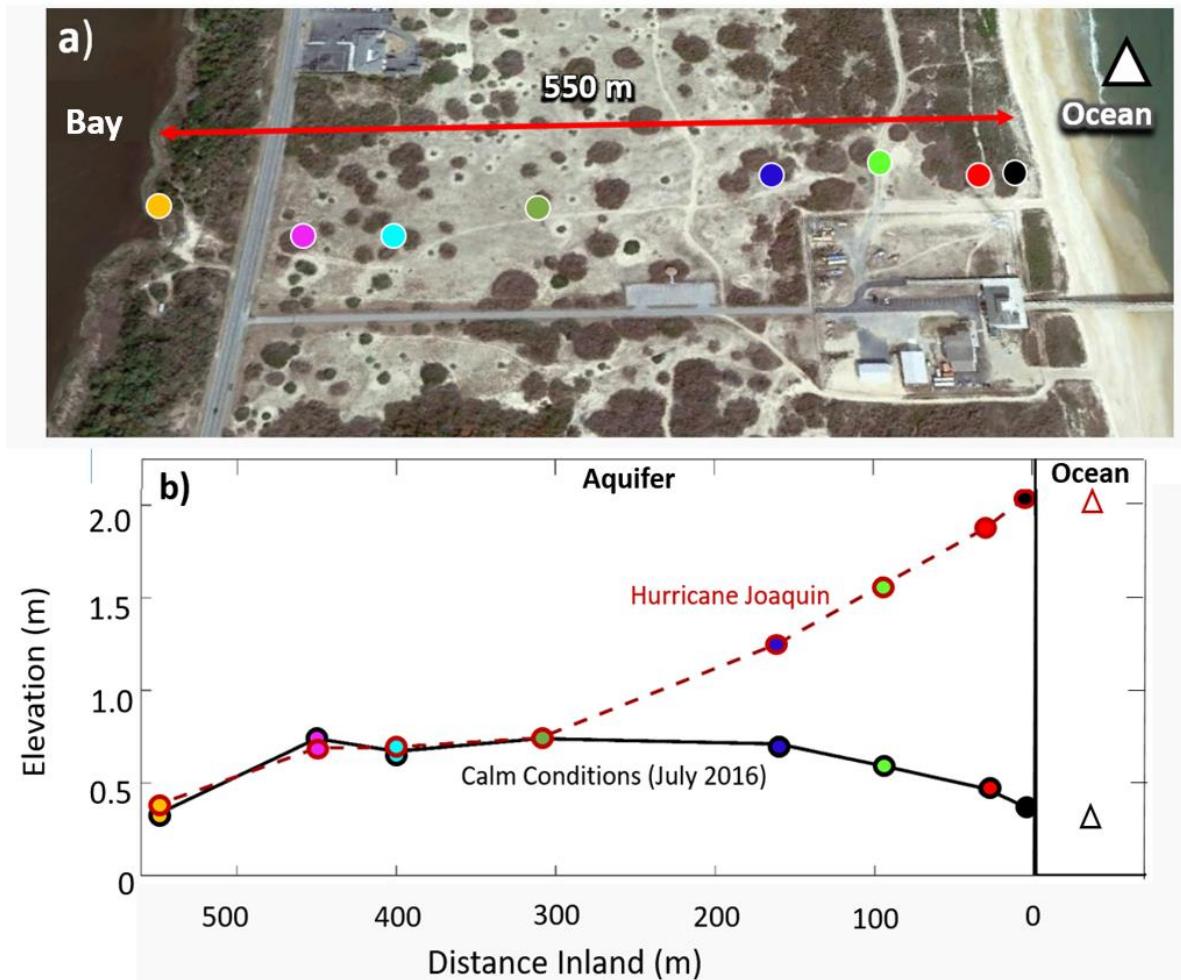
62 Apart from waves, storm surges also could alter the onshore hydrogeological regime and potentially reduce the stability of the
63 beach surface, yet surges have not been explored in this context. This work focuses on the influence of alongshore topography
64 and hydrogeological factors on geotechnical impacts near the shoreline owing to ocean surges driven by coastal storms, which
65 are projected to intensify and become more frequent in the future (Chini & Stansby, 2012; Tebaldi et al., 2012). In particular,
66 the three-dimensional dynamics of surge-induced inundation and the resulting shore-parallel distribution of pore-pressure
67 gradients in sandy beach areas are not well understood. Specific questions addressed in this work are: (1) Can surge-induced
68 pore pressure changes promote sediment liquefaction of the uppermost sediment layers (<5 m), and which areas across the
69 beach are the most vulnerable? (2) What is the relationship between beach morphology and the spatio-temporal evolution of
70 pore pressure gradients? (3) How do the hydrogeological properties (hydraulic conductivity, groundwater recharge) of the
71 coastal system affect the potential for failure? Field evidence is presented for the effect of storm surges on coastal groundwater
72 heads (Section 2), a criterion is derived (Section 3) for momentary soil liquefaction for beach slopes with groundwater
73 discharge based on existing solutions (Briaud, 2013), and a model framework is described (Section 4) and used to simulate
74 surges in theoretical beach settings and to examine their effect on sediment stability (Section 5).

75 **2 Field evidence for hydraulic head changes during storm surges**

76 Groundwater observations collected every 10 min from October 2014 to November 2017 in 8 wells deployed across a 500-m
77 wide barrier island on the Outer Banks of NC, near the town of Duck (Figure 1a) indicate that coastal storm waves and surge
78 significantly affect the freshwater equivalent heads from the beach to more than 310 m inland of the beach (Housego et al.,
79 2018). The study period included 27 storm events (including 4 hurricanes) in which wave heights measured in 26-m water
80 depth (NDBC Station 44100) often exceeded 3.5 m, surge (NOAA tide gauge 8651371) was between 0.5 and 1.0 m, and 36-
81 hr-averaged (to remove fluctuations owing to tides and wind wave motions) shoreline water levels increased from about 0.6
82 to 2.4 m owing to surge and wave-driven setup (included in the simulated surge height). In response to the increased ocean
83 water levels, the groundwater level under the ocean dunes rose 0.5 to 2.0 m. For example, following the passage of Hurricane
84 Joaquin in 2015, which caused offshore wave heights of 4.7 m (and <1 cm of rainfall), head levels under the ocean dunes and
85 25, 90, 160, and 310 m farther inland increased 1.6, 1.4, 1.2, 0.9, and 0.5 m above pre-storm levels, respectively (Figure 1b).
86 These and other storm-driven increases in head levels changed the direction of the hydraulic gradient from toward the bay
87 (inland) during calm conditions to toward the ocean during storms (compare black and red points in Figure 1b under calm
88 conditions with those during the storm). After the shoreline water level returns to pre-storm conditions, the water table behind
89 the dune remains elevated and groundwater discharges back out through the beach as the water table recovers. During the



90 storm, the horizontal location of the shoreline remained more than 10 m seaward of the dunes, and thus there was no inundation
91 from overtopping, which could increase groundwater levels even farther inland. Changes in hydraulic gradients, including the
92 effects of inundation, are investigated in Section 4 with a numerical model that does not mimic the conditions in this field site,
93 but is a generalized representation of coastal hydrogeological systems.



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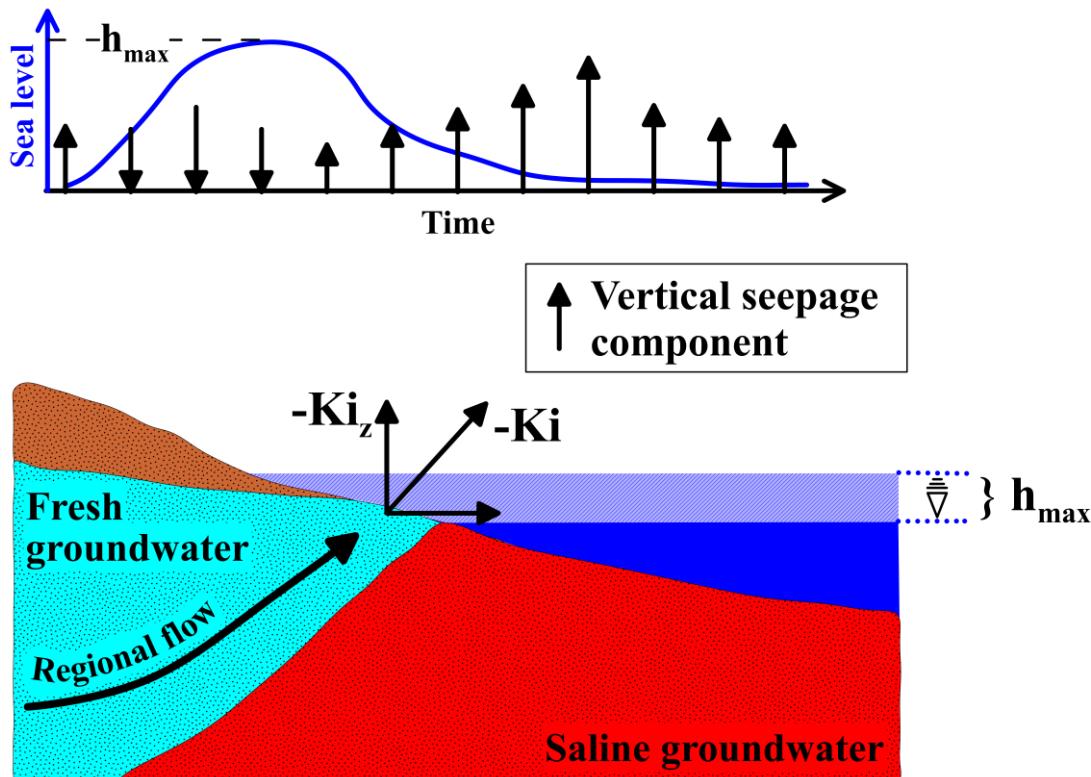
95 **Figure 1:** a) © Google Earth image of the Outer Banks near Duck, NC, with the locations of groundwater wells (colored circles). b)
96 Elevation of the ocean level (triangles) and 36 hr-avg. freshwater equivalent groundwater heads (circles) vs. inland distance from
97 the dune ($x=0$ m). Colors correspond to colors of symbols in (a)) for the average of the calm conditions in July 2016 (black triangle
98 and circle outlines connected by black lines) and at the peak of Hurricane Joaquin (red triangle and circle outlines connected by red
99 dashed lines).

