



1 **Effective diffusivity of sulfuric acid in Antarctic ice cores**

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8

9 **Abstract**

10 Volcanic deposition of sulfuric acid in ice cores is important both for understanding past volcanic activity and for
11 synchronizing ice core timescales. Sulfuric acid has a low eutectic point, so it can potentially exist in liquid at grain
12 boundaries and veins, accelerating chemical diffusion. A high effective diffusivity would allow post-depositional
13 diffusion to obscure the climate history and the peak matching among older portions of ice cores. Here, we use
14 records of sulfate from the EPICA Dome C (EDC) ice core to estimate the effective diffusivity of sulfuric acid in
15 ice. We focus on EDC because multiple glacial-interglacial cycles are preserved, allowing analysis for long
16 timescales and deposition in similar climates. We calculate the mean concentration gradient and the width of
17 prominent volcanic events, and analyze the evolution of each with depth/age. We find the effective diffusivities for
18 interglacials and glacial maximums to be $5 \pm 2 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$, an order of magnitude lower than a previous estimate
19 derived from the Holocene portion of EDC (Barnes et al., 2003). The effective diffusivity may be even smaller if
20 artificial diffusion from the sampling is accounted for. Effective diffusivity is not obviously affected by the ice
21 temperature until about -10°C , 3000m depth, which is also where anomalous sulfate peaks begin to be observed
22 (Traversi et al., 2009). Low effective diffusivity suggests that sulfuric acid is not readily diffusing in liquid-like
23 veins in the upper portions of the Antarctic ice sheet and that records may be preserved in deep, old ice if the ice
24 temperature remains well below the pressure melting point.

25

26 **Summary**

27 We use the oldest Antarctic ice core to estimate the rate of diffusion of sulfuric acid. Sulfuric acid is a marker of past
28 volcanic activity and is critical in developing ice-core timescales. The rate of diffusion is uncertain and is important
29 to know both for selecting future ice core locations and interpreting ice-core records. We find the effective
30 diffusivity of sulfate is 10 times smaller than previously estimated, indicating the sulfuric acid signals will persist for
31 longer.

32

33



34 1 Introduction

35 Ice cores preserve unique records of past climate (e.g. EPICA, 2004; NEEM 2013; WDPM, 2013). Interpreting the
36 records requires knowledge of any post-depositional alteration such that the quantity measured in ice that was
37 deposited in the past can be interpreted as what was deposited at the surface of the ice sheet at that time (e.g. Cuffey
38 and Steig, 1998; Pasteur and Mulvaney, 2000; Bereiter et al., 2014). In general, ice cores are excellent at preserving
39 records because of the cold temperatures and the lack of processes that might disturb the record. Many records –
40 water isotopes and chemical compounds – undergo some amount of post-depositional alteration in the near surface
41 snow, firn column, or solid ice (e.g. Dibb et al., 1998; Johnsen et al., 2000; Aydin et al., 2014; Osman et al., 2017).
42 In the case of post-depositional diffusion of water isotopes, the amount of diffusion can be estimated and used for
43 paleoclimate reconstruction (Gkinis et al., 2014; Kahle et al., 2020); however, for most soluble impurities and
44 atmospheric gases, the primary goal of understanding post-depositional processes is to ensure accurate
45 reconstructions.

46 Here we investigate the diffusivity of sulfuric acid using a record of sulfate from the EPICA Dome C (EDC) ice core
47 (Severi et al., 2007). We focus on sulfate for two reasons: 1) the deposition of sulfate from volcanic events provides
48 distinct features to identify and assess the change of shape with depth/age and 2) the importance of volcanic
49 matching in synchronizing timescales (e.g. Ruth et al., 2007; Severi et al., 2007; Fujita et al., 2015; Buizert et al.,
50 2018; Winski et al., 2019; Svensson et al., 2013; 2020; Sigl et al., 2013, 2015, 2021). Stable soluble impurities, such
51 as sulfate, can remain as distinct layers in deep ice cores (e.g. Zielinski et al., 1997; Fujita et al., 2015); however,
52 chemical peaks due to post-depositional alteration in deep, and warm, ice have also been identified (e.g. Traversi et
53 al., 2009; Tison et al., 2015). Sulfuric acid has a eutectic temperature of -75°C , such that if it is present at the grain
54 boundaries and triple junctions, a liquid vein network may exist even in the cold near surface ice of the East
55 Antarctic plateau (Dash et al., 2006). Some studies indicate that sulfuric acid is located at the triple junctions
56 (Fukazawa et al., 1998; Mulvaney et al. 1988) particularly in specimens with high concentrations of sulfate (Barnes
57 and Wolff, 2004), although there are questions about whether the sample preparation caused the sulfuric acid to
58 concentrate there; the sulfate may be in the form of salt micro-inclusions in grains (Ohno et al., 2005). Thus,
59 whether sulfuric acid is primarily located at the veins, grain boundaries, or within the ice lattice remains an open
60 question and the mechanisms driving diffusion are uncertain.

61 Addressing the diffusion of sulphate in veins, Rempel et al. (2001) showed that impurity fluctuations may retain
62 their amplitude but migrate towards warmer temperatures and away from the ice the impurity was deposited with;
63 however, Ng (2021) showed that including the Gibbs-Thompson effect due to the grain curvature may cause these
64 peaks to diffuse rapidly and therefore not retain their amplitude nor migrate. The sulfate and electrical conductivity
65 measurements (ECM) of Antarctic ice cores also show reduced amplitude (Barnes et al., 2003, hereafter B03; Fujita
66 et al., 2015; Fudge et al., 2016) in older ice. B03 found an effective diffusivity of sulfate to be $4.7 \times 10^{-8} \text{ m}^2 \text{ a}^{-1}$ at -
67 53°C for Holocene ice at EPICA Dome C (EDC), which is two orders of magnitude greater than the self-diffusion of
68 ice at -53°C ($3 \times 10^{-10} \text{ m}^2 \text{ a}^{-1}$), which they considered an upper limit for the diffusivity of sulfate. They proposed two
69 mechanisms of connected and unconnected vein networks to explain the effective diffusivity. Fudge et al. (2016)
70 estimated an effective diffusivity $2.2 \times 10^{-8} \text{ m}^2 \text{ a}^{-1}$ at -30°C at the West Antarctic Ice Sheet Divide (WDC), a
71 considerably smaller value despite a warmer temperature.

