



Meteor Radar vertical wind observation biases and mathematical debiasing strategies including a 3DVAR+DIV algorithm

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Abstract. Meteor radars have become a widely used instrument to study atmospheric dynamics, in particular in the 70 to 110 km altitude region. These systems have been proven to provide reliable and continuous measurements of horizontal winds in the mesosphere and lower thermosphere. Recently, there have been many attempts to utilize specular/transverse scatter 5 meteor measurements to estimate vertical winds and vertical wind variability. In this study we investigate potential biases in vertical wind estimation that are intrinsic to the meteor radar observation geometry and scattering mechanism, and introduce a mathematical debiasing process to mitigate them. This process makes use of a spatio-temporal Laplace filter which is based on a generalized Tikhonov regularization. Vertical winds obtained from this retrieval algorithm are compared to UA-ICON model data. This comparison reveals a good agreement in the statistical moments of the vertical velocity distributions. Furthermore, 10 we present the first observational indications of a forward scatter wind bias. It appears to be caused by the scattering center's apparent motion along the meteor trajectory when the meteoric plasma column is drifted by the wind. The hypothesis is tested by a radiant mapping of two meteor showers. Finally, we introduce a new retrieval algorithm providing a physically and mathematically sound solution to derive vertical winds and wind variability from multistatic meteor radar networks such as the



Nordic Meteor Radar Cluster and the Chilean Observation Network De meteOr Radars (CONDOR). The new retrieval is called
15 3DVAR+DIV and includes additional diagnostic such as the horizontal divergence and relative vorticity to ensure a physically consistent solution for all 3D winds in spatially resolved domains. Based on this new algorithm we obtained vertical velocities in the range of $w = \pm 1 - 2$ m/s for most of the analyzed data during two years of collected data, which is consistent to the values reported from GCMs for this time scale and spatial resolution.

20 1 Introduction

Vertical wind in the mesosphere/lower thermosphere (MLT) is a key parameter because it is directly related to the vertical transport of momentum, energy, and constituents that drive the global meridional circulation, which is related to almost all dynamical processes in the global atmosphere (e.g. Smith et al., 2010; Qian et al., 2017; Guo and Liu, 2021). However, measuring vertical wind is one of the most challenging remote sensing tasks. The main reason is that the magnitude of long-term mean vertical wind is very small, often beyond the accuracy achievable with any instruments, while instantaneous, or short-duration vertical wind can be large but requires measurements at high temporal and spatial resolutions. Models predict vertical motions on seasonal time scales, at their typical horizontal grid resolution of about 100-200 km, on the order of 0.1 to a few cm/s, e.g., in the Kuehlungsborn Mechanistic Circulation Model (KMCM) and the Whole Atmosphere Community Circulation Model (WACCM) (Becker, 2012; Smith, 2012). At higher solutions, the models are able to resolve smaller scale 25 gravity waves and produce larger vertical winds. In Liu et al. (2014), the high resolution WACCM at 0.25° horizontal resolution produced vertical wind of 7-8 m/s in the lower thermosphere above a tropical cyclone. In a more recent study using the High-Altitude Mechanistic Circulation Model (HIAMCM) with a horizontal resolution of about 55 km, vertical wind velocities up to 3 m/s are reported at an altitude of about 80 km (Becker and Vadas, 2018). High resolution observations such as those made with a sodium lidar also measured vertical wind, showing that tidal perturbation in vertical wind can reach tens of cm/s 30 (Yuan et al., 2014). On the other hand, models and observations also indicate that the horizontal wind magnitudes at the MLT are typically one to two orders of magnitude larger (Miyoshi et al., 2017; McCormack et al., 2017; Borchert et al., 2019; Hocking et al., 1997; Batista et al., 2004; Hoffmann et al., 2007; Jacobi et al., 2009; Wilhelm et al., 2019; Stober et al., 2019). This large difference in the magnitudes between the horizontal and vertical wind component poses an additional challenge to 35 the observational methods, measurement analysis, and parameter estimation of vertical wind due to the requirement of clear separation between vertical and horizontal components.

During the past decades there have been many attempts to measure vertical wind velocities using High-Power-Large-Aperture radars such as EISCAT (Fritts et al., 1990; Hoppe and Fritts, 1995a, b). These EISCAT observations, with a temporal resolution of seconds, showed vertical velocities up to ± 10 m/s in the MLT and indicated the presence of a systematic vertical wind bias. Although the EISCAT campaign was conducted during the summer months using polar mesospheric summer echoes



45 as tracers, the mean vertical velocities showed a downward motion, which is contrary to what models suggest for this time of
the year. The systematic deviation was attributed to gravity wave motions interacting with the tracer. More recently, Gudadze
et al. (2019) presented vertical wind observations over two full summer seasons with the Middle Atmosphere Alomar Radar
System (Latteck et al., 2012) and confirmed the presence of a mean vertical wind bias and examined potential error sources
in the data analysis. Gudadze et al. (2019) concluded that the mean wind bias of a net downward motion in the center of the
50 PMSE layer can be explained by the sedimentation speed of the ice particles. Removing this sedimentation speed resulted in a
effectively zero wind speed or a very small upward motion in the order of a few cm/s.

In addition to these direct vertical wind observations using line of sight velocities, there are also indirect methods. For example, Vincent et al. (2019) derived mean vertical wind velocities by exploiting cross-calibrated MF-radar winds, and considering
55 the horizontal divergence between the pole and the latitude of the observations. This study reported the summer time mean
vertical motions of a few cm/s using measurements between 1994-2018. The magnitude and sign of these vertical winds were
in agreement to the values obtained by GCMs. Radiometers also offer an indirect methodology by measuring trace gases such
as water vapor or ozone (Schranz et al., 2019). Straub et al. (2012) estimated the vertical motion of air parcels from water
vapor observation during Sudden Stratospheric Warmings and obtained vertical velocities of a few mm/s at 70-80 km altitude.
Such trace gas observations are suitable for inferring vertical motions, which are too small to be observed by direct line of
60 sight measurements, which often do not reach a sufficient sensitivity to detect such small velocities within the instrument error
bounds.

Meteor radar observations have been widely used to measure horizontal winds and atmospheric waves (Hocking et al., 2001;
Holdsworth et al., 2004; Jacobi et al., 2007; Fritts et al., 2010b; Meek et al., 2013; Andrioli et al., 2013; Liu et al., 2013;
de Wit et al., 2014). Horizontal winds are often derived from meteor radar observations assuming a zero vertical wind, which
65 apparently results in reliable wind speeds compared to meteorological analysis data such as the Navy Global Environment
Model - High Altitude (NAVGEM-HA) (Eckermann et al., 2018; McCormack et al., 2017). However, there were also some
attempts to fit vertical winds to the observations (e.g., Egito et al., 2016; Chau et al., 2017; Conte et al., 2021; Chau et al.,
2021, and references therein), which resulted in spurious and apparently very fast vertical motions of up to 20 m/s over several
70 hours or up to 10 m/s over several days. Considering the large observational volumes of about 350 km in diameter in the
mesosphere, these values are unlikely to be representative of typical atmospheric motions. For such high vertical velocities to be
sustained over hours or even days would require large energy reservoirs, and would be accompanied by strong adiabatic cooling
(heating) for upwelling (downwelling) motions, which so far has not been confirmed by co-located satellite observations or
other temperature measurements.

In this study, we investigate potential biases of meteor radar wind measurements and present mathematical approaches to
75 minimize their impact on the estimated parameters with a particular emphasis on vertical winds. We present observations
from monostatic meteor radars as well as from multistatic meteor radar networks such as the Nordic Meteor Radar Cluster
and CONDOR (Chilean Observation Network De meteOr Radars) (Stober et al., 2021a). The vertical wind bias is discussed
considering the trail physics and scattering geometry (Poulter and Baggaley, 1977; Jones and Jones, 1990; Stober, G. et al.,
2021). Furthermore, fragmentation of meteoroids plays a role in the trail formation and, thus, could lead to biases due to the



Table 1. Technical parameters of the Nordic Meteor Radar Cluster, CONDOR (ALO), Tierra Del Fuego (TDF) and Collm (COL).

	TRO	ALT	SOD	KIR	TDF	ALO	COL
Freq. (MHz)	30.25	31	36.9	32.50	32.55	35.1	36.2
Peak Power (kW)	7.5	8	7.5/15	6	64	48	15
PRF (Hz)	500	430	2144	2144	2144/625	430	2144/625
coherent integration	1	1	4	4	4/1	1	4/1
pulse code	4-bit complementary	4-bit complementary	mono	mono	7-bit Barker	4-bit complementary	7-bit Barker
sampling (km)	1.8	1.8	2	2	1.5	1.8	1.5
latitude	69.59°N	70.0°N	67.4°N	67.9°N	53.7°S	30.3°S	51.3°N
longitude	19.2°E	23.3°E	26.6 °E	21.1°E	67.7°W	70.7°W	13.0°E

80 more complicated trail physics (Subasinghe et al., 2016; Vida et al., 2021). However, as it is not feasible to analyze all these physical processes for each individual meteor, it is nearly impossible to correct these effects for each meteor. Thus, we propose mathematical approaches to reduce potential biases by introducing mathematical parameterizations of these effects.

2 Meteor Radar observations and sampling biases

Meteor radars have been widely used to investigate atmospheric dynamics as well as meteor astronomy over the past decades
85 (Hocking et al., 1997, 2001; Portnyagin et al., 2004; Brown et al., 2008a; Fritts et al., 2010a; McCormack et al., 2017; Stober et al., 2012, 2021a; Janches et al., 2015). The systems have been proven to be reliable and suitable for long-term continuous and automated observations of MLT winds and tides (Larsen et al., 2003; Franke et al., 2005; Jacobi et al., 2007; Wilhelm et al., 2019; Stober et al., 2021b; de Araújo et al., 2020). In this study, we use data from two multistatic meteor radar networks, which are the Nordic Meteor Radar Cluster (NORDIC) and CONDOR, as well as the single station meteor radars at Collm 90 (COL) and Tierra del Fuego (TDF). The Nordic Meteor Radar Cluster consists of 5 monostatic systems at Svalbard (SVA), Tromsø (TRO), Alta (ALT), Kiruna (KIR) and Sodankylä (SOD). CONDOR makes use of the monostatic radar at the Andes Lidar Observatory (ALO), and two passive receiver systems at the Southern Cross Observatory (SCO) and at Las Campanas Observatory (LCO). Table 1 contains an overview of the geographic location of all systems and the corresponding experiment settings.