100



101 3 Conceptual model and governing equations

102 A conceptual model of a coastal system (Figure 2) includes infiltration of rain that recharges the aquifer with freshwater,
103 resulting in fresh groundwater flow toward the ocean. In the nearshore area (typically within meters of the shoreline), an
104 inclined freshwater-saltwater transition zone develops between the saline groundwater underlying the seafloor and the
105 terrestrial fresh groundwater. The density gradient at the transition zone deflects the fresh groundwater flow upward, and
106 produces focused groundwater discharge near the coastline that can be amplified by an order of magnitude or more relative to
107 the average flow rate in the aquifer (Paldor et al., 2020). In phreatic aquifers, submarine groundwater discharge typically occurs
108 within tens of meters of the coastline, depending on the recharge rates and aquifer properties (Bratton, 2010). In systems where
109 the discharge is into a body of freshwater (e.g., a lake), the bottom of the lake is a constant head boundary, and thus the seepage
110 is, by definition, perpendicular to the lakebed. This assumption is widely adopted in geotechnical calculations of groundwater
111 discharge magnitudes. For example, in flow net solutions for classic dam and levee problems, the bottom of the river on both
112 sides of the dam or levee is considered an equipotential line (Briaud, 2013). However, along the bottom of a saltwater body
113 the freshwater-equivalent head is variable with bathymetry, and hence the seepage is not necessarily perpendicular to the
114 seafloor and possibly represents a complex, three-dimensional problem with high spatiotemporal variability. To assess the risk
115 of liquefaction in the context of the freshwater-saltwater transition zone and during coastal inundation events, the vertical
116 component of the hydraulic gradient is computed to evaluate the potential for liquefaction (as will be derived in the following
117 section) with the application of the variable-head boundary condition and the inclusion of variable-density flow solutions. It
118 should be highlighted that in the current work, no effects of long-term loading and residual liquefaction were investigated.
119 Hereinafter, the vertical hydraulic gradients will be discussed rather than the pore pressures or heads. In the next section the
120 equations for soil failure potential in terms of the head gradients are derived based on previous derivations (Briaud, 2013). The
121 magnitude of the hydraulic head gradient (Figure 2), which according to Darcy's law is the magnitude of the seepage vector
122 divided by the hydraulic conductivity, is denoted i . Other variables used in the following calculations are shown in Figure 2
123 and summarized in Table 1.



124

125 Figure 2: A typical coastal hydrogeological system. Regional fresh (light blue) groundwater flows to the sea and upward due to
126 variable-density flow along the freshwater-saltwater (red) interface. In the nearshore area, focused groundwater discharge occurs
127 either into the sea (blue) or along a seepage face onshore. As shown in the top of the figure, when the surge begins, the direction of
128 flow reverses (infiltration), and when the sea level reaches its maximal level (h_{max}) the surge retreats and the direction reverts back
129 (exfiltration). The upward (positive vertical component) of flow reaches a maximum when the sea level is back to pre-surge level,
130 before decaying to the steady-state magnitude.

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141 **Table 1: Variables used in the theoretical calculations and numerical simulations.**

| Parameter | Symbol | Value | Unit | Source |
|---|---------------------|-----------|-------------------|------------------------|
| Hydraulic conductivity | K | 10-100 | m/d | Freeze & Cherry (1979) |
| Anisotropy | K_x/K_z | 10 | | |
| Seawater density | ρ_{sw} | 1025 | Kg/m ³ | |
| Freshwater density | ρ_{fw} | 1000 | Kg/m ³ | |
| Local water density | ρ_w | 1000-1025 | Kg/m ³ | |
| Solid material density | ρ_s | 2650 | Kg/m ³ | |
| Freshwater influx | q_0 | 0.01-0.04 | m/d | |
| Aquifer storativity | S_s | 10^{-4} | 1/m | Freeze & Cherry (1979) |
| Porosity | n | 0.3 | | |
| Longitudinal/Transverse Dispersivity | α_L/α_T | 1/0.1 | m | Gelhar et al. (1992) |
| Maximum surge height | h_{0max} | 3 | m | Chini & Stansby (2012) |

142

143 **3.1 The criterion for liquefaction under groundwater seepage**

144 Some publications distinguish between the terms “liquefaction” and “quick sand”, with the former being used for earthquake-
 145 induced fluidization of the soil, and the latter being related to failure due to upward flow (Briaud, 2013). However, the physical
 146 meaning of the two is the same – geomaterial becoming weightless, which can result in erosion and sediment mobilization, or
 147 loss of support of any infrastructure built into the soil. Here, the term liquefaction is used, although the analysis refers to surge-
 148 induced changes in the subsurface flow rather than seismically induced flows. Following Briaud (2013), sand liquefaction
 149 occurs when the pore pressure (u_w) at a certain depth (z) exceeds the total stress (σ), i.e. when the effective stress (σ') goes to
 150 zero:

$$\sigma' = \sigma - u_w \leq 0 \quad (1)$$

151 Neglecting the possibility that gas is still trapped in the pores and assuming a submerged unit weight can be applied, the
 152 criterion for localized, momentary liquefaction in inundated regions can be written in a gradient form (Goren et al., 2013), in
 153 which the vertical pore pressure gradient (positive downward gradient generates upwards flow) exceeds the submerged unit
 154 weight of the soil (γ_{sub}):



$$\gamma_{sub} + \frac{\partial u_w}{\partial z} \leq 0 \quad (2)$$

155 where

$$\gamma_{sub} = (1 - n) \cdot (\rho_s - \rho_{fw}) \cdot g \quad (3)$$

156

157 in which ρ_s is the density of the beach material (sand), and ρ_w is the density of the local water, which has a value between that
158 of seawater ($\rho_{sw} \approx 1025 \text{ kg/m}^3$) and freshwater ($\rho_{fw} \approx 1000 \text{ kg/m}^3$). This failure criterion is similar to Yeh and Mason
159 (2014), who studied liquefaction of a fully saturated sediment following a tsunami.

160 The constant value of porosity ($n=0.3$) is typical for sandy soils, but neglects localized variations in sand bulk density in the
161 simulated area. Furthermore, it is noted that the use of the submerged unit weight of soil is likely an underestimate of the actual
162 unit weight for soils under storm-surge conditions, since saturated conditions may prevail prior to inundation and the saturated
163 unit weight is higher than the submerged ($\gamma_{sub} = \gamma_{sat} - \gamma_{fw}$). However, this work aims to harness a hydrologic modeling
164 framework to assess the spatio-temporal distribution of surge-induced changes in hydraulic gradients. To that end, the
165 liquefaction assessment is limited to the effects of vertical pressure gradients, momentary liquefaction, and the application of
166 the submerged unit weight. It should be noted that studies have shown partially saturated sediments (e.g., in inundation areas)
167 are typically prone to momentary liquefaction (Mory et al., 2007; Yeh and Mason, 2014). Mory et al. (2007) showed that even
168 a 6% air content may alter the potential for momentary liquefaction. For the gradient-form criterion to hold, this condition
169 would need to be met continuously from the surface to the depth of the liquefied layer (Goren et al. 2013), as accounted for in
170 the analysis below.

171 Here, the momentary liquefaction criterion is related to vertical components of seepage vectors to compare the results of the
172 groundwater model with the failure criterion. The 3D model considered here (see below) could be used to examine the
173 horizontal components too, and to analyze the potential for shear failure, not only for momentary liquefaction (Zen et al.,
174 1998). However, for the sake of simplicity and in the interest of focusing on the questions addressed here, such an expansion
175 is not attempted in the current study. It would require further assumptions on the soil characteristics (internal friction, cohesion)
176 and a localized analysis of the local slopes for each point in the domain. According to Darcy's law the vertical flow velocities
177 (v_z) are equal to the product of the (local) vertical head gradient and the vertical hydraulic conductivity K_z :

$$v_z = -K_z \left(\frac{1}{\rho_{fw}g} \frac{\partial u_w}{\partial z} + 1 \right) \quad (4)$$

178

179 thus, the vertical pressure gradient becomes

$$\frac{\partial u_w}{\partial z} = -\rho_{fw}g \left(\frac{v_z}{K_z} + 1 \right) \quad (5)$$

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181 Substituting Equations 3 and 5 into Equation 2 yields:

$$(1 - n) \cdot (\rho_s - \rho_{fw}) \cdot g - \rho_{fw}g \left(\frac{v_z}{K_z} + 1 \right) \leq 0 \quad (6)$$

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183 From Equation 6, the value of the critical vertical head gradient (i_c) is that above which the effective stress is zero or less:

$$\left(\frac{v_z}{K_z} \right)_c \equiv i_c = (1 - n) \cdot \frac{\rho_s - \rho_{fw}}{\rho_{fw}} - 1 \quad (7)$$