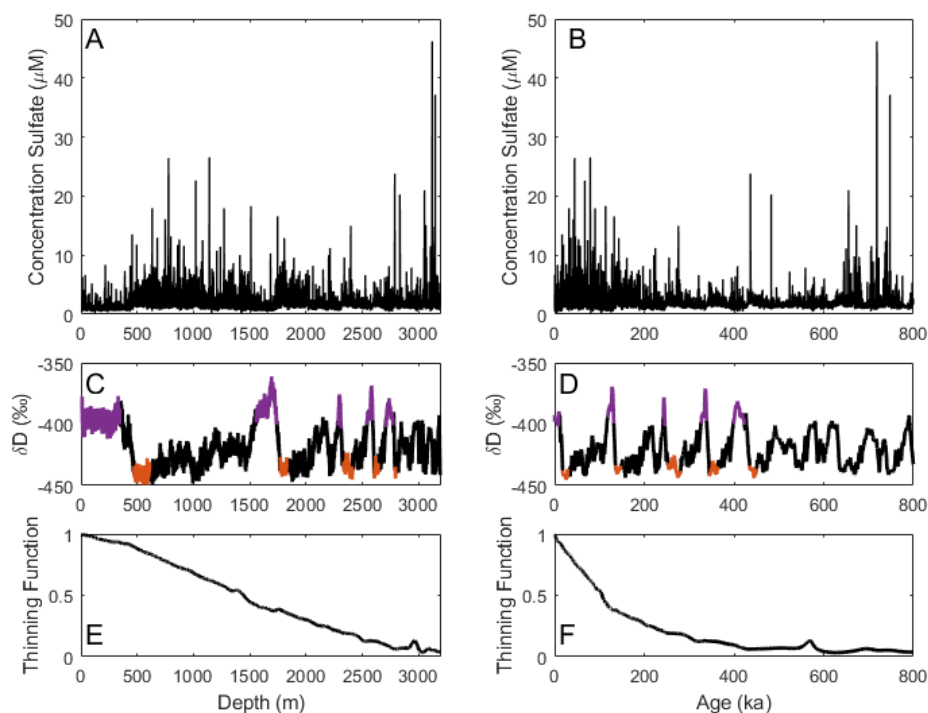
72 Here we assess the effective diffusivity of sulfate at EDC with two methods. First, we follow B03 in calculating the
73 scaled mean gradient. We extend their analysis to the full depth of the EDC ice core. We focus on only the
74 interglacial and glacial maximum time periods to reduce the influence of the climate at the time of deposition.
75 Second, we identify the widths of volcanic peaks throughout the core and evaluate the widening of them with a
76 numerical diffusion model.

77

78 2 Methods

79 2.1 EDC sulfate data set

80 The sulfate record for EDC was measured with fast ion chromatography (Severi et al., 2015). The sampling
81 resolution is approximately 4 cm in the Holocene and increases to approximately 2 cm by 44ka where it remains for
82 the rest of the core. The EDC sulfate record has been widely used for volcanic matching (e.g. Severi et al. 2007;
83 Severi et al., 2012; Buizert et al., 2018) and calculating volcanic forcing (e.g. Sigl et al., 2015; 2022). We assume
84 that vertical variations in the sulfate record greatly exceed lateral variations such that the diffusion can be treated as
85 a one-dimensional problem.



86
 87 *Figure 1: Sulfate (A,B), Deuterium (C,D) and the thinning function (E,F) plotted by depth (A, C,E) and age (B,D,F).*
 88 *The most-recent five interglacials are plotted in purple and glacial maximums are plotted in orange on the*
 89 *deuterium plots (C,D). Earlier interglacials do not reach the “warm” levels of most recent 5.*

90

91 2.2 Scaled mean gradient

92 Diffusion will reduce the amplitude of variations and broaden peaks in the sulfate record. To calculate the effective
 93 diffusivity, we first need a quantitative measure of the sulfate variability and then we need an expression for
 94 inferring the effective diffusivity from the change in the sulfate variability. One method for calculating the
 95 variability is uses the difference between successive points over fixed depth intervals. The mean difference between
 96 successive points will decrease more quickly with age for greater effective diffusivities. This technique has the
 97 advantage of using all the available data in a section of interest and was first described by B03. In B03, this mean
 98 difference was termed the scaled mean gradient. As the name scaled mean gradient suggests, B03 took the additional
 99 step of scaling the sulfate data to account for variations due to the climate at deposition. B03 investigated only the
 100 Holocene data. When we extended the analysis to glacial climates, we found that the scaling did not account for
 101 differences between different climates at deposition. We therefore decided to only compare similar climate states:
 102 the interglacials to interglacials and the glacial maximum periods to glacial maximums. We also found that the
 103 inference of effective diffusivity was not sensitive to whether we scaled the data or left it unscaled. We use the
 104 scaling, as described by B03, to be consistent with previous results. Below we focus on two primary equations – for
 105 the scaled mean gradient and for the effective diffusivity - and provide the supporting equations to explain the
 106 calculations performed. A full description can be found in B03.

107 The first primary equation is for calculating the scaled mean gradient (B03, Eq. 1) which is an estimate of:

$$108 \bar{m}_z = \frac{1}{\Delta z} \sum_z^{z+\Delta z} |\delta c'(z)| \quad (1)$$



109 where c' is the scaled sulfate profile, z is depth, Δz is a fixed depth interval (10 m), $\delta c'(z)$ is the difference between
 110 successive points. The sulfate profile is scaled to account for differences in the amplitude of variations caused by
 111 variations in climate rather than due to diffusion:

$$112 \quad c'(z) = \left(\frac{c(z) - c_0}{\bar{c}_z - c_0} \right) (\bar{c} - c_0) + c_0 \quad (2)$$

113 where c_0 is calculated as the intercept of the total area under the peaks plotted against the mean concentration. We
 114 perform the scaling to be consistent with the B03 methodology.