95 MLT winds are obtained from meteor radar observations by applying a so-called all-sky fit (Hocking et al., 2001; Holdsworth et al., 2004), which minimizes the projection of all measured radial, or line-of-sight, velocities onto a mean 3D wind within an altitude-time bin in a least-squares sense. The radial wind is often written as;

$$v_r = u \cos(\phi) \sin(\theta) + v \sin(\phi) \sin(\theta) + w \cos(\theta) \quad . \quad (1)$$



Here v_r is the line of sight velocity, u , v and w represent the 3D wind velocities in the zonal, meridional and vertical direction, 100 θ denotes the off-zenith angle and ϕ is the azimuth angle counterclockwise from East. In general, the vertical wind is assumed to be negligible ($w=0$ m/s), which simplifies the equation to the horizontal wind components. Obviously, this assumption is justified considering the good agreement of the obtained horizontal winds when compared to meteorological analysis data (McCormack et al., 2017; Stober et al., 2019; Liu et al., 2020) and the large observation volume of about 350 km in diameter as well as the typical temporal resolution of one hour.

105 Although it appears to be legitimate to make the simplification and to remove the vertical wind from the radial wind equation, there is a need for a mathematical justification. Therefore, we investigate the bias that is intrinsic to meteor-radar wind estimates by implementing different data analysis pipelines to the COL and TDF meteor radar using three months of data from January to March 2020. The first data analysis applies a least squares fit using all three wind components, a non-linear error propagation and WGS84 geometry. The wind components are estimated by a singular value decomposition as solver (Press 110 et al., 1992). The second data analysis leverages the same observations, but all radial velocities were replaced by synthetic data sustaining the spatial and temporal sampling of the original measurements and their corresponding statistical errors. The synthetic wind field is composed of an altitude dependent mean wind, planetary waves, and tides plus some gravity waves. However, the vertical wind component was set to zero for all waves and the mean wind at all altitudes and times.

Figure 1 shows four histograms of hourly fitted vertical winds applying the classical least squares approach solving the radial 115 wind equation. The left histograms present the vertical winds from our 'naive' data analysis. The right panels visualize the radial velocity distribution for the synthetic data where we put a zero vertical wind component for all waves. The histograms indicate rather large 'apparent' vertical velocities. In particular, the analyzed synthetic data demonstrates that there are substantial biases. However, the synthetic data also exhibits a reduced standard deviation compared to the 'naive' least squares 120 solutions suggesting that there is at least some sensitivity left to 'infer' a residual vertical wind from the observations. The difference between TDF and COL for the synthetic data is only related to the observational statistics. TDF has about twice the number of detections during this part of the season.

There are many reasons for the intrinsic bias in the meteor radar vertical winds. Some of them are almost impossible to address due to the lack of information provided by the current generation of meteor radars. For instance, the question arises how fragmentation affects the radial velocity measurement and the interferometric solution. Trajectory information to correct for 125 geometric offsets due to the specular/transverse scattering geometry is often not available. Recent studies of high resolution optical observations indicated that almost 90% of the observed meteors exhibit signs of fragmentation (Subasinghe et al., 2016; Vida et al., 2021). There is also the question whether strong wind shear or turbulence induce an apparent motion of the scattering center along the trail axis. Most meteor radars lack the capabilities to investigate and quantify these effects in detail. Only very few systems provide multistatic trajectory measurements, which are required to remove most of the wind shear and 130 geometric effects (Webster et al., 2004; Jones et al., 2005; Brown et al., 2008b; Fritts et al., 2010b; Panka et al., 2021).

However, our synthetic data analysis points out that there are also mathematical and geometrical reasons causing an intrinsic bias in the vertical velocities due to the spatial and temporal sampling. The synthetic data does not suffer from any disturbances related to the meteor trail physics. All radial velocities and their interferometric locations in the WGS84 coordinates are exactly

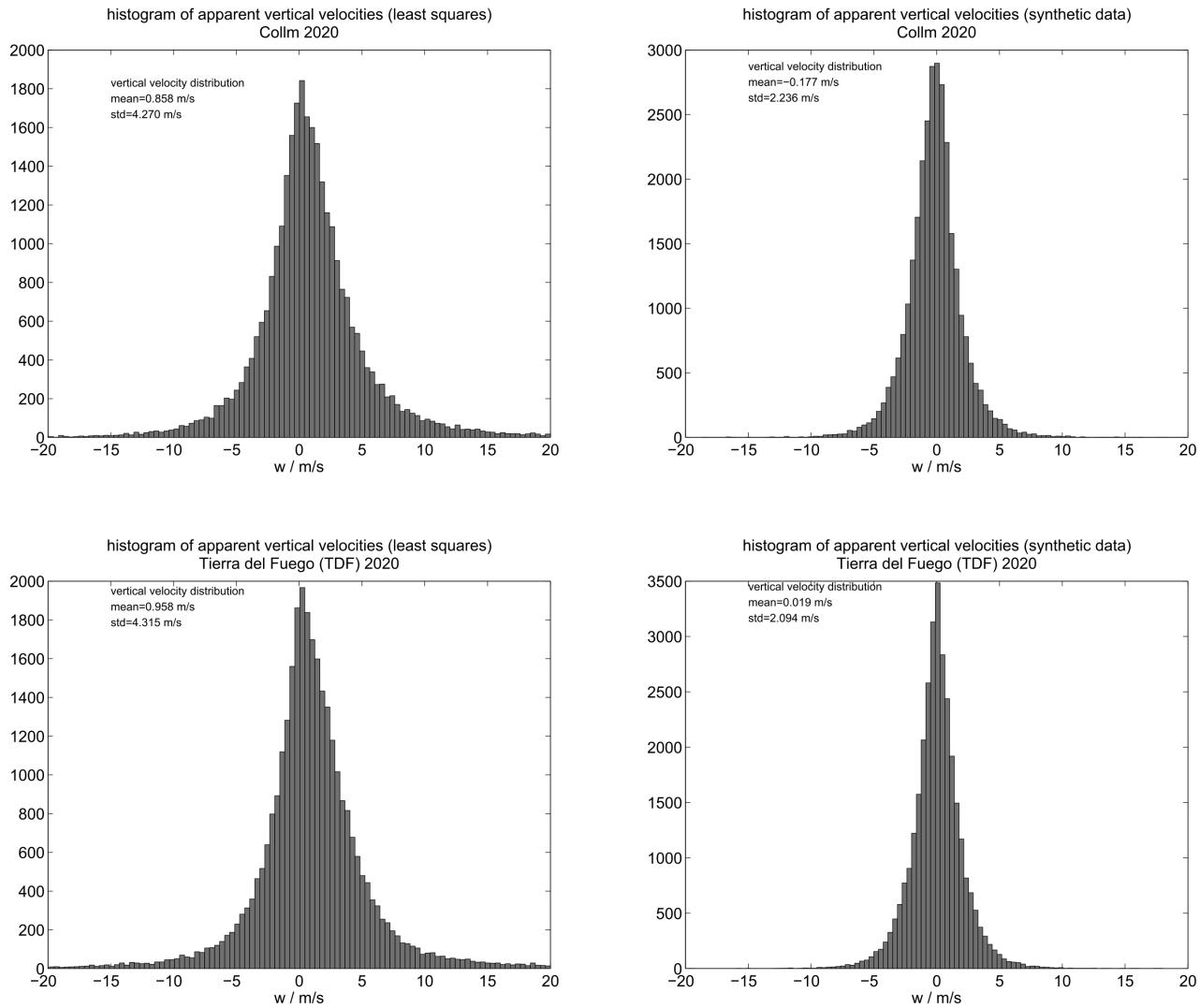


Figure 1. Histograms of the residual bias vertical velocities derived from the COL and the TDF meteor radar using observations from January to March 2020. The left histograms shows the results of the hourly residual bias vertical velocities applying a least squares fit. The right panels show the resulting vertical velocities applying the same algorithm using the COL and TDF detections (volume sampling), but with synthetic data based on mean winds, planetary waves, and tides and a zero vertical velocity.



135 determined and, thus, only numerical and sampling aspects due to the time-altitude binning contribute to the standard variation of the distribution shown in Figure 1 (right panel). Furthermore, the radial wind equation is linear in all three wind components, which results in a weighted measurement response of the sin and cos-terms for the off-zenith angles. Typical meteor radars detect most of the meteors at off-zenith angles between 55-65° corresponding to a scale factor of 1.2 to 1.3 in the geometric measurement response between the horizontal components and the vertical wind. In addition, it is worth to consider that the magnitude of the horizontal wind velocity is often more than a factor of 10 larger compared to the vertical wind magnitude.

140 The consequence of these scaling-terms is also reflected in the statistical uncertainties of the fitted wind coefficients, which range between 2-12 m/s or occasionally more than 15 m/s for each coefficient. These statistical uncertainties are reasonable for horizontal winds, which very often exceed 20-40 m/s as mean wind speed, but are too large to retrieve physical and statistically sound solutions for the vertical velocities.

3 Bias related to scattering geometry

145 Transverse scatter or specular meteor radars are highly sensitive to the observation geometry. Full wave scattering simulations point out that there is a strong polarization dependence between the trail alignment and the polarization of the incident radio wave (Poulter and Baggaley, 1977; Stober, G. et al., 2021). The concept of meteor radar wind observation is based on the assumption that most of the backscattered energy originates from the specular point, which is assumed to be a well-defined location along an infinitely long ambipolar diffusing plasma column. However, the scattering point describes the motion of

150 the scattering center rather than a well-defined location of the meteor trail. Thus, depending on the observing geometry, the measured Bragg vector denotes the motion of this scattering center, which is composed of the trail motion and apparent changes caused by the scattering geometry. These changes in the geometry are related to horizontal or vertical winds and wind shears, and electron line density variations caused by turbulence, fragmentation (Subasinghe et al., 2016; Vida et al., 2021), or differential ablation (Vondrak et al., 2008).