184 This result is similar to that derived by Briaud (2013), but here it is derived for saturated groundwater flow, which is the
185 appropriate formulation for the scenario of surge-induced changes in the groundwater flow regime. Using Darcy's law in this
186 context assumes that during the surge the groundwater flow remains largely laminar, which is likely for storm-surge conditions
187 and is a common assumption in similar studies(Abdollahi & Mason, 2020; Guimond & Michael, 2021; Paldor & Michael,
188 2021; Yang et al., 2013; Yu et al., 2016). For convenience, the magnitude of negative (destabilizing) vertical head gradients
189 which initiate positive vertical velocities, is hereinafter denoted i_z and presented in positive values. Using typical values for
190 porosity, solid particle density, and freshwater density for beach material ($n = 0.3$; $\rho_s = 2650 \text{ kg/m}^3$; $\rho_{fw} =$
191 1000 kg/m^3 , respectively), Equation 7 suggests the critical value of vertical head gradient is about $i_c = 0.15$. The following
192 analyses use this value as a threshold for liquefaction, with simulated values of i_z normalized by the critical value $i_c = 0.15$
193 as the seepage-liquefaction factor (SLF):

$$SLF = \frac{i_z}{i_c} \quad (8)$$

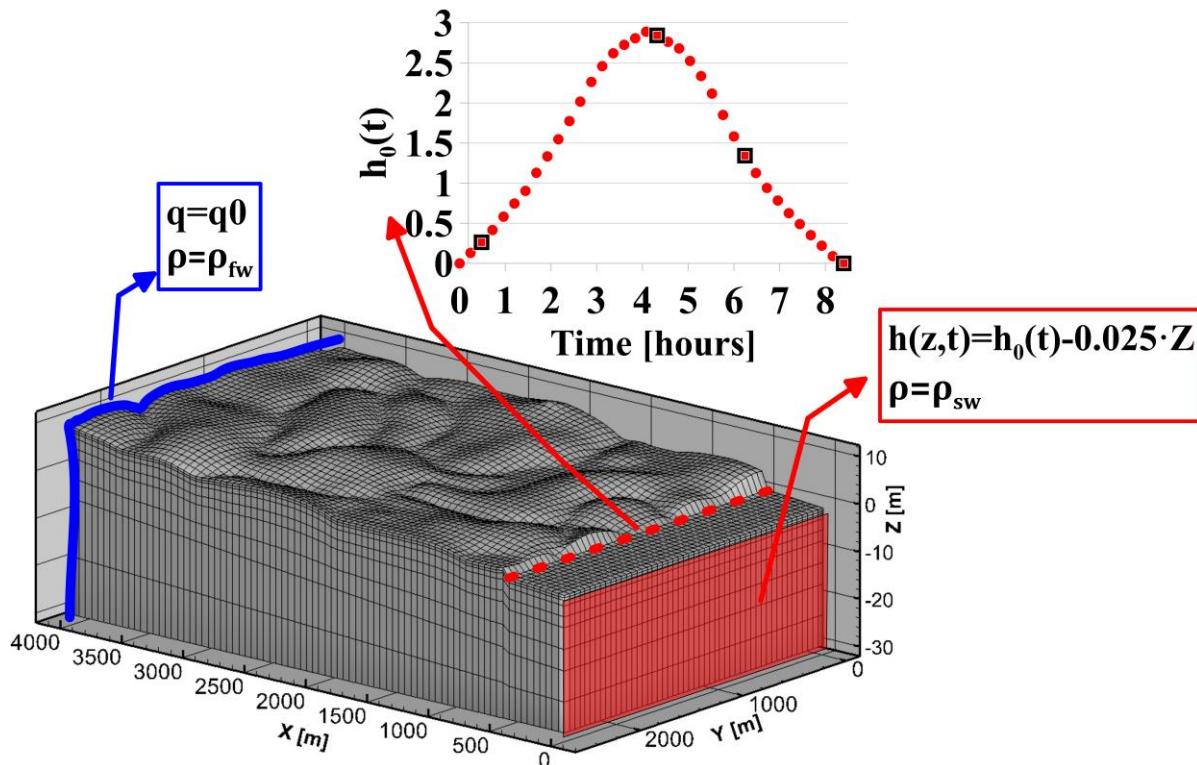
194 In Equation 8, i_z is the actual simulated or observed vertical head gradient, defined as $i_z = -\frac{v_z}{K_z}$ (Eq. 4) and i_c is the theoretical
195 liquefaction threshold (Eq. 7). Thus, any point in space and time in which simulated SLF is close to 1 is potentially nearing
196 liquefaction. A layer in which SLF approaches 1 continuously from the surface to a depth Z_l is considered a "critical layer" of
197 thickness Z_l .

198 4 Hydrologic model

199 The effect of storm surges on groundwater flow is simulated using Hydrogeosphere (HGS) – a 3D numerical code that couples
200 surface and subsurface flow and solute transport (Therrien et al., 2010). For the surface flow, HGS solves the Saint-Venant
201 equations (also known as nonlinear shallow water equations), and for the variably saturated subsurface flow it solves the
202 Richards equation. The salt transport equation is solved in its advective-dispersive form, and the variable-density flow solution
203 is coupled to the transport solution through a linear equation of state. Hydrogeosphere has been successfully employed to
204 simulate storm surges in several recent studies (Guimond & Michael, 2020; Yang et al., 2013, 2018; Yu et al., 2016), and here
205 it is applied to assess the risk for sediment liquefaction and erosion from surge-induced pore water head gradients. This is a
206 novel interdisciplinary approach, applying a robust 3D hydrologic model in the context of coastal geomechanics.



207 The model domain (Figure 3) is 4000 m (cross-shore, X) by 2500 m (alongshore, Y), extending to a depth of 30 m below the
208 mean sea level ($Z=0$). The terrestrial extent of the domain is 3550 m ($450 < X \leq 4000$), with the ocean spanning $0 \leq X \leq 450$ (Figure
209 3). The elevation at the ocean side boundary is $Z(X=0)=-1$, so the seafloor slope is $1/450 \approx 0.0022$. This slope is representative
210 of U.S. Atlantic and Gulf coastal systems averaged over large cross-shore distances (e.g., from the beach to the mid continental
211 shelf). Although local slopes in the surf and beach often are much steeper than those used here, this study is focused on the
212 liquefaction in and near the inundated dune system. The average surface elevation inland ($X=4000$ m) is 5 m, so that the
213 average land surface slope is $5/3550 \approx 0.0014$. Thus, there is a change in average slope at the coastline, as the offshore portion
214 is steeper (~ 0.0022) than the onshore (0.0014), as in many coastal areas. A simulation with a -0.5 m sea level (i.e., still water
215 shoreline at $X=225$ m) indicates that critical vertical hydraulic gradients occur near this change in overall slope irrespective of
216 the shoreline location (Figure A1 in the Appendices). A simulation with a larger beach slope ($Z(X=0) = -6$; $slope = 6/450 = 0.0130$) resulted in similar vertical hydraulic gradients as the baseline slope (0.0022) (Figure A2 in the
217 Appendices), indicating that although the baseline slope is lower than typical, the analysis based on it is also valid for steeper
218 slopes. The domain of the finite difference model consists of 44,000 rectangular cells, where the cell sizes in the X and Y
219 direction are 25 and 50 m, respectively. The cell size in the Z direction varies from 8 m in the bottom of the domain to about
220 0.5 m in the top 2 m to balance between computation time and the resolution necessary to resolve the dynamics close to the
221 surface (Figure 3). The homogenous hydraulic conductivity K_x is 50 m/d for the baseline simulation and K_x varied between
222 10 and 100 m/d in sensitivity analyses. In all simulations, the anisotropy was 10 (i.e., the vertical hydraulic conductivity, K_z ,
223 was 10 times lower than the horizontal hydraulic conductivity, K_x). This range of hydraulic conductivity with a porosity, n , of
224 0.3 is typical for sandy beach environments (Freeze and Cherry, 1979). Although a change in K could be associated with a
225 change in n for some sediments and mixtures, due to the potentially complex relationships between porosity and the sediment
226 textural properties, including grain size distributions, shapes, and K , the porosity was kept constant in the simulations presented
227 here.



229

230 **Figure 3:** Hydrogeosphere model domain as a function of the vertical Z , cross-shore X , and alongshore Y dimensions, boundary
 231 conditions (red and blue boxes), and the surge height evolution curve (inset). The blue curve is the terrestrial freshwater recharge
 232 boundary, the red rectangle is where a fixed seawater head and concentration are applied to the subsurface domain, and the red
 233 dashed line is where the sea level height boundary condition ($h_0(t)$) is applied on the surface domain. For the steady-state
 234 simulations $h_0(t)=0$, and for the transient surge simulations the curve in the inset is applied. The black squares in the inset mark the
 235 times plotted in Figure 5.

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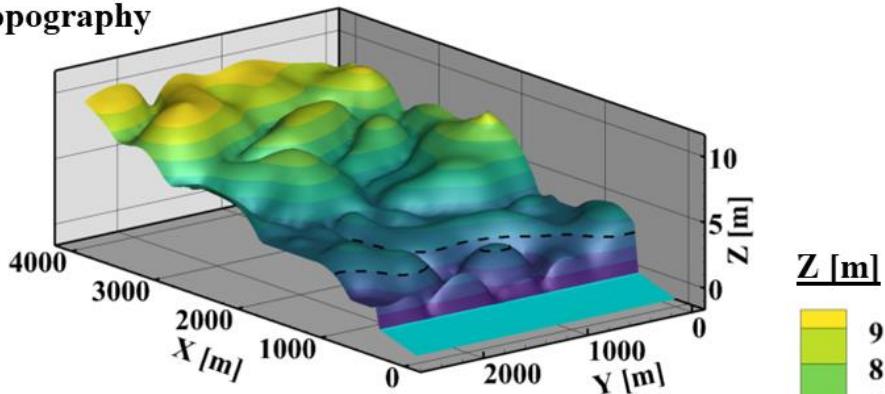
237 The boundary conditions in the simulations were applied in two stages – a steady-state period and a transient surge period. For
 238 the steady-state simulations, terrestrial boundary conditions of constant freshwater specific recharge ($q=q_0, \rho=\rho_{fw}$) were
 239 applied on the vertical wall at the inland edge of the subsurface domain at $X=4000$ (blue curve in Figure 3) (Ataie-Ashtiani et
 240 al., 2013; Yang et al., 2018; Yu et al., 2016). The opposite edge of the domain at $X=0$ (red wall in Figure 3) was a typical sea
 241 boundary condition with depth-dependent head and saline ocean water ($h=-0.025 \cdot Z; \rho=\rho_{sw}$). On the surface domain the only
 242 boundary condition is applied on the coastline $X=450$ m, red dashed line in Figure 3) as a fixed, time-dependent head ($h=h_0$
 243 (t)) and seawater density ($\rho=\rho_{sw}$). The applied head on the coastline was held at zero through the steady-state simulations.
 244 For the transient surge simulations, the coastline head was varied over 8.5 hours between zero and a 3 m maximum surge
 245 height (inset in Figure 3). A sea level of 3 m above the mean represents a combined high-tide and surge event with a projected