115 The second primary equation finds the effective diffusivity using the scaled mean gradient B03, Eq. 9):

$$116 \quad D_{eff} = -\frac{1}{k^2 t} \log \left(\frac{m(z)}{m_0} \right) \quad (3)$$

117 where D_{eff} is the effective diffusion coefficient in solid ice, k is the wave number, t is time, and m_0 is the initial
 118 scaled mean gradient. This equation is derived by noting that sulfate concentration varies primarily with depth such
 119 that horizontal variations can be neglected, yielding a one-dimensional process in depth:

$$120 \quad \frac{\partial c'}{\partial t} = \chi(t) D_{eff} \frac{\partial^2 c'}{\partial z'^2} \quad (4)$$

121 where $\chi(t)$ accounts for the vertical thinning including densification of firm. This differential equation can be solved
 122 following Johnsen et al. (2000) by relating a diffusion length, l , to $D_{eff}(t)$:

$$123 \quad \frac{dl^2}{dt} - 2\dot{\epsilon}_z(t)l^2 = 2D_{eff}(t) \quad (5)$$

124 where $\dot{\epsilon}_z(t)$ is the vertical strain rate. Here, we first destrain the ice based on the thinning function of the Antarctic
 125 Ice Core Chronology 2012 (AICC2012; Bazin et al., 2013, Veres et al., 2013) such that $\dot{\epsilon}_z(t)$ has been removed.
 126 The amplitude can then be defined in relation to the initial amplitude – calculated for the Holocene and Last Glacial
 127 Maximum (LGM); age durations for intervals are given in Appendix A. Assuming a harmonic cycle with a wave
 128 number, k

$$129 \quad H = H_0 \exp \left(-\frac{1}{2} k^2 l^2 \right) \quad (6)$$

130 where H is the amplitude after some time, and H_0 is the initial amplitude. The wave number, k , is related to the
 131 mean peak width

$$132 \quad k \sim \frac{2\pi}{\bar{w}} \quad (7)$$

133 where \bar{w} is the mean peak width for the full Holocene or LGM profile and can be calculated from the mean absolute
 134 gradient as

$$135 \quad \bar{w} = \frac{4}{\bar{m}_{total}} \sum_0^{z_{total}} |c'(z) - \bar{c}| \delta z \quad (8)$$

136 where \bar{m} is the mean absolute gradient and z_{total} is the full depth for either the Holocene or LGM.

137

138 **2.3 Width of volcanic peaks and numerical modeling of diffusion**

139 An alternate technique for measuring the evolution of chemical signals is to identify individual volcanic peaks (e.g
 140 Fudge et al., 2016). We identified peaks based on their prominence - the amplitude between the peak and the nearest
 141 local minimum in both up- and down-core directions with the findpeaks functions in Matlab. Finding the
 142 prominence allows a standardized approach that accounts for different background levels of sulfate deposition that
 143 occur in different climate states. Our approach compared well with studies that used an exceedance of the median
 144 absolute deviation (Traufetter et al., 2004; Sigl et al., 2015; Nardin et al., 2020; Cole-Dai et al., 2021) in sliding
 145 windows.

146 The volcanic events were identified in age on the AICC2012 timescale (Bazin et al., 2013; Veres et al., 2013). The
 147 findpeaks algorithm calculates the width of the peak at the half-maximum of the sulfate peak amplitude. The width
 148 is in years and can be converted to depth using the depth-age relationship. As for the scaled mean gradient, we
 149 compare the interglacial and glacial maximum periods separately. For each period, we use the largest 25 events.



150 The effect of diffusion is modeled with a one-dimensional numerical model (e.g. Eq. 4). The model is described in
 151 Fudge et al. (2016). At each time step, the diffusion is calculated first, and then the ice is thinned. The amount of
 152 thinning at each time step is calculated from the thinning function (Veres et al., 2013; Bazin et al., 2013). The ice
 153 temperature and effective diffusivity can be varied at each time step, although we only show results using a constant
 154 effective diffusivity here. We model the evolution of a gaussian peak. The initial width of the peak is determined by
 155 the sulfate data in the Holocene (0.128 m) and LGM (0.135 m). The initial thickness are similar but, because the
 156 accumulation rate in the LGM is a factor of 2 lower, the duration of the events are twice as long in the LGM. The
 157 amplitude of the peak is three times the background concentration although in practice the amplitude is not critical
 158 because we calculate the relative change.

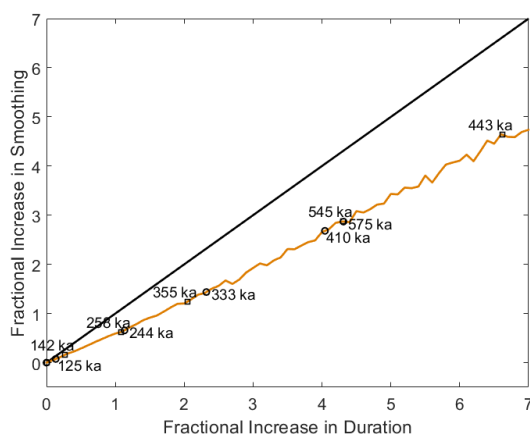
159

160 2.4 Estimating artificial diffusion

161 The calculation of effective diffusivity may be affected by artificial diffusion due to the measurement duration. The
 162 sampling resolution in depth for the sulfate record is approximately 4 cm for the first 45 ka and 2 cm for older ages.
 163 The increasing thinning with depth results in each measurement averaging over a longer duration which can
 164 artificially smooth the mean gradient by failing to fully resolve peaks and troughs. If the sampling resolution is short
 165 relative to the signal, such as a multi-year volcanic event with seasonal sampling, the artificial smoothing will be
 166 negligible. On the other hand, if the sampling resolution is long relative to the signal, significant smoothing will
 167 occur.

168 We estimate the amount of artificial smoothing for older ice by resampling the Holocene data. We select the period
 169 2.3 to 11.3 ka which has relatively even sampling durations and is well resolved due to the small amount of vertical
 170 thinning. We resample the data at increasing durations and compare the mean gradients of the resampled data with
 171 the mean gradient of the original data. The ratio of the original mean gradient to the resampled mean gradient is the
 172 size of the correction that should be applied (Figure 2). The amount of artificial smoothing increases at about 70% of
 173 the rate of the increase in duration.

174 The potential artificial diffusion for each interglacial and glacial maximum periods is shown in Figure 3. The Last
 175 Interglacial Glacial (LIG) and Penultimate Glacial Maximum (PGM) sampling interval is 2cm, approximately half
 176 that of the Holocene and LGM, such that the shorter sampling interval largely offsets the layer thinning. For older
 177 ages, the artificial smoothing becomes larger, reaching a maximum of 440% for the glacial period centered at 443 ka
 178 (Marine Isotope Stage 12). This is an outlier, due to the small annual layer thickness during that period, but
 179 highlights that the correction for artificial smoothing becomes quite large and potentially dominates the estimate of
 180 effective diffusivity.