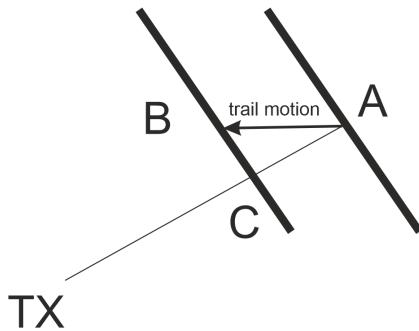
155 Figure 2 schematically illustrates how these apparent motions of the specular point relate to purely horizontal or vertical movements of the trail. The letter 'A' describes the position of the specular point along the trail after the meteoroid passed the t_0 -point (closest distance to the transmitter, see also McKinley (1961); Mazur et al. (2020)) and 'B' labels the location of this specular point if it stays 'glued' to the trail, while the meteoric plasma column is drifted by the neutral wind. 'C' labels the position of the scattering center considering the trail motion, but sustaining the geometry regarding the transmitter and receiver

160 location (TX/RX). Although the concept of the specular point as a reflection center is already a substantial simplification of the scattering process, the scheme visualizes the basic geometric problem. A more realistic approach considers that the scattering actually occurs from an extended section of the meteoric plasma trail along the meteor flight path containing several Fresnel zones around the specular point.

165 The later point is of particular concern for multistatic or forward scatter meteor radar observations. Due to the more slant incident radiowave the scattering section along the trail is much longer. Already Stober and Chau (2015) demonstrated that the forward scatter angle corresponds to a frequency shift to lower frequencies and, thus, to even larger Fresnel zones. Hence,

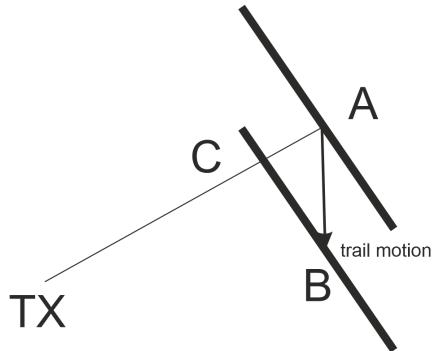


trail moves horizontally



So the trail moves from A to B, creating an apparent vertical motion of the scattering point from A to C.

trail moves vertically



So the trail moves from A to B, creating an apparent horizontal motion of the scattering point from A to C.

Figure 2. Idealized schemes of the specular scattering geometry indicating the apparent motion of the specular point or scattering center along the trail due to the drift of the meteor plasma column by neutral winds. The label 'A' shows the position of the specular point at the first detection, 'B' denotes the location of the scattering center assuming it stays glued to the trail and 'C' shows the position of the reflection point sustaining the transmitter and receiver geometry.

changes of the electron line density within the scattering section along the trail act as an additional weighting and lead to even more pronounced apparent motions of the specular point, which can slide along the meteor trajectory. This sliding can be caused by changes of scattering geometry due to winds and wind shears or by modifications of the electron line density that

170 are associated with fragmentation and differential ablation.

We evaluate the above described hypothesis by performing a normal wind analysis using all three data sets provided by the CONDOR network in Chile. The network is unique in the sense that it combines monostatic and forward scatter passive receivers in a fairly compact geographic region. Figure 3 shows a comparison of the zonal (left column) and meridional (right column) winds using only the northernmost site at LCO, the standard meteor radar at ALO and the passive receiver at SCO.

175 A geographic map of all three sites can be found in Stober et al. (2021a). The observation volumes are basically overlapping, and, thus, it is reasonable to expect that a climatological comparison should result in an almost identical mean wind behavior. However, zonal winds exhibit large differences especially during May to September and at altitudes below 85 km. Above 85 km, discrepancies appear to be much smaller. The excess of the zonal wind magnitude between the monostatic (ALO) and the forward scatter stations is about a factor of 2 around 80 km and below. There is no geophysical reason why in such a narrow 180 latitudinal band the zonal wind should show such significant changes. We reproduced these results with a commercial software to rule out any issues caused by the retrieval algorithm that is described in detail in section 4. It is evident from Figure 3 that the zonal wind appears to be significantly stronger above the passive forward scatter receivers. Meridional winds seem to be much

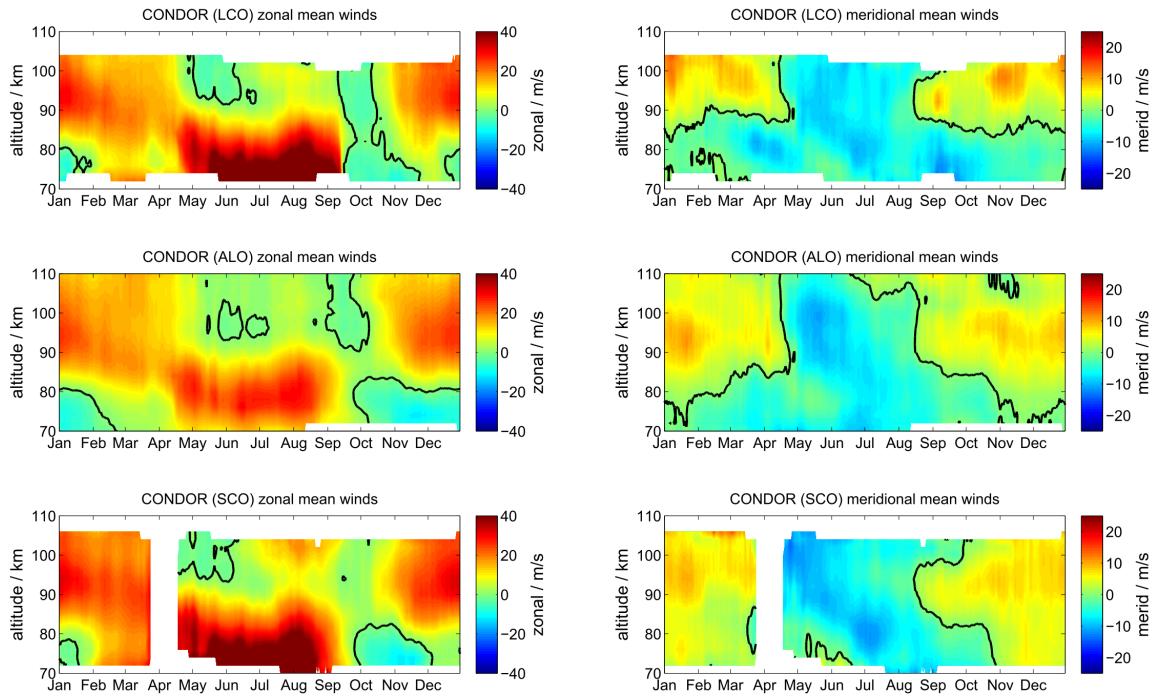


Figure 3. Comparison of zonal and meridional winds for the forward scatter receiver stations LCO and SCO and the monostatic radar in ALO. The left column shows the zonal wind component and the panels on the right the meridional wind. The panels are sorted according to their geographic latitude with the northernmost sampling volume at LCO on top, the ALO in the center and the southernmost station SCO in the lowest row.

less affected, although there are substantial difference between the northernmost and southernmost location, which are only separated by 3° in latitude. Our preliminary analysis thus already reveals that there is a considerable and altitude dependent difference in the wind magnitude between monostatic and passive receiver systems.

Finally, we investigate whether the magnitude difference manifests also in the Bragg vector pointing direction between the forward scatter receivers at SCO and LCO relative to the monostatic radar at ALO. In Figure 2 we hypothesized that the Bragg vector pointing direction is not affected by the trail motion due to the wind, which is described by a parallel translation, and, thus the Bragg vector pointing is supposed to remain perpendicular to the meteor trajectory for underdense meteors, whereas the length of the Bragg vector is a measure of the total path of the scattering center over successive radar pulses, which includes the motion of the trail due to the neutral wind plus an apparent sliding of the scattering center (specular point) along the trail. We computed the source radiant of two well-known and long lasting (several degrees in solar longitude) meteor showers applying a modified single station radiant mapping algorithm (Jones and Jones, 2006). The meteor source radiant maps for SCO, LCO and ALO were obtained by implementing a revised version of the algorithm applied in Stober et al. (2013). The new generalized radiant mapping is based on the WGS84 geometry for each individual meteor. There have been already several meteor shower catalogues published in the literature covering the northern and southern hemispheres (Brown et al., 2010;



Janches et al., 2013; Pokorny et al., 2017) and, hence, it was easy to pick some of the established meteor showers for the solar longitudes of concern. Figure 4 shows six radiant activity maps for LCO, SCO and ALO. At the beginning of May all three sites exhibit an increased activity at the source radiant of the *eta*-Aquariids (ETA), which are visible at right ascension $\alpha = 337^\circ$ and declination $\delta = -0.9^\circ$. This meteor shower is active between solar longitudes of $\lambda_{sol}=30-60^\circ$, corresponding to end of April until May (Brown et al., 2010; Janches et al., 2013). The second shower that we found was the Daytime zeta Perseids (ZPE), which is visible at right ascension $\alpha = 56.6^\circ$ and declination $\delta = 23.2^\circ$. Daytime zeta Perseids are active between solar longitudes $\lambda_{sol}=56-90^\circ$, corresponding to end of May until June (Brown et al., 2010; Schult et al., 2018). The right ascension and declination coordinates are provided for the days around the maximum meteor shower activity. These radiant activity maps indicate no systematic differences between the forward scatter stations at SCO and LCO and the monostatic radar at ALO. Thus, the Bragg vectors are correctly determined for all stations and reflect no substantial deviation of the source radiant for these two meteor showers. In particular, the Daytime zeta Perseids have a geocentric velocity of $v_g = 28 - 32$ km/s and, hence, can penetrate deep into the atmosphere and reach the altitudes where we already see significant differences in the wind magnitudes. In summary, we were not able to identify a similar deviation in the source radiant mapping of two major meteor showers between the forward scatter receiver stations and the monostatic radar that corresponds to or explains the magnitude offset that is evident in the zonal winds.