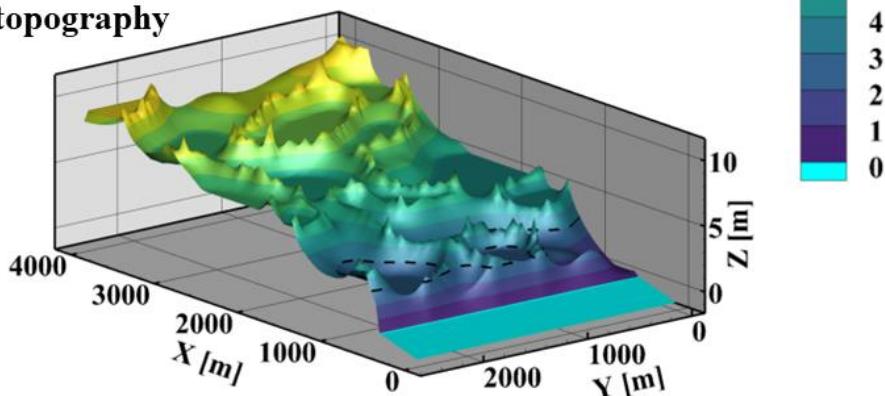
246 return period of 100 yr by the year 2050 in the East Coast of the United States (Tebaldi et al., 2012). The ocean surface was
247 assumed to be spatially constant at any time, and effects of wind waves were not simulated.
248 The sensitivity of the results to the topography and hydrogeologic parameters was tested, including freshwater influx ($0.01 < q_0 < 0.04$ m/d, Figure 3 and Table 1) and hydraulic conductivity ($10 < K_x < 100$ m/d, Table 1, typical values for sandy
249 beaches (Freeze & Cherry, 1979)). For the baseline hydraulic conductivity ($K_x=50$ m/d) the range of overall (land-to-sea)
250 hydraulic gradients, calculated as q_0/K_x , was 0.0002 and 0.0008, on the lower side of typical coastal settings (roughly
251 around 0.0010), and so the calculated hydraulic gradients in the current analysis are considered a conservative estimate. Two
252 topographies (Figure 4) (Yu et al., 2016) were generated with ARCMAP 10.0 Geographic Information System (GIS) software
253 (ESRI, 2011), using multigaussian random fields that were transformed (Zinn & Harvey, 2003) to connect either topographic
254 highs or lows rather than the median topographic values as in the non-transformed multigaussian fields. The first topography,
255 named “River” (Figure 4a), is characterized by surface depressions that connect to the sea. The topographic lows are connected,
256 forming “river”-like patterns in the surface morphology), superimposed on the background slope of 0.0014. The second
257 topography, “Crater” (Figure 4b), features connected crests surrounding disconnected surface depressions, such that the highs
258 are connected, forming “crater” like shapes. The two topographies do not mirror each other (Figure 4), but represent reverse
259 alongshore trends near the shoreline ($450 < X < 500$ m) in which the area around $0 < Y < 300$ m ($2200 < Y < 2500$ m) is the highest
260 (lowest) for the River topography and lowest (highest) for the Crater topography. Comparisons with real topographies of the
261 Delaware coastal plains (Yu et al. 2016) suggested that the River topography best represents real-world meso-topography.
262 However, the Crater topography provides important insights to how meso-topography controls the evolution of head gradients
263 during storm surges. In extreme flooding events (e.g., tsunami), large-scale changes in surface morphology (e.g., landslides)
264 may alter the pore-pressure distribution. These effects were excluded from the current work, as the simulated surface was
265 considered constant throughout the simulation. Additionally, soil deformation and the resultant stress re-distribution were not
266 considered in this model, as the hydrologic model (HGS) assumes constant porosity.



a. River topography



b. Crater topography



268

269 **Figure 4: (a) River and (b) Crater topographies as a function of the vertical Z, cross-shore X, and alongshore Y coordinates. Light**
270 **blue is the offshore bathymetry, and the coastline is at X=450 m. The overall slope accounting for macro-topography is the same for**
271 **both topographies, the average elevation at X=4000 m is ~5 m, making it a slope of 5/3550≈0.0014. The dashed black curve marks**
272 **the Z=3 m contour, which is equal to the maximum surge-induced sea level (hmax).**

273

274

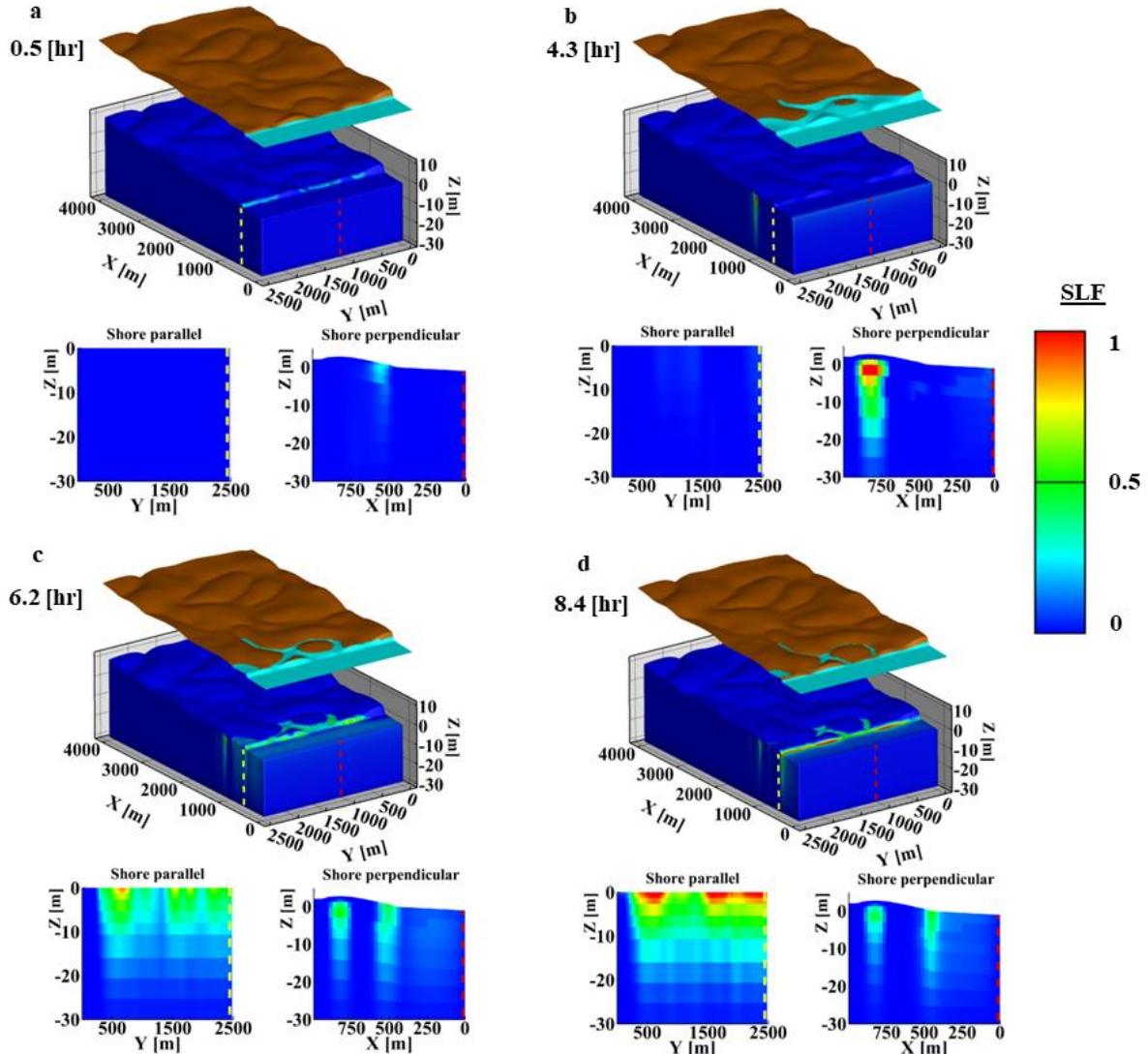
275 For each simulation, the vertical hydraulic gradients (i_z in Equation 8) are calculated over a vertical slice along the coastline,
276 i.e., the plane defined by $X=450$, and normalized by the threshold defined by Equation 7 (i_c) to calculate the SLF (Equation
277 8). As explained in Section 3 above, values of SLF that approach 1 are considered critical for liquefaction. When $SLF \ll 1$ the
278 simulated surface theoretically is stable. Only upward, destabilizing velocities (exfiltration) are considered, and so negative
279 velocities were assigned a value of $i_z=0$.

280



281 5 Results

282 The baseline case ('River' topography with $q_0=0.02 \text{ m/d}$; $K_z=5 \text{ m/d}$) includes a 3 m surge and simulates the resultant
283 changes in head gradients (Figure 5). During the inundation stage when sea level is increasing, the head gradients increase
284 landward in front of the moving surge, and in the flooded zone there is infiltration (head decreases downward, $\nabla h > 0$). After
285 the peak of the inundation, when the high-water levels begin to recede, downward gradients (i.e., head increases downward,
286 potentially destabilizing) develop underneath the still-water shoreline ($X=450 \text{ m}$). These downward gradients increase in
287 magnitude as the water level recedes, and the subsurface system relaxes back to background levels (not shown in Figure 5)
288 within ~50 days for the high-K aquifers to ~500 days for the low-K aquifers, similar to prior simulations of storm impacts
289 (Robinson et al. 2014). The peak alongshore variation of the vertical hydraulic gradients occurs at the end of the inundation
290 ($t=8.4 \text{ hr}$, Figure 5d). The vertical hydraulic gradients onshore of the inundation front during run-up (Figure 5b) develop in
291 subaerial areas, and therefore the calculated SLF for these zones is based on the saturated unit weight ($\gamma_{\text{sat}}=\gamma_{\text{sub}}+\gamma_{\text{fw}}$) of
292 sediments rather than the submerged unit weight (γ_{sub} , Equation 3), and the model-predicted liquefaction may not occur in
293 real systems because saturated soils are more stable than submerged ones (Briaud, 2013).