181

182 *Figure 2: Estimate of the impact of artificial smoothing due to sample size in each interglacial or glacial maximum*
 183 *period. The amount of artificial smoothing (brown) increases more slowly than the increase in average sample*
 184 *duration. The one-to-one line is shown in black.*



185

186 **3 Results**

187 **3.1 Effective diffusivity from the scaled mean gradient**

188 The characteristics of the sulfate record vary with the climate state at deposition. The sulfate deposited in
 189 interglacials has a lower concentration than during the glacial periods. This is most evident in Figure 1A where the
 190 increase in sulfate concentration is visible at about 450 m. Comparing the sulfate gradients among different climate
 191 states is problematic. We found that scaling the mean gradients did not fully remove the differences due to the
 192 climate at deposition (section 2.1). Therefore, we compare the sulfate gradients among periods with similar climate
 193 characteristics, specifically the interglacial periods and the glacial maximum periods separately.

194 We follow B03's methods (Section 2.2) and repeat the calculations for the initial sulfate parameters for the
 195 Holocene (0-11.3 ka). We find c_0 is 0.62 compared to the B03 value of 0.54 and the same value of \bar{w} of 0.19 m. The
 196 small difference is likely due to updates in the sulfate data set and thinning function. We also find initial parameters
 197 for the LGM (18-30 ka). For both the Holocene and LGM, we assume that these initial values at deposition were the
 198 same for previous interglacial and glacial maximum periods (Table 1).

199 We first calculate the effective diffusivity using the most recent glacial cycle by comparing the Holocene and LIG,
 200 and the LGM and PGM. The difference between the interglacial and glacial maximums of the most recent cycle
 201 provide a useful duration (~120 ka) for diffusion to have operated without the ice having become too greatly
 202 thinned. In addition, the sample resolution in depth was smaller for the older periods which yielded a relatively
 203 small increase in the duration (age span) of each sample compared to the most recent periods.

204 We use Eq. 3 to calculate the effective diffusivity, where m_z is the median scaled mean gradient of the LIG or PGM
 205 and m_0 is the median scaled mean gradient of the Holocene or LGM. The value of k is calculated from the Holocene
 206 or LGM data. The resulting effective diffusivities are $6.1 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ (Table 1) using the interglacials, and 5.1×10^{-9}
 207 $\text{m}^2 \text{ a}^{-1}$ using the glacial maximums.

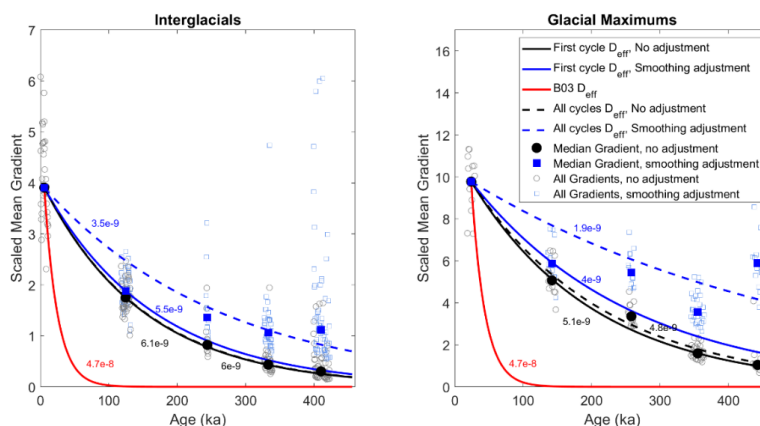
208 For the LIG, we can calculate an effective diffusivity corrected for artificial smoothing by increasing the mean
 209 absolute gradient by the 7.5% calculated above; this decreases the effective diffusivity to $5.5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. For the
 210 PGM, the 16% adjustment for artificial smoothing results in a decrease in the effective diffusivity to $4.0 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$.
 211 The calculated effective diffusivities are shown in Table 1 and Figure 3.

212

213 Table 1: Effective diffusivities calculated using the two most recent interglacial and glacial maximum periods

	Interglacials		Glacial Maximums	
c_0	0.62		-0.4	
w	0.19		0.19	
Median Gradient Initial	3.9	Holocene	9.71	LGM
Median Gradient of next glacial cycle	1.75	LIG	5.18	PGM
k	33.3		32.7	
Effective Diffusivity ($\text{m}^2 \text{ a}^{-1}$)	6.1×10^{-9}		5.1×10^{-9}	
Adjustment	7.5%		16.0%	
LIG Adjusted Median Gradient	1.85		5.57	
Adjusted Effective Diffusivity ($\text{m}^2 \text{ a}^{-1}$)	5.5×10^{-9}		4.0×10^{-9}	

214



215
 216 *Figure 3: The effective diffusivity, D_{eff} , is labeled in units of $\text{m}^2 \text{a}^{-1}$ for the different inferences from the scaled mean*
 217 *gradients in interglacial and glacial maximum periods. Scaled mean gradients for 10 m sections are shown by open*
 218 *symbols; the median scaled mean gradients for each interglacial or glacial maximum are shown by solid symbols.*
 219 *Solid lines show inference of effective diffusivity from the first glacial cycle; dashed lines are adjusted for artificial*
 220 *smoothing. Black lines and symbols have no adjustment for artificial smoothing; blue lines and symbols are*
 221 *adjusted for artificial smoothing. Red line is previous estimate of effective diffusivity from B03 using Holocene data.*

222
 223 The scaled mean gradients from the past five glacial cycles can also be used to estimate the effective diffusivity. We
 224 perform a linear fit to the log of the medians for each period for both the unadjusted and adjusted values (Figure 3,
 225 Table 2); we fit the medians instead of all the points to account for the different number of points in each
 226 interglacial. The fit for the past five glacial cycles finds similar effective diffusivities for the unadjusted values for
 227 both the interglacial and glacial maximum periods; however, the fit finds significant smaller effective diffusivities
 228 for the adjusted values. The inferred interglacial effective diffusivity is approximately 40% smaller than that inferred
 229 for the Holocene-LIG, while the glacial maximum value is approximately 55% smaller.

230 For both the most recent glacial cycle and the past five glacial cycles, we calculate the uncertainty of the effective
 231 diffusivity based on the 95% confidence interval of the linear fit to the log of the scaled mean gradients; instead of
 232 using the medians, we use the individual data points. We report this as a \pm value by taking the 95% confidence range
 233 and dividing by two. The uncertainties are given in Table 2.