4 Mathematical debiasing strategies

After we introduced the intrinsic bias of the vertical wind estimates in meteor radar observations, we are going to briefly discuss mathematical debiasing strategies. The most straightforward method is to implement a Tikhonov regularization in the least squares fitting (Wilhelm et al., 2017; Stober et al., 2017). However, this approach leads to a brute force norm reduction and depends on an empirically determined Tikhonov matrix and Lagrange multiplier;

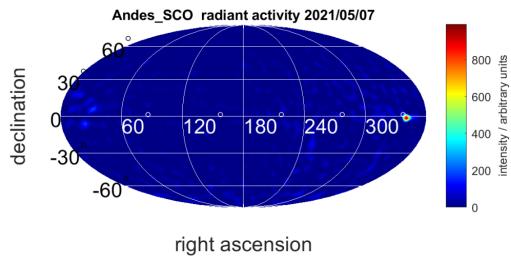
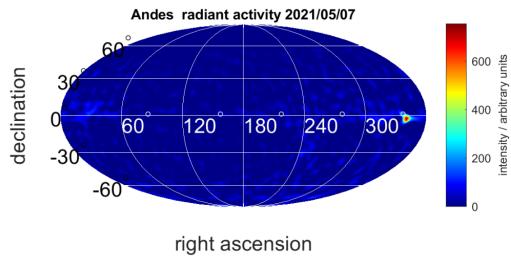
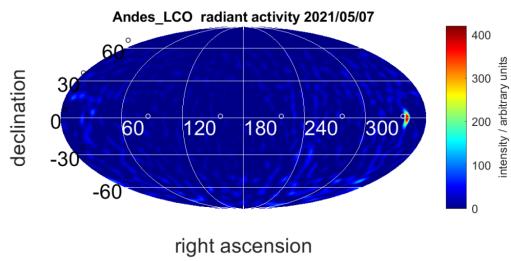
$$\|Ax - b\|^2 + \lambda\|\Gamma x\|^2 \quad . \quad (2)$$

Here A is the Jacobian matrix of the problem, x is our state vector, b are the observations, Γ denotes the Tikhonov matrix, λ describes the Lagrange multiplier (here and further on $\lambda = 1$), and the superscripts denote the Euclidean norm. It is now possible to construct a Tikhonov matrix in such a way that $\lim_{\Gamma_w \rightarrow \infty} \Gamma x \rightarrow \infty$, which results in $w=0$ m/s for all solutions and, thus, is equivalent to the assumption of a negligible wind. The infinite growth of the right hand side enforces a norm reduction for the vertical wind and, hence, the vertical wind solution converges to zero. However, it is also possible to insert a solution of $\Gamma_w \in [0, \infty)$, which in consequence leads to a strong damping of the vertical velocities. The most straightforward approach is to use the unit matrix as Tikhonov regularization.

Although a Tikhonov regularization is suitable to enforce small vertical velocities, we are going to outline an even more complex approach to solve for the vertical wind. To this end, we modify the Tikhonov regularization to a filter function, which is also known as generalized Tikhonov regularization. Due to the implemented spatio-temporal Laplace filter in the meteor radar retrievals, it is straightforward to estimate a predictor for the state vector x_a for each time-altitude bin (Stober et al.,



Eta Aquariids



Daytime zeta Perseids

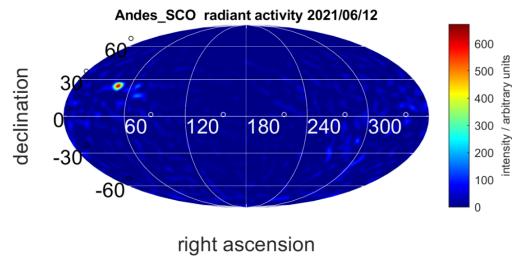
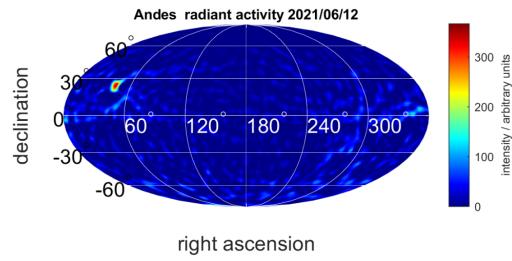
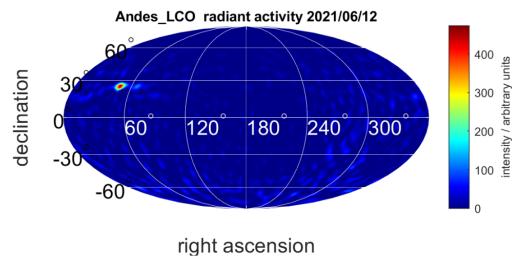


Figure 4. Meteor radiant activity maps derived from CONDOR for LCO, ALO(Andes) and SCO (the top, middle and bottom row, respectively). The left column shows the source radiant activity for the η -Aquariids and the right panels present the Daytime zeta Perseids. The meteor showers are identified from the catalogues presented in Brown et al. (2010); Janches et al. (2013).



2019, 2021a). Furthermore, we can insert constraints to the error covariance for the state vector accounting for the above
230 described scaling effects between the horizontal and vertical wind components. Thus, we now solve the problem using the
form;

$$\|Ax - b\|_P^2 + \lambda \|\hat{\Gamma}(x - x_a)\|_Q^2 . \quad (3)$$

Here, P denotes the inverse covariance matrix of b , and Q is the inverse covariance of x including a scaling term for the
vertical wind component to remove the bias. The advantage of the new norm reduction is that for small differences $x - x_a$ and
235 reasonable covariance errors the solution is identical to the least squares fit as the right hand term of Eq. 3 basically vanishes.
By construction the right hand term permits a certain part of the solution to pass through the spatio-temporal Laplace filter
depending on its covariance. The larger the statistical uncertainties the stronger and more important becomes the right hand
term, which often results in smaller vertical velocities.

Furthermore, the spatio-temporal Laplace filter is also beneficial for the horizontal wind components to compensate and reduce
240 effects caused by the random and irregular spatial and temporal occurrence of meteors within the sampling (observation)
volume of the radar. Sometimes, small or even tiny measurement errors in the location of a meteor may induce large projection
errors in the final solution of the retrieved wind components, which is minimized when applying the spatio-temporal filter.

Figure 5 shows the vertical velocity histograms based on the retrieval algorithm applying the spatio-temporal Laplace filter
and the empirical bias correction based on the scale analysis. The left panel shows the inferred vertical velocities based on
245 the original COL observations. The histogram in the right panel is obtained when the synthetic data set, with all vertical
wind values being zero, is analyzed with the retrieval algorithm. The remaining width of the distribution is caused by the
sampling window in time and space (vertical bin size) and the other atmospheric waves. However, already this simple de-
biasing approach, where we just consider the scale analysis described above substantially reduced the offset that was inherent
250 when only a 'standard' least squares wind fit (Press et al., 1992) was applied (Figure 1 right panel). Although, however,
generalized Tikhonov regularizations or filtering functions such as the spatio-temporal Laplace filter can help to reduce the
intrinsic bias in the meteor radar wind analysis to determine vertical winds by comparing idealized synthetic data, we are still
not able to prove the reliability of the derived vertical winds beyond their statistical properties due to a missing ground 'truth'.

5 Statistical comparison to the non-hydrostatic UA-ICON model

A direct comparison of the retrieved vertical winds to other observations is not feasible due to the lack of such measurements.
255 Therefore, we prepared a statistical comparison to a recently developed state of the art non-hydrostatic General Circulation
Model (GCM). The Upper Atmosphere ICOsahedral Non-hydrostatic (UA-ICON, Borchert et al., 2019) extends the vertical
coverage of the ICON numerical weather prediction model from 80 km to about 150 km altitude. A detailed description of the
upper atmosphere physics is given in Borchert et al. (2019). The upper atmosphere version leverages the numerical weather
prediction physics packages (Zängl et al., 2015; Giorgetta et al., 2018; Crueger et al., 2018). Here we made use of a 21
260 year free running climate simulation without any nudging and parameterized gravity waves on a so-called R2B4 grid with

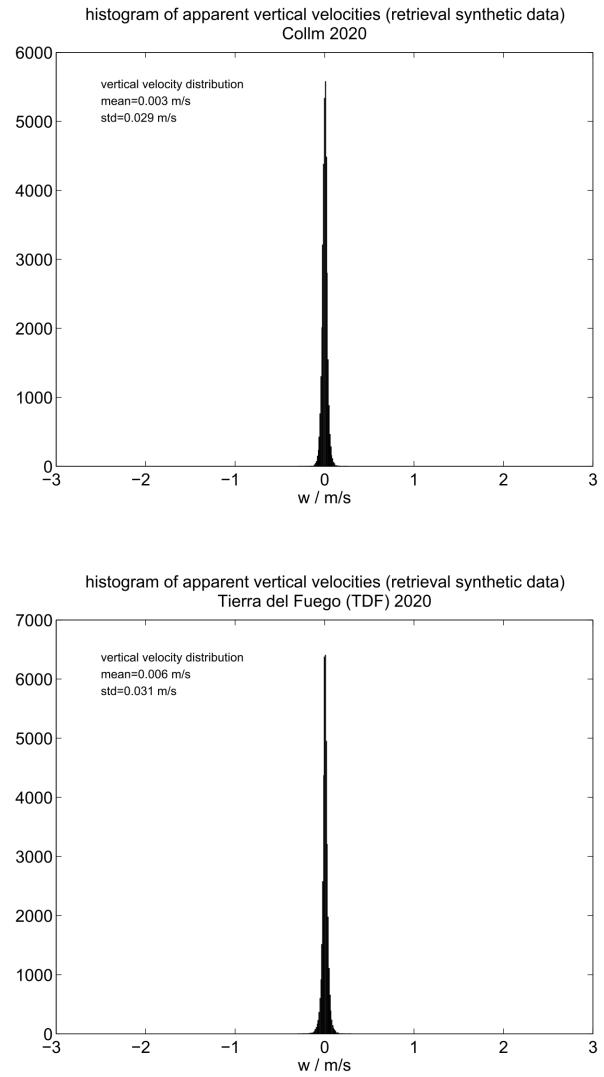
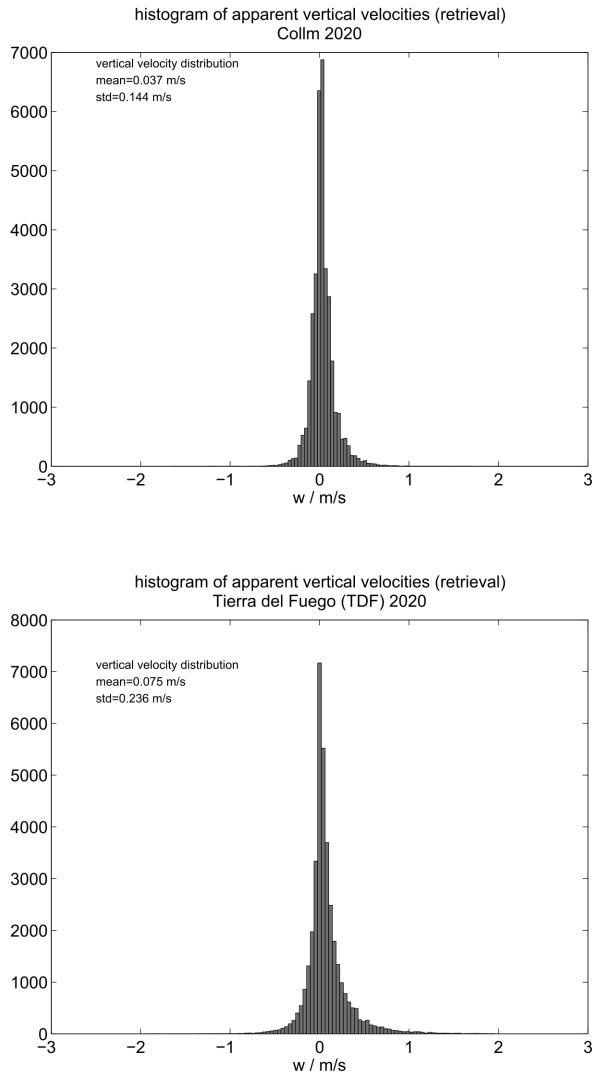


Figure 5. The same as Figure 2, but the hourly vertical winds are obtained by applying the retrieval algorithm including the spatio-temporal Laplace algorithm. The x-axis scale or w -axis scale was reduced to show the remaining variability.



a horizontal resolution of 160 km (Borchert et al., 2019; Giorgetta et al., 2018). Above 120 km altitude the model applies Rayleigh damping on the vertical winds. The benefit of the UA-ICON model for such a comparison is that vertical winds are available on a geometric vertical coordinate grid. The UA-ICON horizontal winds and tides have been compared to WACCM-X(SD), GAIA and data from six meteor radars (Stober et al., 2021b). Similar to this study we extracted vertical winds by 265 considering the instrument observation volume.