294

295 **Figure 5:** Surface inundation and vertical hydraulic gradients at (a) 0.5, (b) 4.3, (c) 6.2, and (d) 8.4 hr after the simulated surge
 296 begins (for the surge height at these times refer to Figure 3). In each panel, the surface domain is shown on top, the subsurface 3D
 297 domain and vertical gradients are shown below, and two cross sections through the subsurface are shown: shore-parallel (left in
 298 each panel) and shore-perpendicular (right). The locations of the sections are shown on the 3D plot as red dashed lines (for shore
 299 perpendicular) and yellow dashed lines (for shore parallel). The upper two panels are during the run-up stage and the lower are
 300 during the retreat stage. Refer to Figure 3 for the surge height at each time shown here. Note that downward gradients (head
 301 increases downward) are plotted as positive and upward gradients (head increases upward) are plotted as zero.

302

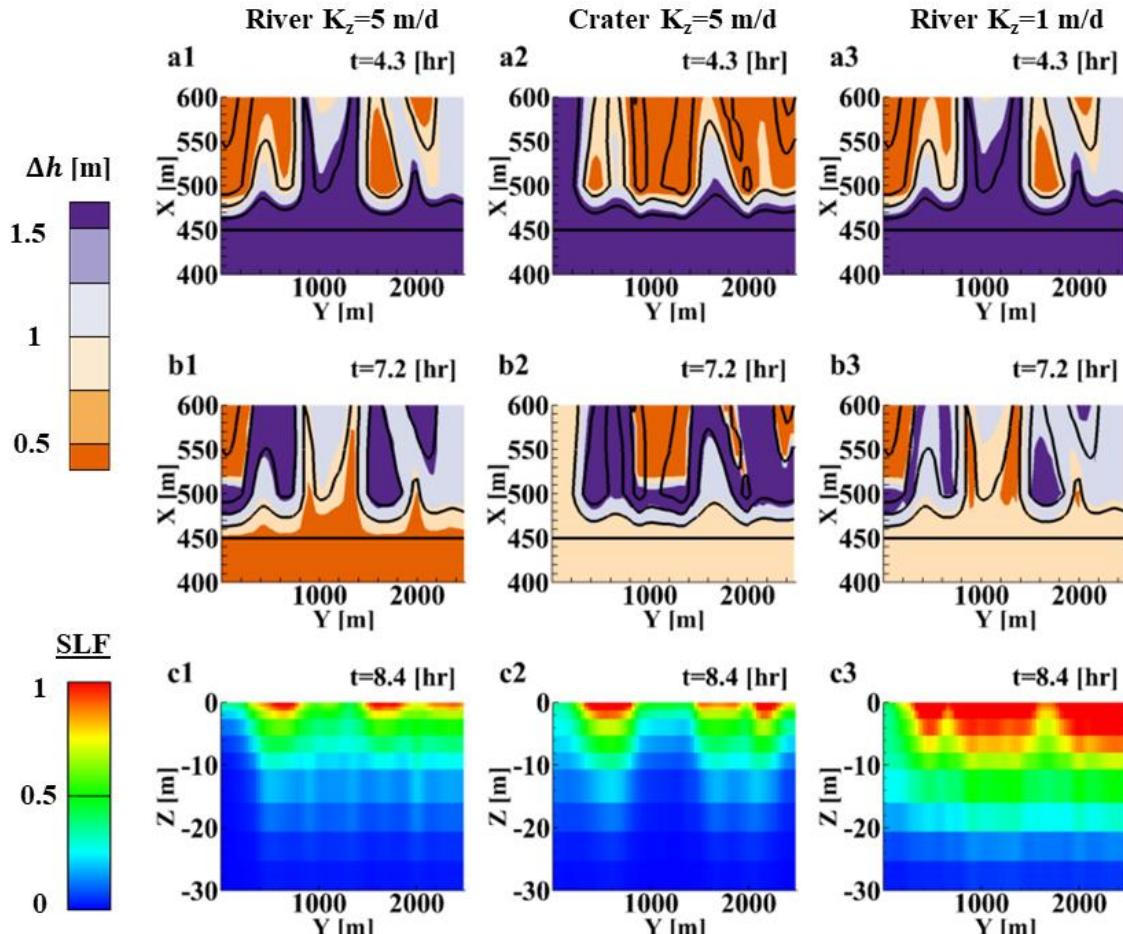
303 The head changes (Δh in Figure 6) between the steady state and the peak of the inundation inversely follow the topography
 304 (black contours in Figure 6a and b). For the highest topographic elements (Y=0 m for the “River” and Y=2500 m for the



305 “Crater”), which are not inundated, the simulated heads are approximately equal to the maximum ocean level at the dune crest
306 ($X \sim 460$ m), and decay inland over ~ 100 m, roughly consistent with field observations (Figure 1). The maximum head changes
307 (purple colors in Figure 6a) inland of the shoreline ($X > 475$ m) at peak surge occur in the inundated topographic lows. Toward
308 the end of the simulated surge ($t=7.2$ hr, Figure 6b) the surge-induced overpressures are released in the topographic lows (low
309 values of Δh in Figures 6b). The head differences also are low in the topographic highs because the heads there did not rise
310 significantly during inundation. In contrast, the intermediate topographic features show high head differences (dark purple in
311 Figure 6b). The lowest near-shore ($450 \leq X \leq 500$ m, $900 \leq Y \leq 1200$ m) topography undergoes similar head changes during the
312 peak surge for high and low K (compare Figure 6a1 with 6a3). However, in the low K case (Figure 6a3, 6b3), the heads are
313 not released effectively as the surge recedes, and significant excess heads of ~ 1 m difference remain near the end of the surge
314 (compare Figure 6b3 with 6b1 for $X \sim 450$ m).

315 When the surge has retreated ($t=8.4$ hr), the head gradients at the dune toe (initial shoreline) ($X = 450$ m) reach their maximum
316 (Figure 6c1-c3). In all simulations critical gradients ($SLF \rightarrow 1$, red zones in Figure 6 c1-c3) are simulated at some locations
317 below the shoreline, supporting the findings of several recent field studies in which momentary liquefaction was observed in
318 response to inundation events (Sous et al., 2016; Yeh & Mason, 2014). The alongshore distribution of the surge-induced
319 gradients is insensitive to the freshwater influx (q_0), even though the antecedent local hydraulic gradients differed by up to a
320 factor of 4 between simulations (Figure A3 in the Appendices, note that the values of the antecedent local gradients are about
321 an order of magnitude lower than the peak gradients). The depth and alongshore locations of the areas prone to liquefaction
322 (i.e., $SLF \sim 1$) are sensitive to the topography (compare Figures 6 a1,b1,c1 with a2, b2, and c2) and the hydraulic conductivity
323 (compare Figures 6 a1,b1,c1 with a3, b3, and c3). The two topographies exhibit a similar spatial pattern of SLF (Figure 6c1
324 and c2) even though the differences in topography (Figure 4) cause significant differences in the surge-induced head changes
325 (Figure 6 a1 and a2). For example, the area to the left of the domain ($Y \leq \sim 300$ m) is a topographic low in the Crater topography
326 and undergoes significant head changes at the peak of the inundation (Figure 6a2), whereas for the River topography there is
327 a topographic high for $Y \leq \sim 300$ m, which is not as strongly affected by the surge (Figure 6a1). However, in both cases this area
328 is where the least significant vertical head gradients develop (Figure 6c1 and c2).

329 The hydraulic conductivity has a significant effect on the simulated surge-induced gradients (Figure A4 in the Appendices).
330 Decreased hydraulic conductivity causes higher peak vertical gradients and changes the spatial (shore-parallel) distribution of
331 the gradients (compare Figure 6c3 with 6c1, especially near $Y = 1000$ m, and also see Figure A4). Furthermore, decreasing
332 hydraulic conductivity alters the depth Z_1 of “critical layers” with $SLF = 1$ (Equation 8) (compare Figure 6c3 with 6c1). In
333 the high-K simulations (Figure 6c1 and c2), the depth Z_1 of these “critical layers” with $SLF \sim 1$ ranges between 0 and 2.5 m,
334 and in the low-K simulation (Figure 6c3) Z_1 is up to ~ 5 m.



335

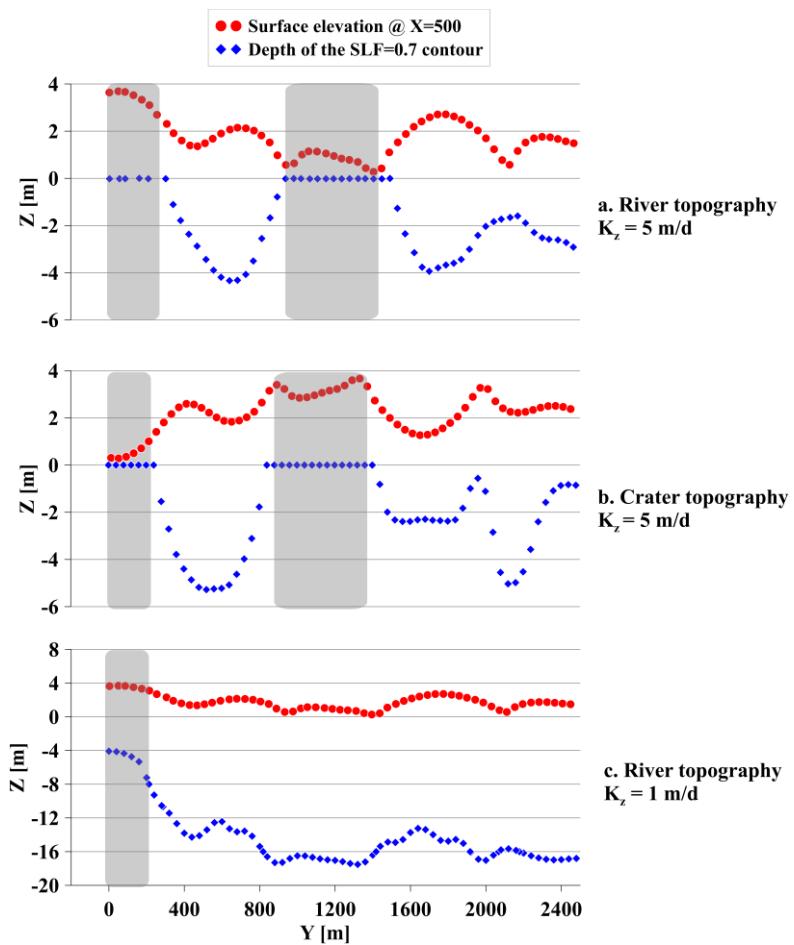
336 **Figure 6:** Top row (a1-a3): maps of the maximum near-surface head differences between those at the peak of the inundation and the
 337 initial, pre-surge values (denoted Δh_1) as a function of the cross-shore X and alongshore Y coordinate. Middle (b1-b3): maps of
 338 the maximum subsurface head differences between those near the end of the surge ($t = 7.2$ hr, Figure 3) and the initial, pre-surge
 339 heads (denoted Δh_2) as a function of X and Y. Bottom (c1-c3): Liquefaction potential SLF at the shoreline, $X = 450$ m, as a function
 340 of the vertical Z and alongshore Y coordinate. These 3 metrics are plotted for River topography with $K_z=5$ m/d (left, a1-a3), Crater
 341 topography with $K_z=5$ m/d (center, a2-c2) and River topography with $K_z=1$ m/d (right, a3-c3). In the upper and middle panels (map
 342 views a1-a3 and b1-b3) the black contours are surface elevation with 1 m intervals. The horizontal line at $X=450$ is the coastline
 343 ($Z=0$). The lower panels are plotted for $t=8.4$ hr, the time at which the vertical gradients peaked in all simulations all along the
 344 coastline.