234 Table 2

	Scaled mean gradient				Volcanic width	
	Unadjusted	Adjusted	Unadjusted	Adjusted	N/A	
	Last cycle	Last cycle	Past 5	Past 5	Last cycle	Past 5
Interglacial	6.1±0.7	5.5±0.7	6±0.5	3.5±1.9	6.0±NA	1.6±NA
Glacial Maximum	5.1±0.8	4.0±0.8	4.8±0.9	1.9±2.5	5.5±NA	4.0±NA

235 Effective diffusivity values (D_{eff}) in units of $10^{-9} \text{m}^2 \text{a}^{-1}$ for interglacial and glacial maximum periods. Italics
 236 indicate two outliers which help define the lower bound.

237

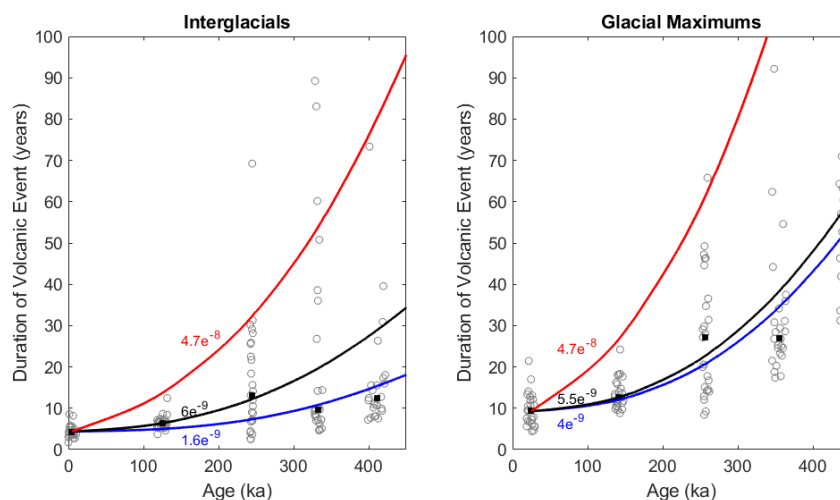
238 3.2 Effective diffusivity using widths of volcanic events

239 The duration of the 25 largest volcanic events, and their medians, for the interglacial and glacial maximums of the
 240 past 5 glacial cycles is shown in Figure 4. The duration of the largest events increases with age although there is a lot
 241 of scatter. The median duration of a volcanic event in an interglacial increases relatively little from about 5 years in



242 the Holocene to about 10 years for interglacials older than 200 ka. The glacial maximums, in contrast, show a large
243 increase in duration, from 10 years in the LGM to over 50 years at ~450 ka. The widths of the largest volcanic
244 events are less sensitive to artificial smoothing due to the sample duration than the scaled mean gradient because the
245 volcanic signal is long relative to the sampling interval; therefore, we do not add a correction for artificial smoothing
246 (section 2.4).

247 The numeric model (section 2.3) is used to infer the effective diffusivity by minimizing the misfit to the median
248 widths of each period. As with the scaled mean gradients, we infer a constant effective diffusivity and start with only
249 most recent glacial cycle. For the LIG-Holocene, an effective diffusivity of $6.0 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ is the best fit; for the
250 LGM-PGM, the best fit is $5.5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. The best fits for all five glacial cycles are $1.6 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ for the
251 interglacial periods and $4 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ for the glacial maximums. We do not calculate an uncertainty for the
252 inferences and will discuss the overall uncertainty in the inference of effective diffusivity in the following section
253 (3.3).



254

255 *Figure 4. Duration of the 25 largest volcanic events (grey circles) for each interglacial and glacial maximum*
256 *period. Medians for each period are shown as black squares. Solid lines are the modeled volcanic event durations*
257 *using constant effective diffusivities. Black lines are best fit to Holocene-LIG or LGM-PGM; blue lines are best fit to*
258 *all 5 interglacial or glacial maximum periods; red is the B03 estimate from Holocene data.*

259

260 3.3 Evaluating uncertainty in effective diffusivity

261 The effective diffusivities inferred from the most recent interglacial and glacial maximum periods agree well
262 between the scaled gradient method and the volcanic width method (and the adjusted and unadjusted for the scaled
263 mean gradient method). The inferred values range from $6.1 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ to $4.0 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. The inferred values
264 using the past five glacial cycles agree less well, with a minimum inference of $1.9 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. The uncertainties
265 calculated for each inference are smaller than the differences among the estimates. This indicates that the calculated
266 uncertainties are too small. Therefore, we discuss qualitative uncertainty bounds.

267 There are twelve different inferences of the effective diffusivity (Table 2). Ten of these (83%) fall within the range
268 of $6.1 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ to $3.5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. The other two values are both smaller, reaching a minimum of $1.6 \times 10^{-9} \text{ m}^2$
269 a^{-1} . Seven of the inferred diffusivities are between 6.1 and $4.8 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$, and this is the range that we suggest is
270 the most likely value for the effective diffusivity. Those seven estimates have a mean of $5.5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$; all of the
271 other five estimates are smaller. For simplicity, we suggest of a value of $5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ given the bias towards
272 smaller values and the lack of certainty which does not warrant an additional significant figure in the estimate.

273 The Dome Fuji ice core is most similar to Dome C, with a similar ice thickness, temperature profile, depth-age
274 profile, and modern accumulation rate. While we were unable to find a publicly available sulfate data set to perform



275 a similar analysis, we were able to use the ECM record (Fujita et al., 2002). Calculating the scaled mean gradient as
 276 for EDC, the inferred effective diffusivity using the Holocene and LIG was $6.6 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$, and for the LGM and
 277 PGM was $3.2 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. While ECM and sulfate records respond similarly to volcanic deposition of sulfuric acid,
 278 the ECM typically also responds to neutralization of acids where dust is in higher concentrations. Therefore, we
 279 have more confidence in the effective diffusivities inferred from the EDC sulfate record but are encouraged that the
 280 Dome Fuji results are in reasonable agreement. This suggests that the inferred effective diffusivity from EDC is
 281 applicable to similar East Antarctic sites.

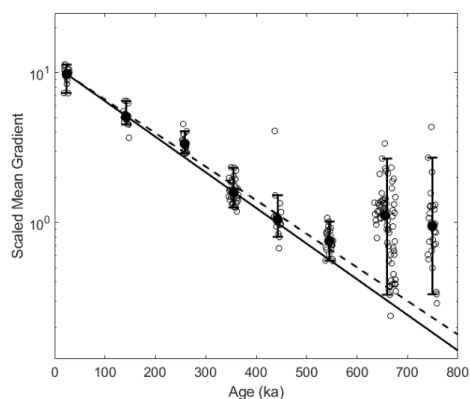
282 Selecting an uncertainty value is challenging. The correction for artificial smoothing appears too large for the older
 283 ice, which may be biasing some of the effective diffusivity estimates too low; however, we cannot rule out that the
 284 effective diffusivity is smaller than our main range of estimates. The largest estimate of the effective diffusivity is
 285 from the interglacial estimate of the most recent cycle. This is the estimate we are most confident in since it has the
 286 best resolved data and the smallest potential corrections. While do not have firm constraints on the uncertainty, we
 287 suggest using $5 \pm 2 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ based on the available estimates.