Figure 6 shows a statistical comparison of hourly retrieved vertical wind velocities for the COL and TDF meteor radars. These histograms are obtained using the entire available data set for both systems, which covers 16 years in the case of COL and about 12 years for TDF. The left column presents the observations from both meteor radars and the right column shows the corresponding UA-ICON data. The histograms exhibit a remarkable agreement of the inferred debiased vertical 270 velocities. The observations, however, indicate more variability compared to the GCM. However, the overall agreement of the vertical velocity distribution between the observations and UA-ICON data reveals that at least the statistical moments of the distributions have significantly improved compared to the least squares derived vertical winds. Furthermore, it is possible to use the skewness of the histogram to estimate potential systematic issues of the radar either due to irregular detections within the radar beam volume or issues in the interferometric solution (e.g., technical problems). Although the debiasing seems to 275 provide reasonable results, we cannot assess the reliability of individual observations or identify other systematic effects due to the more complicated scattering process (e.g., fragmentation, differential ablation and so forth). Thus, we intend to go beyond these simple approaches and further refine the retrievals to implement physically and mathematically consistent solvers to infer more reliable vertical wind velocities and vertical wind variability.

6 3DVAR+DIV retrieval

280 Recently, a 3DVAR algorithm was introduced to retrieve spatially resolved 3D winds using multistatic meteor radar observations from the Nordic Meteor Radar Cluster and CONDOR (Stober et al., 2021a). This 3DVAR algorithm already included the retrieval of vertical winds, but required a Tikhonov regularization to reduce the numerical instabilities, which often arise for parameters with low/poor measurement response. Due to the much worse statistics per grid cell, the quality of each radial velocity measurement comes even more into play and we have to consider the representativeness of a single measurement.

285 This is achieved by introducing a smoothness constraint or variable correlation lengths inside the domain. Such correlation lengths are described by the averaging kernel. However, the zonal and meridional wind components exhibited a reasonable measurement response inside the retrieval domain with values beyond 0.6 and more, indicating short correlations or narrow averaging kernels. Another benefit of the 3DVAR approach was the possibility to add additional constraints by expanding the cost function, e.g., for data assimilation of other observations.

290 The new 3DVAR+DIV algorithm was revised and expanded by adding a divergence constraint to the cost function. For this, we implemented diagnostics to estimate the horizontal divergence and relative vorticity for each grid cell. We consider that an air parcel that is moved by neutral winds should satisfy the continuity equation;

$$\frac{\partial \rho}{\partial t} + \rho \cdot \operatorname{div}(\mathbf{u}) = 0 \quad . \quad (4)$$

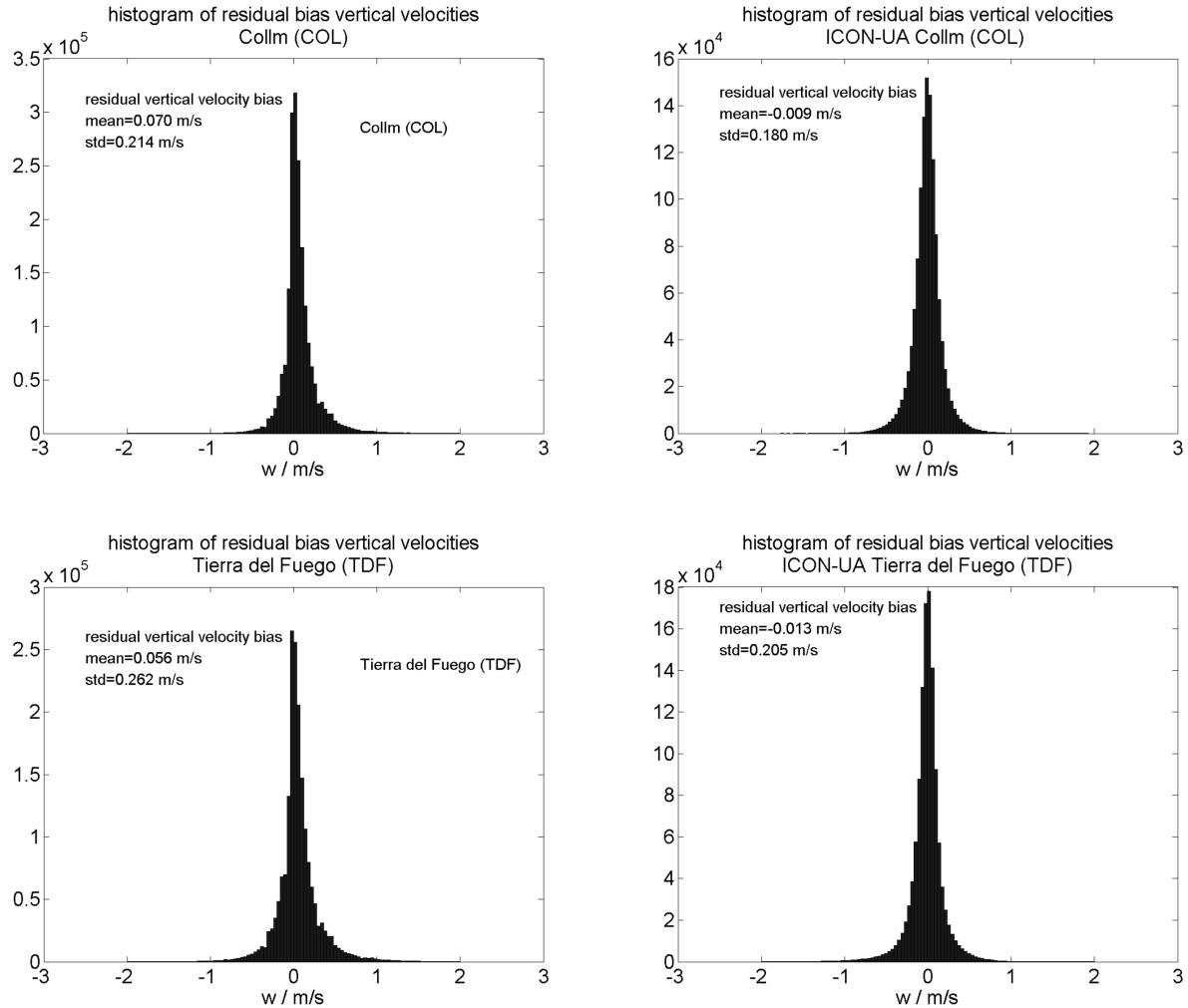


Figure 6. Histograms of the residual vertical velocity for the available data at COL and TDF including the debiasing from the spatio-temporal Laplace filter. The left panels show the meteor radar observations. The right panels visualize the corresponding UA-ICON velocities for a typical meteor radar sampling volume.



295 Here ρ is the mass density of the air. Furthermore, it is reasonable to assume a stationary process for each time step, which is equivalent to $\partial\rho/\partial t = 0$. Thus, the continuity equation simplifies and we only have to derive the divergence for each voxel. The divergence is given in Cartesian coordinates by;

$$\text{div}(\mathbf{u}) = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \quad . \quad (5)$$

296 In the 3DVAR algorithm, variable domain geometries could be used (Stober et al., 2021a). Therefore, the numerical solution of the derivatives to diagnose the horizontal divergence uses a first order approximation of the elliptical integrals for the WGS84 300 reference coordinate systems (National Imagery and Mapping Agency, 2000), which appears to be sufficient for most of the typical voxel sizes of a few tens to hundreds of kilometers or a few degrees in latitude and longitude.

Assuming an incompressible flow, we can estimate the change of the vertical velocity Δw between two vertical layers and for each grid cell by;

$$\Delta w_i = \int_{z_1}^{z_2} \text{div} \cdot v_{h_i} \, dz \quad . \quad (6)$$

305 Here the index i denotes the grid cell within a layer, and z_1 and z_2 are the upper and lower boundaries, respectively, describing the layer thickness.

306 The 3DVAR+DIV algorithm solves all equations through several iterations. The first call is again the standard 2DVAR retrieval, which permits to obtain a first estimate of the horizontal divergence, which can be integrated for each grid cell assuming a lower boundary of the vertical velocity $w(z_0)_{0,i}$. From the second iteration, we include the continuity equation and perform the full 310 3DVAR+DIV retrieval.

315 To solve for the vertical velocity at each altitude and grid cell, we need to integrate equation (6) from below or above, which requires an initial value $w(z_0)_{0,i}$. Equation (6) only provides a relative measure for the change of the vertical velocity between two layers. The standard retrieval estimates this boundary in such a way that the mean vertical velocity (integrated over all altitudes) in each column for a defined domain grid is zero. This is equivalent to the assumption that the mean vertical motion in the column over large areas and a vertical dimension of approximately 20-40 km is close to zero.