345

346 The relationship between coastal topography and the surge-induced liquefaction potential is evident when comparing the
 347 surface elevations 50 m landward of the coastline ($X=500$ m) and the peak vertical gradients below the coastline for different
 348 topographies and K's (Figure 7). Here, the $SLF=0.7$ contour is used for statistical stability (there are more locations with
 349 $SLF \geq 0.7$ than with $SLF=1$). For both topographies, when K is high, SLF typically remains less than 0.7 (in Figure 7 where the



350 blue diamonds = 0) at the shoreline adjacent to the highest ($Z > 3\text{m}$) and lowest ($Z < 1\text{ m}$) topographic elements (marked by
 351 gray rectangles in Figures 7a and b), suggesting the intermediate topographic features may lead to the strongest vertical
 352 hydraulic gradients and liquefaction potential. However, the height of intermediate features that produce high gradients may
 353 be dependent on the site and hydrogeological parameters. For example, in the two simulations with higher K_z , 1-3 m
 354 topographic features are associated with most of the significant surge-induced gradients (Figure 7a and b). For the lower K_z
 355 case, significant gradients occur also below the lowest area (Figure 7c), and only the highest area that is not inundated does
 356 not develop significant gradients (gray rectangle in Figure 7c).



357

358 **Figure 7: Topographic elevation at $X=500\text{ m}$ (50 m onshore of the shoreline, red circles) and depth of the $SLF=0.7$ contour below**
 359 **the shoreline (blue diamonds) versus alongshore coordinate Y for (a) the River topography with $K_z=5 \text{ m/d}$,**
 360 **(b) Crater topography with $K_z=5 \text{ m/d}$, and (c) River topography with $K_z=1 \text{ m/d}$. Deeper locations of the $SLF=0.7$ contour (blue diamonds) mean thicker**
 361 **"critical layers". The places where no significant critical layer develops (i.e., the elevation of the $SLF=0.7$ contour is $Z=0$) are marked**
 362 **by gray rectangles.**



363 6 Discussion

364 6.1 Alongshore variability

365 The simulations suggest that alongshore variability of the magnitudes of the vertical gradients is strongly associated with the
366 coastal topography (Figures 5-7). To induce high gradients and deep critical layers when pressures are released, it is necessary
367 to have inundation resulting in high infiltration and increased heads. Thus, topographic highs that are not inundated cannot
368 develop high gradients (Figures 6 and 7). Meanwhile, overpressures often are released efficiently from inundated areas as the
369 surge recedes. Topographic elements that are low enough to be inundated, but are also high enough to limit the post-surge
370 exfiltration may prevent release of pressures, possibly explaining the correlation of liquefaction potential with intermediate
371 topographic features (1-3 m high for a 3 m surge). This explanation would suggest that the characteristic elevation of
372 “intermediate features” would scale with the surge magnitude. Pressure releases also can be limited by low hydraulic
373 conductivity. Thus, the simulations suggest the areas most susceptible to destabilization (i.e., deep critical layers) are those
374 where topography is low enough to be inundated widely, and high enough that the pressure release is limited. The range of
375 susceptible topographic elements depends on hydraulic conductivity, which also has a sweet spot of vulnerability: A simulation
376 with even lower hydraulic conductivity ($K_z=0.05$) showed that very low values of K limit the surge-induced infiltration and
377 thus critical gradients develop only to a limited vertical extent and the alongshore variability (i.e., the dependency on onshore
378 topography) diminishes (Figure A5 in the Appendices).

379 6.2 Cross-shore spatiotemporal variability

380 During the flooding stage, negative vertical gradients (infiltration) that do not promote sediment instability occur at and
381 seaward of the moving inundation front. Positive vertical gradients occur landward of the front (top right panel in Figure 5)
382 owing to alteration of the pre-existing steady-state flow field (Figure 2) by the advancing overpressures from the surge.
383 However, the simulated values of $SLF=1$ inland of the inundation front are not necessarily sufficient to liquefy the surface,
384 because the actual weight of the unsubmerged soil is greater than γ_{sub} (Equation 2). Nevertheless, the liquefaction potential
385 calculated here may still represent an underestimate, as Mory et al. (2007) showed that as little as 6% air content in the pores
386 may reduce the required pressure difference to liquefy the sediment by 0.01 m. This highlights the need to consider air contents
387 in future studies. Furthermore, these inland processes, and the potential for liquefaction in these areas, may be affected by
388 vegetation, trapping of gases, hysteresis of wetting and drying, and other processes that have not been considered here.
389 Nevertheless, the presented approach demonstrates the feasibility and a pathway to implement the concept of surge-induced
390 momentary liquefaction in a hydrological model that can predict variable-density groundwater flow in coastal and estuarine
391 environments.

392 The receding water levels after the peak of the surge allow fast release of the elevated heads that developed in the inundated
393 area, because the overlying burden of surge waters is removed abruptly. For all simulations at all alongshore locations, the
394 positive head gradients simultaneously reached a maximum when the water had receded completely ($t=8.4$ hr) and all the



395 inundation water overburden was released. The rate of head release determines the hydraulic gradients that occur in the soil
396 material, so that faster release of the overpressures produces lower positive head gradients. As the water recedes, the highest
397 release rates, and thus overpressures, develop under the beach area, where the slope changes from a terrestrial average slope
398 of 0.0014 to the seafloor slope of ~0.0022 (Figure 3). Thus, the simulations suggest the highest surge-induced gradients might
399 be expected under convex topography, for example near the berm or near a scarp in the beach face.

400 **6.3 Implications for coastal engineering**

401 Most previous studies of extreme wave-induced pressurization in coastal environments focus on cross-shore variability (Sous
402 et al., 2013, 2016; Turner et al., 2016; Yeh & Mason, 2014). Here, it is shown that under realistic hydrogeological conditions
403 (surge height, topography, groundwater flow regime – all based on values that are commonly observed in natural systems)
404 with alongshore varying topography there can be significant differences in storm-induced maximum hydraulic gradients and
405 in the depths of corresponding critical layers over small distances along the coastline (<500 m) (Figure 6). The simulations
406 suggest that beach and dune morphology are important factors determining the spatial variability of high gradients. Although
407 low-lying coastal areas may endure the greatest flooding, the largest hydraulic gradients and the deepest liquefaction layers
408 may occur at the toes of the intermediate-scale (1-3 m high for a 3 m surge) topographic features. While discussing practical
409 implications of the present analysis, it is important to remember that, as noted above, the model adopted here is a hydrological
410 model that does not explicitly simulate the soil dynamics and the surface and subsurface domains were assumed constant with
411 time through the simulations. This assumption overlooks other dynamic controls on the development of stresses, such as soil
412 deformation and surface erosion. Moreover, the analysis presented here isolates the vertical seepage component to calculate
413 the potential for momentary soil liquefaction. In a 3D framework, horizontal seepage components likely come into play and
414 other failure mechanisms, such as shear failure, are likely too (Zen et al., 1998). However, for the conclusions drawn here
415 regarding the spatio-temporal distributions of surge-induced gradients, the hydrologic modeling provides an important tool to
416 study the hydrogeological aspect of the problem. The model could be further expanded to include other components in future
417 work.

418 **7. Conclusions**

419 Field measurements from Duck, North Carolina, show that during Hurricane Joaquin the groundwater flow regime at the ocean
420 side was impacted substantially, and the hydraulic head gradient reversed its direction, followed by a period of recovery during
421 which downward gradients (upward fluxes) were regenerated. This suggests that hydraulic gradients generated by storm surges
422 may substantially affect the stability of beach surfaces. We explored this idea and its generality by harnessing a robust
423 hydrological model to simulate a generalized coastal system and found that in the nearshore area, surge-induced hydraulic
424 gradients may peak to critical levels that could potentially induce sand liquefaction. The locations where these critical gradients
425 occur are transient and depend on the beach morphology and hydraulic conductivity. Both the elevation of topographic features

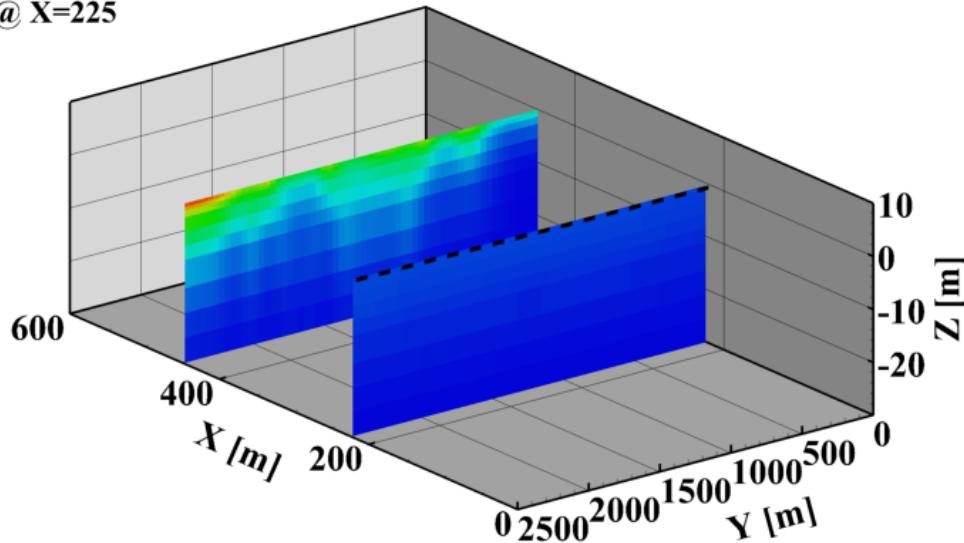


426 and their permeability are important factors in promoting liquefaction. Elevations must be low enough to become inundated,
427 and high enough to retain elevated heads needed to build critical gradients. Similarly, hydraulic conductivity must be high
428 enough to allow floodwater to infiltrate, but low enough that water is not drained immediately such that critical gradients can
429 persist. This alongshore variability has not been observed in field measurements because the common approach in field studies
430 is to measure the cross-shore variability of hydraulic heads during storms. Importantly, this work presents a novel approach to
431 bridge the gap between coastal hydrology and coastal engineering, incorporating robust hydrogeological modeling in a
432 geotechnical framework.