288

289 3.4 Deep ice older than 450 ka

290 We have thus far in this paper restricted the analysis to the most recent five glacial cycles because of the lack of
 291 interglacials with “warm” enough water isotope values to compare with similar characteristics at deposition. The
 292 glacial maximum water isotopes, however, are of similar values (Figure 1), allowing three more periods to be
 293 considered. Figure 5 shows that the scaled mean gradients decrease with age through 545 ka; however, by 656 ka the
 294 median scaled mean gradient has increased rather than continued to decrease. This occurs at about 3000 m depth.
 295 The variability also increases for the oldest two glacial maximum periods. The scaled mean gradients have not been
 296 corrected for artificial smoothing.

297 The increase in scaled mean gradient for the final two glacial maximum periods is at first counter-intuitive. An
 298 increase in effective diffusivity should reduce the scaled mean gradient. The increase in the scaled mean gradient
 299 likely occurs because sulfate has become mobile enough to form peaks unassociated with the sulfate concentrations
 300 at deposition. Traversi et al. (2009) noted that below about 2800 m depth there were sharp spikes in sulfate with
 301 anomalous chemical compositions of low acidity and high magnesium. The scaled mean gradient at 545 ka shows a
 302 slight increase relative to the expectation, which may indicate some contribution of anomalous peaks; however, the
 303 increase is not substantial and suggests that if impurity interactions have begun, they remain limited.



304

305 Figure 5. The scaled mean gradients for eight glacial maximum periods are shown as open black circles. Analysis
 306 includes only glacial periods because the water isotope values for all glacial maximums are similar while the older
 307 interglacials are not as comparable (Figure 1). Medians of each glacial maximum are shown in solid black and the
 308 range, excluding the largest and smallest in each time period, are shown with black whiskers. The most recent five
 309 glacial maximums are the same as in Figure 3B. The black solid and dashed lines the same as in Figure 3B; note the
 310 log y-scale.

311



312 4 Discussion

313 4.1 Relation with temperature and grain size

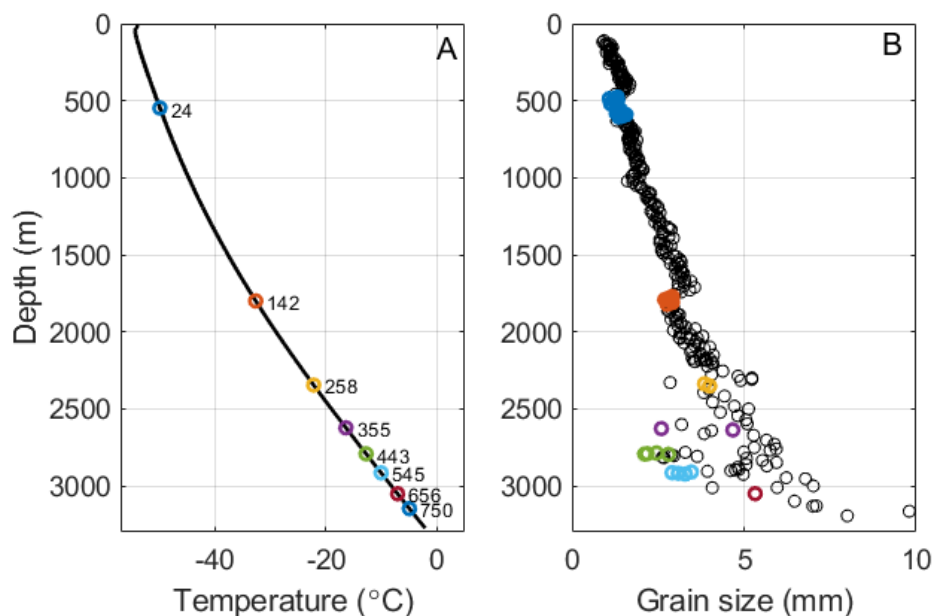
314 The scaled mean gradient and the broadening of the largest volcanic peaks yield a similar result: the effective
315 diffusivity of sulfate on timescales of hundreds of thousands of years is approximately $5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ throughout
316 most of the ice sheet until the ice warms near the bed. Our inferred value of the effective diffusivity is an order of
317 magnitude smaller than that inferred from the Holocene data at EDC (B03). This result is consistent with the ability
318 to find volcanic matches between EDC and Dome Fuji for the past 216 ka (Fujita et al., 2015) which would not be
319 possible if the volcanic signals diffused at the previous estimate.

320 Two different complications may affect our inference of the effective diffusivity. The first complication is artificial
321 smoothing due to averaging over longer durations of time in samples of older ice. We estimated the impact of
322 artificial diffusion (section 2.3) and showed that the inferred effective diffusivity becomes smaller if it is included
323 (section 3.1); however, the artificial smoothing appears to be overestimated in the older ice. The second
324 complication is rearrangement of impurities in deeper and warmer ice (Traversi et al., 2009), leading to an increase
325 in the scaled mean gradient due to anomalous sulfate peaks. The rearrangement of impurities increases the scaled
326 mean gradient and would decrease the inferred effective diffusivity. This process is thought to begin at about 2800
327 m depth (Traversi et al., 2009) and become more pronounced with depth as the ice warms and exceeds 10°C ; the
328 bottom 60m of ice are sufficiently altered that no interpretable climate records are preserved (Tison et al., 2015).
329 The rearrangement of impurities is unlikely to affect the more recent glacial cycle in the ice in the upper part of the
330 ice sheet. Both the artificial smoothing and the possible anomalous peaks are least likely to affect the estimate of
331 effective diffusivity from the most recent glacial cycle (e.g. Holocene-LIG and LGM-PGM), and thus we believe
332 this period provides the most reliable estimate. The effective diffusivity inferred from the most recent glacial cycle
333 fits the first five glacial cycles well (Figure 3) and thus suggests that these two potential complications are not
334 significant until older (deeper) ice. Beyond the past 5 glacial cycles, post depositional alteration of the chemical
335 impurities (Figure 5) becomes the dominant feature of the record and the concept of an effective diffusivity is no
336 longer useful.