320 However, the 3DVAR algorithm already included the full 3D wind solution for each grid and we just removed the Tikhonov regularization, which damped the numerical instabilities, in the new 3DVAR+DIV retrieval. These vertical velocities are called compressible/non-stationary solution, because we permit at least some deviation from zero in equation (5) without defining an explicit threshold. The major advantage of the 3DVAR+DIV retrieval is now given by providing a compressible/non-stationary and incompressible solution for the vertical velocity for each grid cell. The incompressible solution only makes use of the vertical velocity gradient obtained from the horizontal divergence equation to minimize the numerical instabilities caused by the low geometric measurement response in large parts of the domain. Thus, both solutions exhibit very similar morphology and only show some deviations in the absolute magnitude.



325 7 Results

The new 3DVAR+DIV retrieval is now implemented for routine data analysis for the Nordic Meteor Radar Cluster and CONDOR. The main goal was to infer more reliable vertical velocities using a more physical consistent data description in the forward model. The performance of the new algorithm is demonstrated using observations conducted during September 2021 after major upgrades of the TRO meteor radar. During this time of the year the circulation changes from typical summer conditions to the winter regime. There is a moderate gravity wave activity, and enhanced semidiurnal tides are present (e.g., Wilhelm et al., 2019; Stober et al., 2021b).

The results presented herein are based on the 3DVAR+DIV algorithm using the Cartesian geographic grid with 30 km horizontal spacing and WGS84 geometry with a temporal resolution of one hour and a vertical spacing of 2 km. Figure 7 shows four panels. The upper two panels present the zonal (left) and meridional (right) wind components for a single time bin and the altitude centered at 90 km. Black arrows represent the (horizontal) wind in grid cells that have enough meteor detections. The wind magnitude for each component is color coded. Reddish colors refer to eastward and northward winds, whereas bluish colors indicate westward or southward motions, respectively. The lower two panels visualize the corresponding measurement response (Shannon, 1948; Shannon and Weaver, 1949; Stober et al., 2021a). The whiter the color, the higher is the observation density, which allows to achieve high spatial resolution. The bluer the grid cells are, the more information is mixed from long distance correlations beyond the next neighboring grid cells corresponding to broader averaging kernels.

An essential improvement of the new 3DVAR+DIV algorithm is the embedded diagnostics of the horizontal divergence and relative vorticity between grid cells. These values are obtained by spatial derivatives qualitatively and quantitatively for all possible geometries and in both implemented domain grids (geographic and Cartesian (rectangular grid)). We use Euler steps at the domain edges and central differences for all other grid cells within the domain. Figure 8 shows the horizontal divergence (left panel) and relative vorticity (right panel) for the same altitude and time period as the winds shown previously. The horizontal divergence exhibits coherent structures that are likely associated with a superposition of several gravity waves. A more random pattern is reflected by the relative vorticity, which shows a more patchy and irregular structure. Both quantities reach a relative strength of about $\pm 2 \text{ m/s/km}$, and occasionally higher values were also found.

Finally, the retrieved vertical velocities (upper panels) and corresponding measurement response (lower panels) are shown in Figure 9. The absolute vertical velocities are obtained assuming a lower boundary, which was determined in such a way that the mean vertical velocity in the column above each grid cell is zero. The compressible/non-stationary and incompressible solution for the vertical velocities are almost identical, which is very often the case. As our forward model makes use of the continuity and radial wind equation, we have no independent estimate of the measurement response for the compressible/non-stationary solution and only the residuals of the radial wind equation contribute to the final estimate. Similar to the monostatic observations, the geometry of the meteor detections is not favorable to infer reliable vertical winds. Adding the continuity equation compensates for that but also dominates the measurement response and the overall contribution of the finally retrieved 3D winds. This is also reflected by the measurement response for the vertical velocities, which is identical for both solutions for the above mentioned reasons, and is dominated by the horizontal velocity measurement responses. We also investigated the

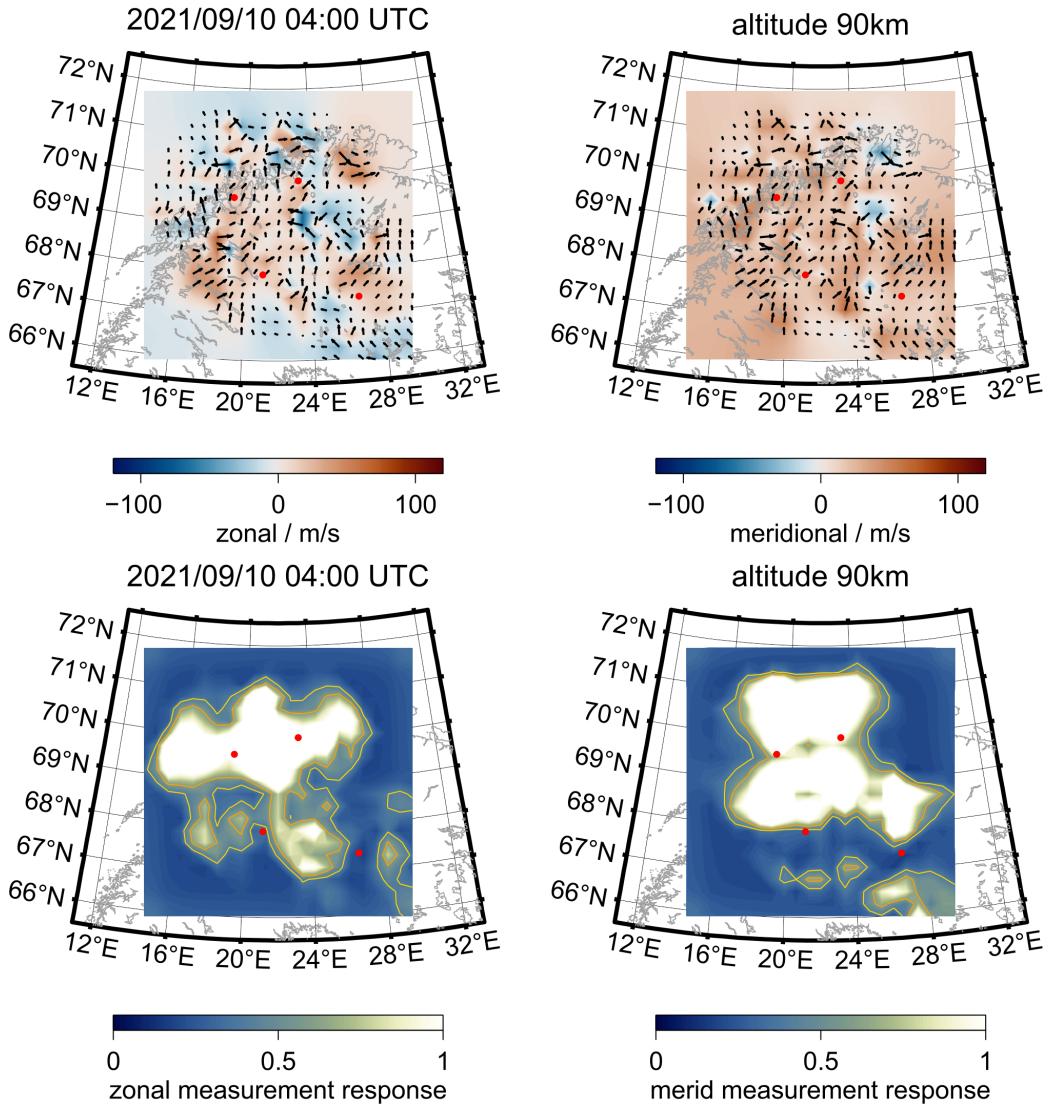


Figure 7. Snapshot of zonal and meridional winds and corresponding measurement response using the 3DVAR+DIV algorithm and measurements from the Nordic Meteor Radar Cluster. The red dots label the locations of the meteor radars.

statistical variability of the 3DVAR+DIV derived vertical velocities. Therefore, we analyzed the year 2021 from the Nordic
360 Meteor Radar Cluster and two weeks of data in March 2020 from CONDOR to estimate the statistical moments of the hourly inferred vertical wind measurements. The corresponding histograms are shown in Figure 10. The histograms only contain results for grid cells with a measurement response larger than 0.5 and the compressible/non-stationary solution. The incompressible solution (vertical(div)) exhibits an approximately 20% reduced standard deviation for the same periods. The offset of the mean from zero is caused by the lower integration boundary condition being determined including all grid-cells and

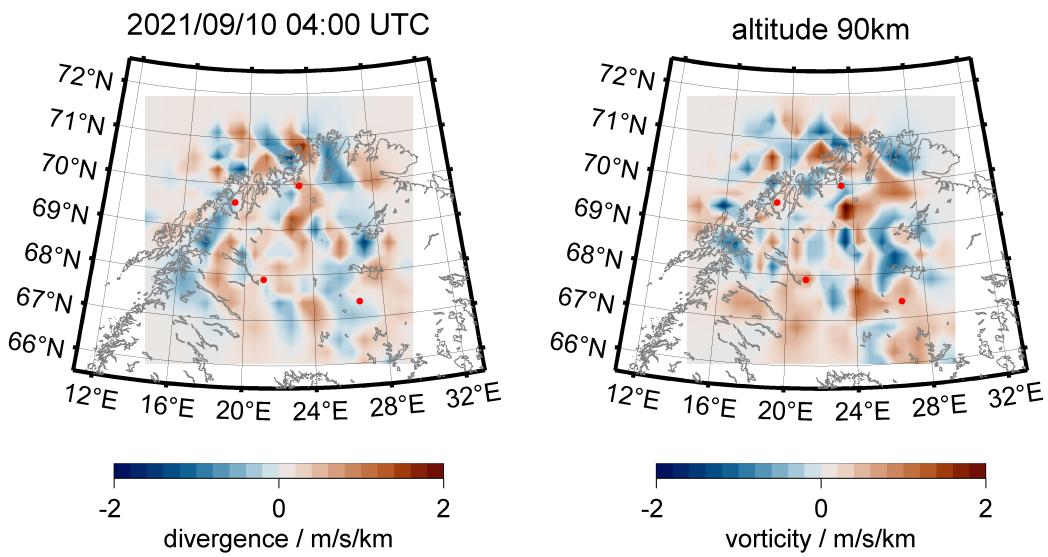


Figure 8. Horizontal divergence and relative vorticity calculated from the 3DVAR+DIV algorithm making use of the horizontal winds. The shown snapshot corresponds to the same period as in Figure 7.

365 altitudes, while the histograms only show a subset. Furthermore, CONDOR shows a much higher variability compared to the
366 Nordic Meteor Radar Cluster, which suggests that there is an increased gravity wave activity above the Andes. Considering the
367 different amount of data included in the histograms, we do not want to put too much focus on this difference in the vertical wind
368 variability. Both histograms provide a sufficient database to infer the order of magnitude of the vertical wind variability for a
369 30 km diameter area. Increased variability is expected since this is significantly smaller than the typical monostatic observation
370 volume.