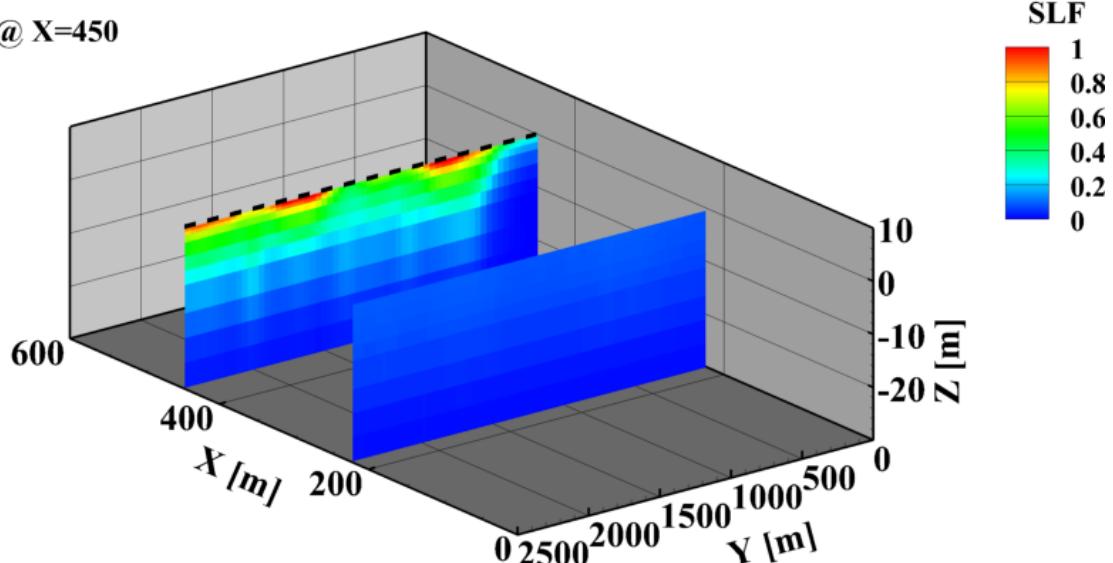


433 8. Appendices

a. Coastline @ X=225



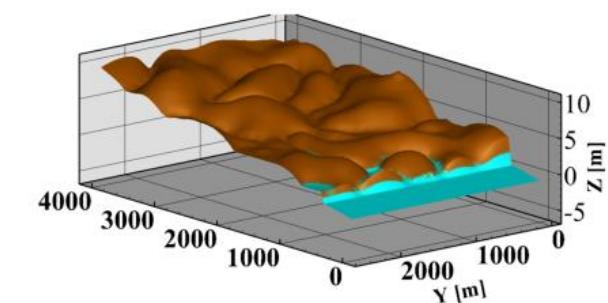
b. Coastline @ X=450



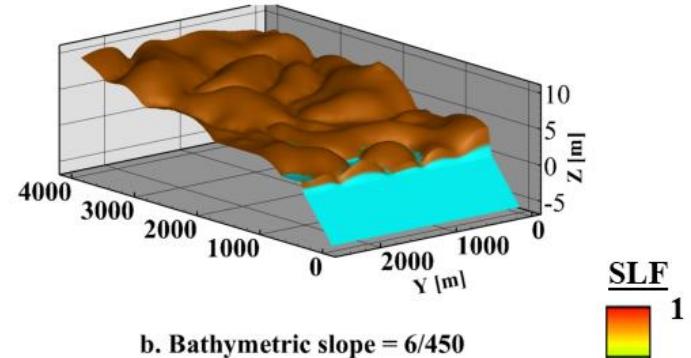
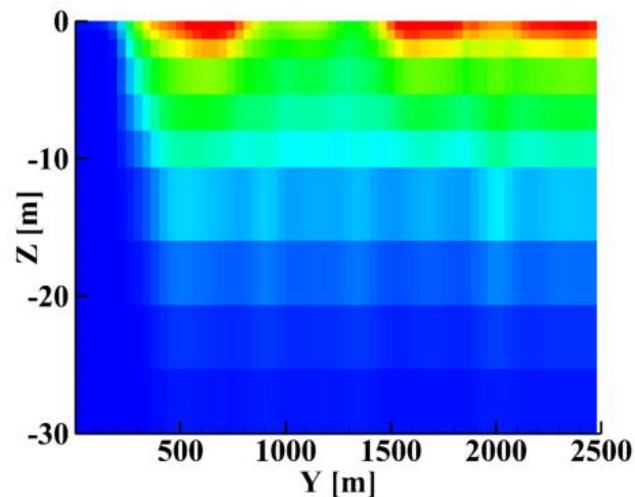
434

435 Figure A1: Contours (color scale on the right) of peak SLF ($t=8.4 \text{ hr}$) as a function of the vertical Z , cross-shore X , and
436 alongshore Y coordinate for (a) a simulation with the coastline at -0.5 m ($X = 225 \text{ m}$) and (b) a simulation with the
437 coastline at 0 m ($X = 450 \text{ m}$). The dashed black lines mark the coastline in each respective simulation. The slice with
438 high SLF values in (a) is not underneath the simulated coastline.