337 The effective diffusivity is relatively consistent between interglacial and glacial maximum periods, suggesting that it
338 is not strongly affected by the total impurity concentration in the ice. There also is not a significant impact of
339 temperature over the range of -55°C at the surface to about -10°C (3000 m depth). The borehole temperature profile
340 at EDC (Buizert et al., 2021) is shown in Figure 6A. The rate of diffusion does not appear to change with age and
341 temperature, until the last two (possible the last three) glacial maximums (Figure 5) where multiple chemical species
342 have become mobile enough to form anomalous sulfate peaks. As suggested by Traversi et al. (2009), there appears
343 to be a transition to a connected vein network which is not present at shallower depths and colder temperatures.
344 While the transition is not a simple decrease in the variability in the sulfate record, it nonetheless suggests a zone
345 where sulfate starts becoming mobile and the primary signal switches from a record of the atmospheric conditions at
346 the time of deposition to a record of post-depositional alteration.

347 The modeling of diffusion of sulfate (e.g. Rempel et al., 2001; 2002; B03; Ng et al., 2021) has focused on sulfate in a
348 vein network. However, the location of sulfate is not well established and will affect the ability for sulfate to diffuse.
349 Our estimate of the effective diffusivity ($5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$) is of the same order as the self-diffusion of ice ($1 \times 10^{-9} \text{ m}^2$
350 a^{-1} at -40°C), suggesting that sulfate may reside in grains at colder temperatures. The self-diffusion of ice should be
351 regarded as an upper limit (B03) for the diffusivity of sulfate, and is still smaller than our estimate which suggests
352 some sulfate likely also resides in grain boundaries and veins as well as in the crystal lattice.

353 The mechanisms for diffusion of sulfate have been linked to grain size and growth (Rempel et al., 2001; B03).
354 Figure 6B shows that the grain size increases relatively little, particularly for the glacial maximum periods, likely
355 due to impurities pinning grain boundaries (Durand et al., 2006). B03 developed two mechanisms, for either a
356 connected or disconnected vein network, that were ultimately driven by grain growth. If net grain growth drives the
357 solute movement, the relatively constant rate of grain growth at EDC until ~ 3000 m may explain why there is little
358 noticeable impact of temperature. At depths below ~ 3000 m and temperatures above -10°C , the grain size starts to
359 increase more rapidly. A temperature of -10°C is unlikely to be a rigid threshold, as post-depositional processes are
360 likely to start above this temperature (e.g. Traversi et al., 2009), and is more likely to be a point at which the
361 impurity interactions accelerate. The larger grains and corresponding larger veins likely allow greater impurity
362 mobility. At EDC, this increase in mobility allows the sulfate to react with magnesium and other impurities, creating
363 peaks unrelated to the concentration of sulfuric acid deposited at the surface.



364

365

366 Figure 6: A) Temperature profile of EDC borehole (Buizert et al., 2021) with the glacial maximum periods plotted.

367 B) Grain size of EDC (Durand et al., 2006) with measurements within the glacial maximum periods plotted. Note

368 that the last glacial maximum period has no grain size measurements within the time period.

369

370 4.2 Implications for ice older than 1 Ma

371 The locations of ice cores targeting 1.5 Ma ice (e.g. Fischer et al., 2013) will have similar site characteristics to
372 Dome C and Dome Fuji. The European, Australian, and Japanese projects are located near the existing EDC and
373 Dome Fuji cores (Lilien et al., 2021; Karlsson et al., 2018). The US Center for Oldest Ice Exploration (COLDEX) is
374 searching for a location in the interior of the East Antarctic plateau while the Chinese effort is at Dome A and the
375 Russian project may be near Vostok. The low inferred effective diffusivity of sulfate suggests the volcanic climate
376 records will be well preserved in these deep ice cores for at least the past few hundreds of thousands of years.

377 However, the preservation of impurity signals for ice older than a few hundreds of thousands of years may depend
378 on the basal temperature and whether the old ice has experienced considerable time at temperature above -10°C
379 allowing large grains and likely connected vein networks to exist. Thus, ice core sites where the old ice remains
380 cold, either due to a basal temperature well below the melting point or a basal ice layer that keeps the old ice farther
381 from the bed, are likely to recover longer impurity records indicative of past climate variations.

382 The other type of location that has preserved ice older than 1 Ma is blue ice ablation areas (Higgins, 2015; Yan,
383 2019). The ice studied to date is not in stratigraphic order and thus records snapshots of past climate that average
384 over an unknown interval of time. This ice has a complicated flow history as revealed by the stratigraphic
385 disturbances. The old ice has thus far been found between 100 m and 200 m deep, where the ice temperature is close
386 to the surface temperature of -30°C for the Allan Hills. While the depths the ice reached on the path to its current
387 location is not known, the maximum depth imaged upstream is about 1200 m (Kehrl et al., 2019). It is likely that the
388 old ice in blue ice regions has never gotten warm enough (e.g. -10°C) for the post-depositional processes found in
389 the deep ice of EDC to become active.

390

391



392 4.3 Application to other ice cores

393 It is more difficult to compare the results from EDC to Antarctic ice cores with different site characteristics. We
394 have investigated WAIS Divide (Fudge et al., 2016; McConnell et al., 2017), SPICEcore (Winski et al., 2019),
395 EDML (Severi et al., 2012), and Talos Dome (Severi et al., 2015). Neither WAIS Divide nor SPICEcore spans a full
396 glacial cycle and the sulfate data sets for EDML and Talos Dome do not extend past the last glacial period. This
397 limits the ability to compare the evolution of the scaled mean gradient or volcanic widths for the same climate at
398 deposition. We attempted to compare the scaled mean gradient and widths of volcanic events deposited in different
399 climate states for EDC but could not find a reliable way to remove the influence of the climate at deposition. Thus,
400 extending this analysis to these other ice cores was not possible.

401 Fudge et al. (2016) used the width of volcanic events at WDC to estimate the effective diffusivity but did not
402 account for the different climate states at deposition. They found an effective diffusivity of $2.2 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ at -30°C
403 using a temperature-dependent relationship; this effective diffusivity is half that of B03, but four times greater than
404 inferred here. While we believe their methodology is not as robust as presented here, it is still worthwhile to
405 consider why the two estimates might differ. The depth-age relationships are very different. The accumulation rate
406 and vertical thinning is much greater at WAIS Divide; for instance, the LGM ice at EDC is 87% of its original
407 thickness while at WAIS Divide it is 31% of the original thickness. The higher rate of vertical thinning at WAIS
408 Divide may lead to a higher effective diffusivity, particularly if the effective diffusivity is linked to grain growth and
409 nucleation processes (B03). The grain size at WAIS Divide (Fitzpatrick et al., 2014) is smaller than EDC for all
410 depths greater than 500 m, so net grain growth would not be the cause of a higher effective diffusivity at WDC.