8 Discussion

Vertical velocity measurements at the MLT are still very challenging. The most reliable observations have been carried out with High-Power-Large-Aperture (HPLA) radars such as EISCAT and MAARSY (Hoppe and Fritts, 1995a, b; Fritts et al., 1990;

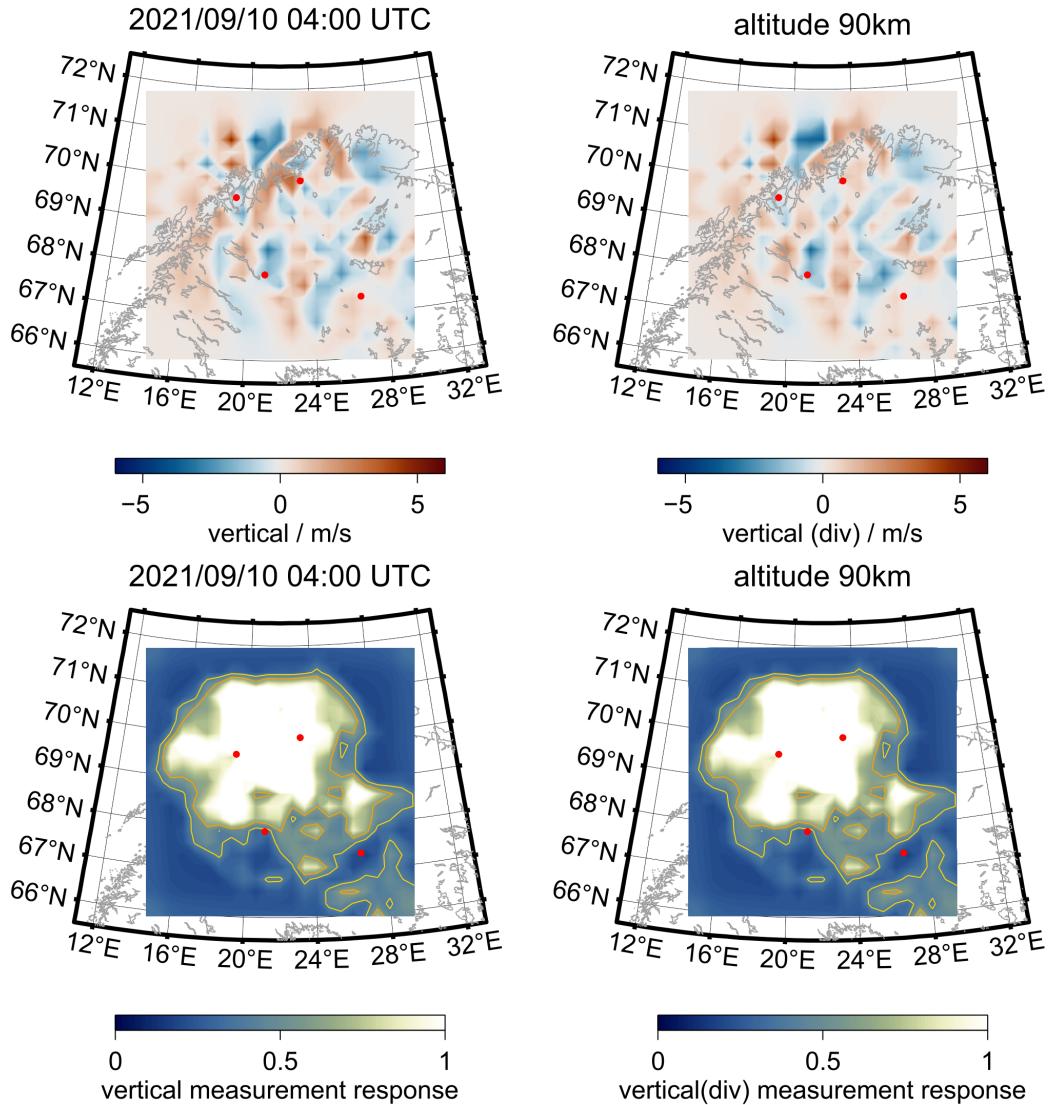


Figure 9. Corresponding vertical wind velocities (upper two panels) and measurement response (lower two panels) obtained by the 3DVAR+DIV algorithm for the same period as Figure 7. Note that the measurement response of the compressible/non-stationary vertical wind solution is dominated by the incompressible solution, which is used in all iterations due to the implemented horizontal divergence constraint.

Gudadze et al., 2019). However, these observations still did indicate biases when absolute magnitudes close to zero were tried 375 to be inferred. Some of these biases appear to be caused by gravity waves, as was reported for the EISCAT measurements. MAARSY results indicated a remaining uncertainty due to scattering from PMSE related to the sedimentation speed of the ice particles. However, HPLA radar measurements provides at least some valuable insights on the vertical wind variability and the

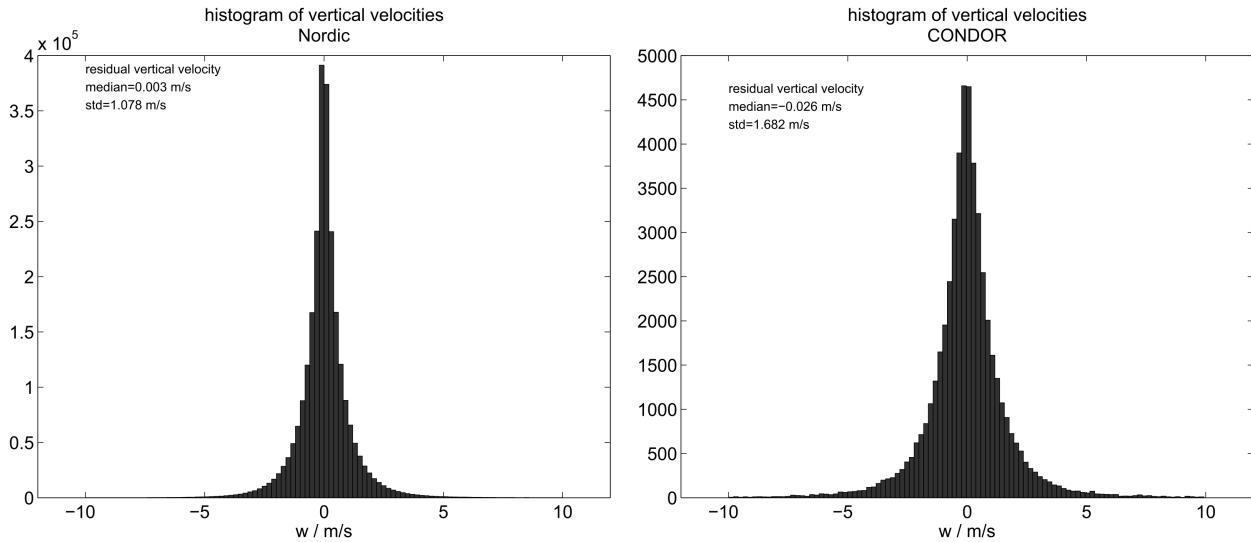


Figure 10. Histograms of hourly vertical winds obtained from the Nordic Meteor Radar Cluster and CONDOR.

magnitude of the vertical winds for the characteristic beam volumes and dwell times of the systems (seconds to minutes). These radars achieve statistical uncertainties down to a few cm/s, which is sufficient for most geophysical processes (Stober et al., 2018), but still leaves some ambiguities when it comes to the very small vertical velocities related to the residual circulation (e.g., Smith, 2012; Becker, 2012, and references therein).

There have been some attempts to derive mean vertical velocities from meteor radar observations applying least squares fits (Egitto et al., 2016; Conte et al., 2021). These meteor radar observations clearly exhibit intrinsic biases that can result in vertical velocities of more than a few m/s. In particular, Conte et al. (2021) reported vertical velocities in the excess of ± 10 m/s over 385 hours. Considering the large observation volume of a few hundreds of kilometers for a typical domain area for multistatic observations these values seem to be very large. Furthermore, based on measurements presented using data from the Nordic Meteor Radar Cluster and CONDOR, we were not able to reproduce these extreme values using the 3DVAR+DIV algorithm analyzing more than 2 years of data. Although there could be various reasons for such large values, we were able to identify some intrinsic biases related to the observation geometry and sampling and present mathematical debiasing strategies for 390 monostatic meteor radars using synthetic data. The proposed Tikhonov regularization and generalized Tikhonov or filter functions provide statistical sound solutions for the vertical winds. However, we also want to point out that since the assumption of a zero vertical wind seems to be justified in the context of these biases, this approach is mathematical equivalent to a Tikhonov regularization. Furthermore, large biased vertical velocities can degrade the quality of the horizontal wind solutions as well.

The comparison of the statistical distribution of the meteor radar inferred vertical velocities and the UA-ICON model gives 395 some confidence that the applied debiasing results in more consistent solutions for this wind component. However, there are still some sources of error left, which let us conclude that the term 'residual bias vertical velocity' seems to be the right term as we can not prove the correctness of individual hourly measurements. Fragmentation of the meteoroids and mean winds and



wind shears can lead to small changes in the scattering geometry, which cause an apparent shift of the scattering center along the trail. Thus, the Bragg vector of the scattered electromagnetic wave is not necessarily defined by the motion of the trail due
400 to neutral winds. Although these changes appear to be small, they affect the vertical component much more than the horizontal winds. In particular, these apparent motions of the scattering center along the trail could occur for transition echoes from over-dense to underdense, which could be caused by fragmentation (Subasinghe et al., 2016; Vida et al., 2021) or by differential ablation (Vondrak et al., 2008).

Almost a decade ago, there was a lidar study on vertical wind magnitudes related to atmospheric tides (Yuan et al., 2014). The
405 climatology exhibited vertical velocities of a few cm/s for large scale atmospheric tidal waves. The lidar observations indicated about 15-20 cm/s vertical wind magnitude for the semidiurnal tide and about 5-10 cm/s for the diurnal tide. These values appear to be consistent with the apparent vertical velocities estimated for the monostatic meteor radars at COL and TDF applying the spatio-temporal Laplace filter. The mid-latitude stations are dominated by semidiurnal tides during the hemispheric winter season and diurnal tides during the summer months (Stober et al., 2021b). However, the vertical wind magnitudes presented by
410 Yuan et al. (2014) and in this study are orders of magnitude lower than other estimates obtained from meteor radar observations (Egito et al., 2016) and multistatic meteor radar data (Chau et al., 2017, 2021; Conte et al., 2021).