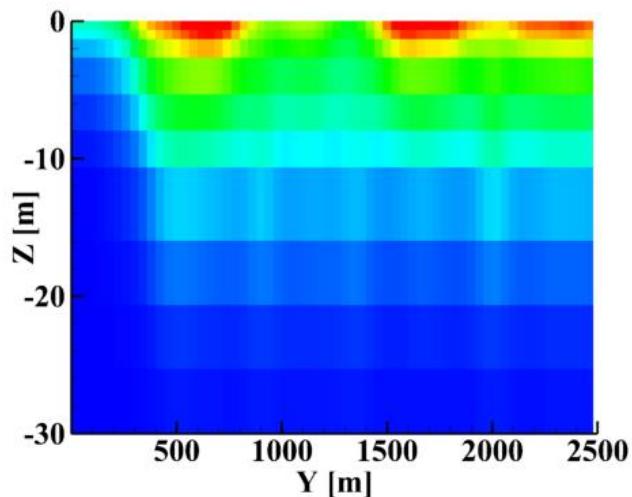
439



a. Bathymetric slope = 1/450



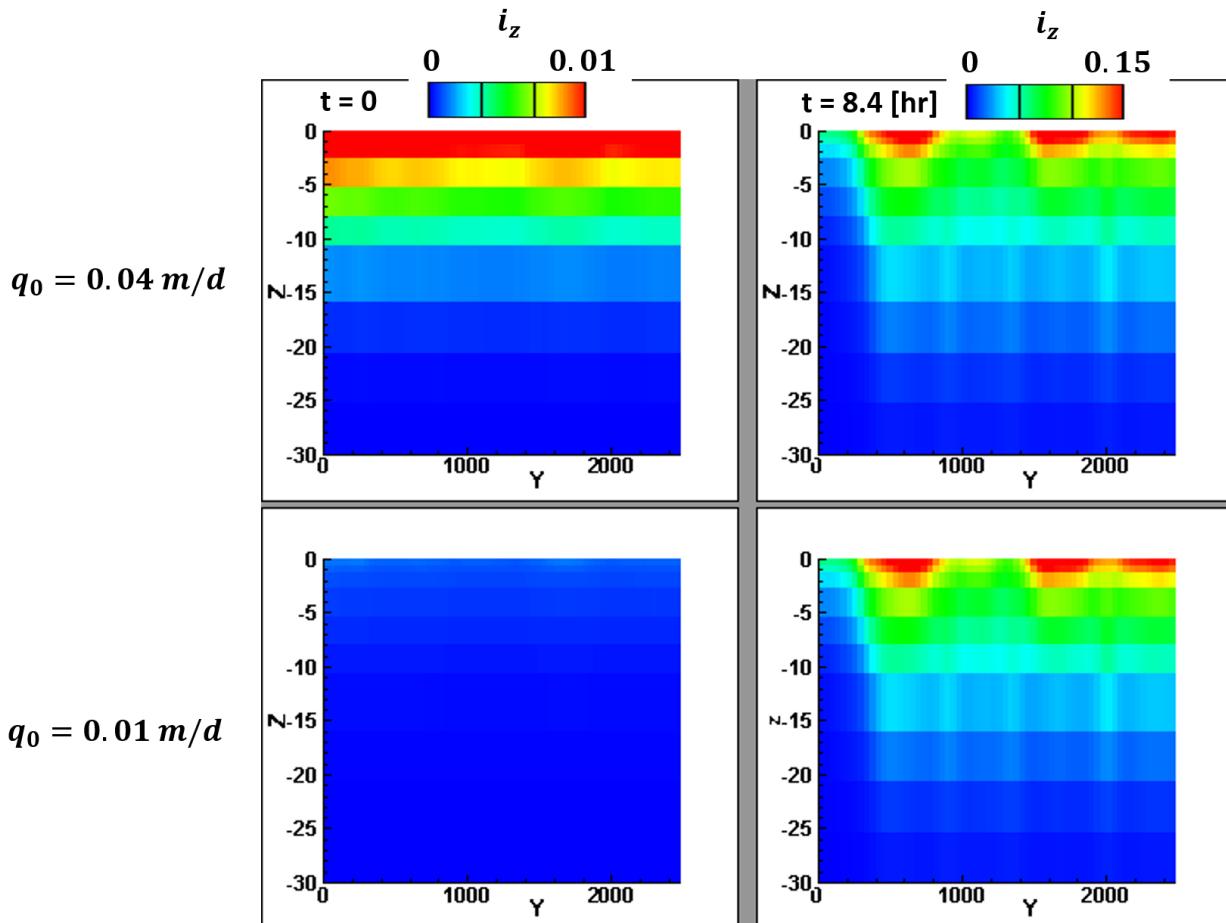
b. Bathymetric slope = 6/450



440

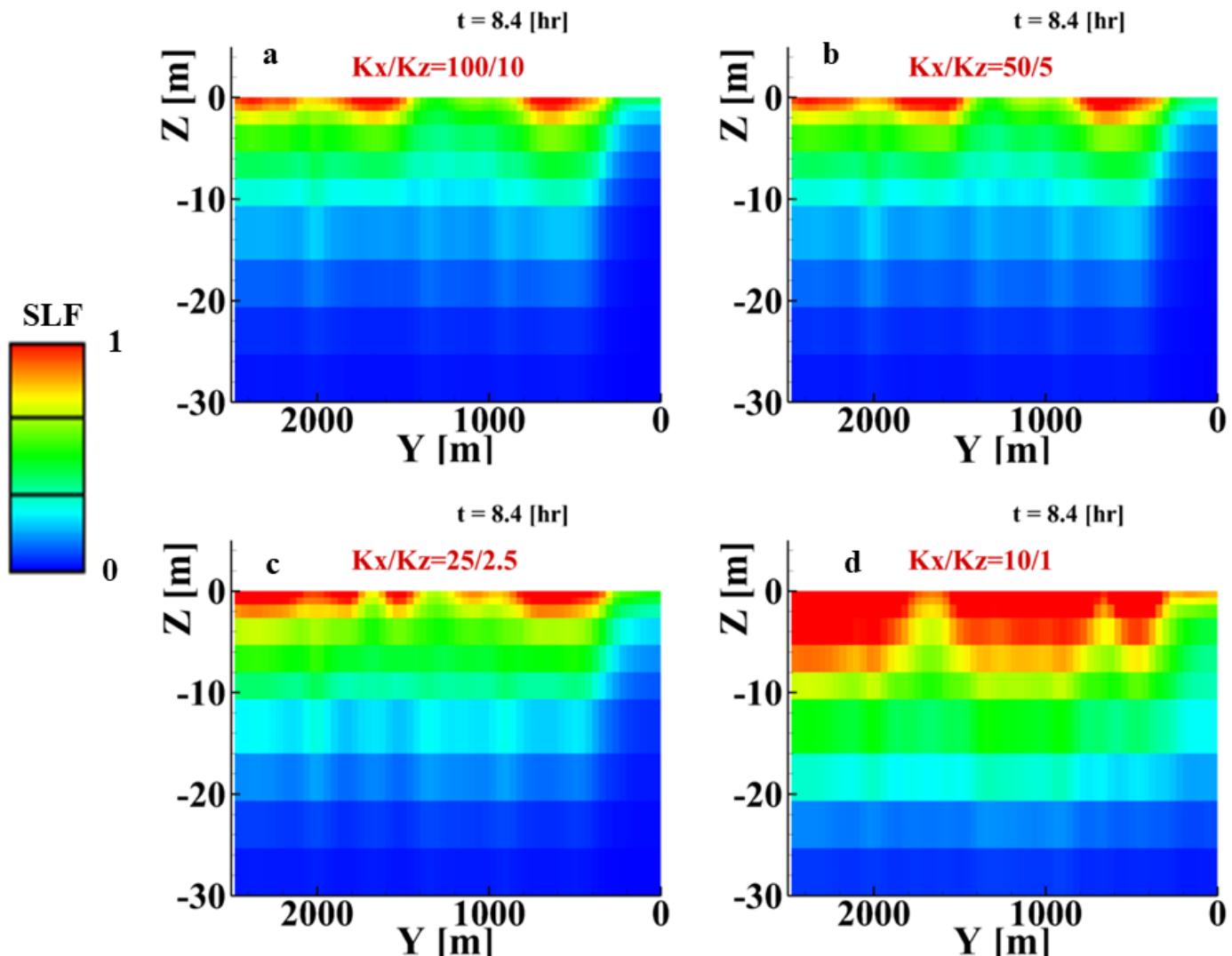
441 **Figure A2:** Contours (color scale on the right) of peak SLF ($t = 8.4 \text{ hr}$) for a simulation with (a) bathymetric slope of
442 $\frac{1}{450} \approx 0.002$ and (b) a simulation with a higher bathymetric slope ($\frac{6}{450} \approx 0.013$). The upper part of each panel shows
443 the surface with the inundation water and the lower part is the vertical slice with the SLF values below the coastline
444 (X=450 m).

445



446

447 **Figure A3:** Contours (color scales on the top) of vertical hydraulic gradients (i_z) at $X = 450$ m (shoreline location) for
448 the pre-surge conditions (left) and the end of the surge when gradients are maximum (right) as a function of vertical Z
449 and alongshore Y coordinates. Note the different color scales between the pre-surge (left) and the peak (right) plots.
450



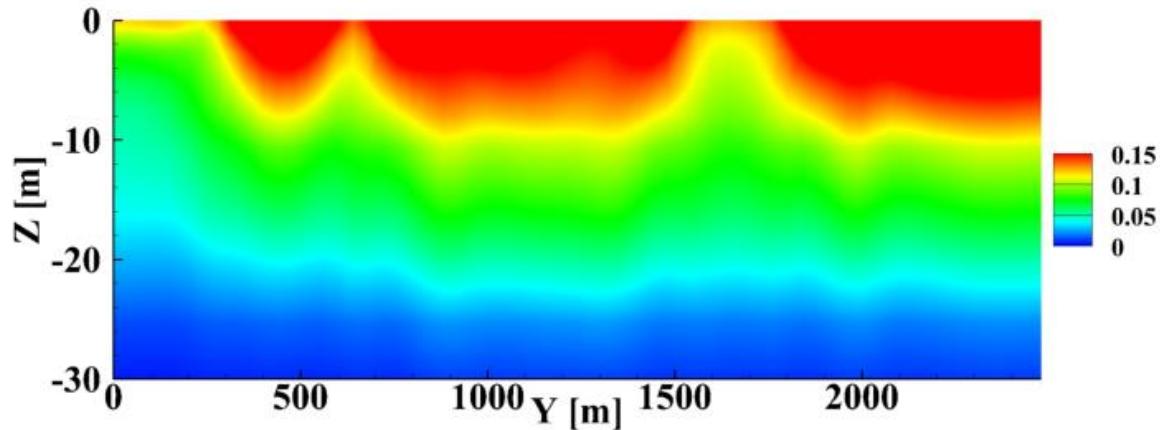
451

452 **Figure A4:** Contours (color scale on the left) of peak SLF ($t=8.4$ hr) vertical slices at the shoreline ($X = 450$ m) for K_x
453 and K_z of (a) 100 and 10, (b) 50 and 5, (c) 25 and 2.5, and (d) 10 and 1 m/d.

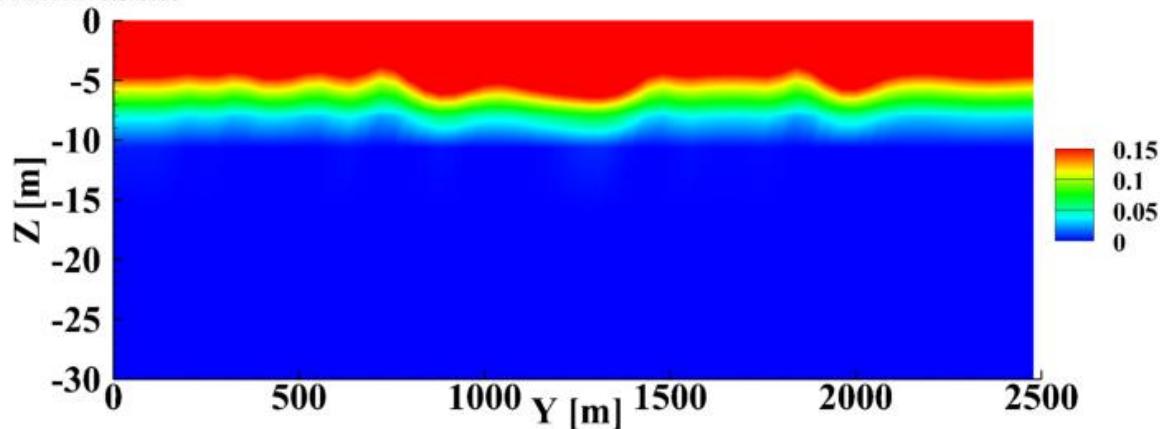
454



a. $K_x/K_z=10/1$



b. $K_x/K_z=0.5/0.05$



455

456 **Figure A5:** Contours (color scales on the right) of the maximum vertical hydraulic gradients (i_z) at $X = 450$ m (shoreline
457 location) for (a) $K_z = 1$ and (b) $K_z = 0.05$) as a function of vertical Z and alongshore Y coordinates.

458 Author contribution

459 AP: conceptualization, investigation, visualization, formal analysis, writing (original draft); NS: conceptualization, formal
460 analysis, writing (review and editing), funding acquisition; MF: formal analysis, writing (review and editing); BR;
461 conceptualization, formal analysis, writing (review and editing), funding acquisition; SE: conceptualization, formal analysis,
462 writing (review and editing), funding acquisition; RH: Data curation, visualization, writing (review and editing); RF formal
463 analysis, methodology; HM: conceptualization, formal analysis, writing (review and editing), supervision, funding acquisition,
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