Furthermore, we investigated systematic differences in the derived neutral wind velocities using data from CONDOR only. The comparison reveals a considerable difference in the estimated total wind magnitude during several months from May to August at altitudes below 85 km. The difference is most prominently visible in the zonal wind, but also the meridional wind is
415 affected, which is less obvious due to the much lower mean wind speeds. However, our radiant activity mapping of two meteor showers supports the above described scheme of a sliding scattering center or specular point along the meteor trajectory due to the motion of the trail by neutral winds. The source radiant maps only depend on the accurate determination of the pointing direction of the Bragg vector and, thus, are not affected by the apparent scaling of its magnitude due to the sliding of the scattering center along the meteor trajectory. Forward scatter receivers are more prone to this effect. Tiny changes in the geometry result in comparably larger apparent motions of the scattering center compared to monostatic systems. Furthermore, the effect increases the longer the trail lasts, corresponding to a slower diffusion, and, thus, mostly the lower altitudes are affected. However, the discovered bias in CONDOR between the monostatic and forward scatter mean winds is worth to be investigated
420 in more detail and opens the question on how to interpret the Bragg vector and corresponding radial motion concerning the specular or transverse scattering and the meteor trail geometry.

425 Multistatic observations are versatile and new approaches can be applied to improve vertical wind measurements. Considering the fast development over the past years from the first multistatic forward scatter meteor radar experiment (Stober and Chau, 2015) to more routine and established networks (Chau et al., 2017; Spargo et al., 2019) underlines the huge scientific potential of such observations. These first observations were analyzed making use of the classical assumptions on the vertical velocity ($w = 0$ m/s) or by fitting a mean value within the observation volume (Stober and Chau, 2015; Chau et al., 2017). However,
430 the retrieval of vertical winds remained challenging even when more advanced methods were applied (Volz et al., 2021). These advanced methods still resulted in vertical wind velocities of up to 10 m/s and more. The 3DVAR retrieval controlled the numerical instability in the vertical velocities by a Tikhonov regularization for each grid cell (Stober et al., 2021a). The new



3DVAR+DIV approach circumvents the need of an additional Tikhonov regularization by extending the forward model with the continuity equation, which permits to estimate horizontal divergence and relative vorticity directly to constrain the vertical 435 velocity solution.

The algorithm permits to obtain a compressible/non-stationary and incompressible solution for the vertical winds. Furthermore, the combined radial wind and continuity equations leverage the good measurement response from the horizontal wind 440 velocities, which significantly increases the measurement response for the vertical velocities as well. Due to the much smaller scales that are resolvable with the 3DVAR+DIV retrieval compared to standard monostatic meteor radars, it is also expectable to observe a higher variability and larger vertical wind magnitudes. The values obtained from the new retrieval fit between the 445 large scale values from the monostatic retrievals and observations using HPLA radars (Hoppe and Fritts, 1995a; Fritts et al., 1990; Gudadze et al., 2019), which represents the limit for the smallest temporal scales of a few seconds (dwell time) and a spatial coverage of 3-4 km (beam diameter). Furthermore, we tested the 3DVAR+DIV retrieval with a much higher temporal resolution of 10 minutes. At this resolution the compressible solution again showed signs of numerical instability due to 450 the much sparser data coverage, which can be compensated by increasing the Lagrange multiplier for the vertical covariance constraint at the cost of smoothing some small scale structures. A similar effect occurs when increasing the vertical bin size beyond the typical 2 km. Due to the large vertical shear often associated due to large scale waves such as tides this increases the tendency for numerical instability that has a negative effect on the reliability of vertical winds.

One aspect is left that is worth to be considered. The vertical integration of the horizontal divergence, which is needed to derive 455 absolute vertical velocities, requires an initial boundary condition either at the bottom or top side of the domain depending on the integration direction. Currently this boundary is estimated assuming that the mean vertical velocity in each column above a grid cell is zero. We tested also domain means and other options. These values for the vertical velocity at the lower boundary are typically smaller than $\pm 0.2 - 0.3$ m/s for hourly winds. These vertical velocities are fairly consistent compared to other studies estimating vertical winds at altitudes between 70-80 km altitude (Straub et al., 2012; Vincent et al., 2019), which are 460 representative for a coarser temporal resolution of several hours to a day. Thus, the new 3DVAR+DIV retrieval provides more reliable values of the vertical wind variability rather than absolute wind values at a specific altitude.

Furthermore, the combined horizontal divergence and vertical velocities present a good additional diagnostics to identify 465 coherent structures in the domain area, which can be associated with gravity waves. This is often more difficult to be achieved from the horizontal winds alone without additional filtering. Zonal and meridional winds are dominated by large scale waves such as atmospheric tides that gain large magnitudes and, thus, lead to apparently smooth color maps and mostly parallel wind arrows in the images.

9 Conclusions

In this study we outlined some of the intrinsic biases that arise when inferring vertical winds from standard and multistatic 465 meteor radar observations. For this purpose, we implemented a data analysis pipeline based on least squares fits with a singular value decomposition solver for real and synthetic data. We demonstrated that even for synthetic data with zero vertical winds in



all atmospheric components including mean winds, planetary waves, tides and gravity waves, a least squares analysis results in vertical winds with a standard deviation of $\pm 2.3 \text{ m/s}$. For real atmospheric soundings the standard deviation had a value of up to $\pm 5 \text{ m/s}$. This bias is caused by the temporal and spatial sampling of meteor radars due to the random occurrence of meteors inside the beam volume. Every meteor observation is representative for a given time period determined by the decay time of 470 ambipolarly diffusing meteoric plasma and the spatial extension of the scattering volume along the trail. Thus, the apparent line of sight velocities are representative for a well-defined area inside the beam volume defined by the Fresnel scattering and for a very short time period, which is typically less than a second.

Considering these sampling aspects for typical meteor radar observations, we introduced two mathematical debiasing strategies to ensure that the estimated wind components are statistically sound solutions for a given spatial and temporal meteor 475 distribution within each time-altitude bin. We showed that the assumption of a zero vertical wind, which is often used in standard meteor radar wind analysis algorithms, is equivalent to a Tikhonov regularization of the solution for an infinitely large vertical wind component in the Tikhonov matrix. Furthermore, we introduced a more complex approach by designing a spatio-temporal Laplace filter with constraints on the error covariance, which can be seen in the broadest sense as a generalized Tikhonov regularization. This retrieval algorithm resulted in a standard deviation for the same synthetic data set of $\pm 3 \text{ mm/s}$.

480 In addition, we analyzed available multi-year meteor observations from COL and TDF and performed a statistical comparison of the inferred vertical winds with those from the UA-ICON model. The mean and statistical moments of the resultant vertical velocity distributions showed a surprisingly good agreement concerning the GCM. However, we are not able to prove, for individual measurements, the geophysical correctness of the computed vertical wind, which is why we conclude that the term 'residual bias vertical winds' still seems to be justified.

485 Although specular or transverse scatter meteor radars have been in use since decades, there is still some debate about the scattering mechanism and whether there are additional geometry effects due to the high aspect sensitivity of meteor trails. Recent quantitative simulations of reflection coefficients with a full wave scattering model have confirmed a significant change of the effective decay time and signal magnitude, which depends on the polarization of the incident radio wave and the meteor trail alignment. We were able to identify another bias in the wind magnitude when comparing forward scatter receiver data and 490 monostatic observations using CONDOR. The bias appear to be most significant below 85 km, and increases with decreasing altitude. We explain this offset by a sliding of the scattering center along the meteor trail when the meteoric plasma column is drifted by the neutral winds. Thus, meteor radars measure the Doppler velocity of the scattering center or specular point, which consists of the 'true' Doppler from the neutral winds and an apparent velocity component caused by an apparent motion of the scattering center along the trail. Source radiant mapping of two meteor showers confirmed that the Bragg vector pointing 495 direction remained unaffected. Most existing meteor radars do not provide information on the meteor orbit or trajectory and, thus, this bias poses an additional challenge to estimate mean vertical winds from monostatic meteor radars.

The new 3DVAR+DIV algorithm for multistatic meteor radar networks was implemented for routine data analysis of CONDOR and the Nordic Meteor Radar Cluster observations. This algorithm provides the first physical and mathematical consistent approach to infer vertical velocities and vertical velocity variability from multistatic networks by combining the continuity and 500 radial wind equations in the cost function. Furthermore, the 3DVAR+DIV retrieval includes new diagnostics such as horizontal



divergence and relative vorticity for each grid cell. In particular, the horizontal divergence benefits from the good measurement response of the horizontal wind components, and thus, the vertical velocities derived from the incompressible solution are also reflecting a high measurement response. The derived vertical velocities are in the range of $w = 1 - 2 \text{ m/s}$ and sometimes (3-4 sigma variance) exceed $3 - 4 \text{ m/s}$ for single grid cells of 30-by-30 km and a temporal resolution of one hour. Due to the vertical 505 integration of the continuity equation, the absolute magnitude is still subject to the assumption that the mean vertical velocity over a large vertical and spatially column is small. Although the mean absolute value still depends on the upper and lower boundary the horizontal divergence and vertical wind variability are robust quantities and provide valuable information about the spatial scales of gravity waves and their horizontal wavelength spectra.

510 *Data availability.* The data are available upon request. Please contact Alexander Kozlovsky (alexander.kozlovsky@oulu.fi) for the Nordic Meteor Radar Cluster and Alan Liu (LIUZ2@erau.edu) for CONDOR to obtain the 3DVAR+DIV retrievals. The Collm Meteor radar data can be requested from Christoph Jacobi (jacobi@uni-leipzig.de). SAAMER data from Tierra del Fuego can be requested from Diego Janches (diego.janches@nasa.gov).

515 *Author contributions.* GS and AL developed the idea to include the continuity equation into the 3DVAR algorithm. CM conceptualized the scattering schemes. AK, AL, ZQ, CJ, DJ, GL supported the implementation of the algorithms and data handling of the Nordic Meteor Radar Cluster and CONDOR as well as the Collm and Tierra del Fuego MR (SAAMER). JK, EB, SN, MT, NM, NG, and ML sustained the observations of the Nordic meteor radars and provided the data. HS guided the use of the UA-ICON data analysis and model description. AK reduced the UA-ICON data and performed the analysis including the observational filter. All authors helped with the editing of the manuscript.

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