411 The WDC volcanic widths were fit with an effective diffusivity with an assumed temperature dependence. We note
412 that the temperature dependence is not well constrained by the data, given the uncertainties in the volcanic widths.
413 Thus, WAIS Divide neither supports nor refutes the lack of temperature dependence of the effective diffusivity at
414 temperature below $\sim -10^\circ\text{C}$. WDC is in agreement that above -10°C the effective diffusivity seems to change. Fudge
415 et al. (2016) did not identify peaks at ages older than 52 ka and the synchronization of WDC with other East
416 Antarctic cores (Buizert et al., 2018) stopped at 58 ka, despite a bottom age of 67 ka at WDC. These depths and ages
417 correspond approximately to a -10°C temperature in the ice sheet (Cuffey et al., 2016). No anomalous sulfate peaks
418 were observed at WDC, which may be due to the much shorter duration that the ice has spent at the warm
419 temperature (i.e. 10s of ka at WDC compared to 100s of ka at EDC). Future work is needed to better constrain the
420 effective diffusivity of sulfate in ice cores with higher accumulation rates than the interior of the East Antarctic
421 plateau.

422

423 4.4 Implications for existing timescales and estimates of volcanic forcing

424 The effective diffusivity of $\sim 5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$ is low enough that volcanic events in the Holocene are not significantly
425 impacted by diffusion. After 10 ka, the amplitude is reduced by only 6%. Since detection algorithms (Cole-Dai et
426 al., 2021; Sigl et al., 2015) take the background level into account and calculate the sulfate flux onto the ice sheet
427 over multiple years, there is likely to be little change in the flux estimates. The flux estimates will not be
428 significantly impacted until the peaks have diffused enough to blend into and increase the background levels.
429 Reconstructions of atmospheric sulfate from bipolar ice core records now extend through the Holocene (Sigl et al.,
430 2022) and may be extended farther as volcanic matches in the glacial period are identified (Svensson et al., 2020).
431 The impact of diffusion on the volcanic flux will be more important to consider for older ages.

432 Our results also show that diffusion is unlikely to substantially impact the synchronization of Antarctic timescale
433 through the last glacial period (e.g. Ruth et al., 2007; Severi et al., 2012; Buizert et al., 2018). After 50 ka, the
434 amplitude has been reduced by only 30%. However, by 100 ka the reduction has increased to 55% such that less
435 than half the original amplitude remains. There is no precise point at which volcanic events cannot be distinguished
436 because it depends on the original amplitude and the background variability, neither of which are constant. It is a bit
437 surprising that the synchronization of EDC and Dome Fuji (Fujita et al., 2015) matched approximately twice as
438 many events in the LIG (~ 10 per 1 ka) as in the Holocene (~ 5 per 1 ka), despite diffusion acting to reduce the peak
439 amplitudes by more than half; the LIG may have had greater volcanic activity, allowing more matches.
440 Synchronizations extending beyond the current 210 ka limit may be possible, particularly with better sampling
441 resolution to limit artificial diffusion in the records.

442



443 4.5 Limitations and Future Work

444 In this work, we have focused on obtaining a single estimate and uncertainty for the effective diffusivity for sulfuric
445 acid. The inferred effective diffusivities range by about a factor of 5 (Table 2). Many of these differences are likely
446 driven by uncertainty in the data and methods; however, the effective diffusivity may also vary for different
447 conditions in the ice sheet. Our inferences are based on average properties or groups of volcanic events, and may be
448 integrating the impacts of multiple diffusive processes. While we did not find an obvious temperature dependence,
449 there is likely an interplay between higher sulfuric acid concentrations and warmer temperatures allowing more
450 premelting (Dash et al., 2006), particularly if the sulfuric acid resides primarily at grain boundaries. More liquid at
451 grain boundaries should change the effective diffusivity locally, although related changes to the grain growth rate
452 may be necessary for a noticeable impact (e.g. Rempel et al., 2001; B03). The sulfuric acid concentration thus may
453 create a nonlinearity where the effective diffusivity depends on the concentration. The lack of knowledge of the
454 portion of sulfuric acid located at grain boundaries compared to within grain lattices is a particular limitation. Where
455 the sulfuric acid is located may depend on the conditions at deposition or the post-depositional processes involved in
456 grain growth and nucleation, such that different diffusive mechanism are more important at different locations
457 within the ice, potentially changing the effective diffusivity through time.

458 The challenges in observing where sulfuric acid resides (Mulvaney et al., 1988; Fukuzawa et al., 1998) limits the
459 ability to accurately model diffusive processes. Isolation of different diffusive mechanisms may be possible with
460 lab-grown ice doped with sulfuric acid (e.g. Hammonds and Baker, 2018) if the ice is grown such that the acid is
461 located either within the grains or at the grain boundaries, and then deformed and changes in location identified.
462 Estimates of effective diffusivity can also be extended to other impurities, as B03 originally did for chloride and
463 sodium. The differences in effective diffusivity among impurities may also help to distinguish processes that affect
464 all impurities from those that are specific to an individual impurity. With multiple international projects focused on
465 interpreting ice of 1 Ma and older, which has been strained to a small fraction of its initial thickness, more work in
466 understanding the preservation and post-depositional alteration of impurities, as well as gases and water isotopes,
467 will be necessary for confident interpretation of past climatic variations.

468

469 5 Conclusion

470 The long temporal record at EDC allows the effective diffusivity of sulfate to be estimated using multiple glacial
471 cycles. By comparing the characteristics of the sulfate record exclusively in either interglacial or in glacial
472 maximum periods, the effect of differences of the climate at deposition can be minimized. We found an effective
473 diffusivity of $\sim 5 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. That value is an order of magnitude lower than the value that was previously inferred
474 at EDC using only Holocene data (B03). The ECM data from Dome Fuji for the past glacial cycle agrees well with
475 our estimate from EDC. The uncertainty is difficult to quantify, but we suggest a value of $\pm 2 \times 10^{-9} \text{ m}^2 \text{ a}^{-1}$. Our
476 estimate of the effective diffusivity is applicable for timescales of hundreds of thousands of years and for
477 temperatures colder than -10°C . In the deep ice warmer than -10°C , the variability in the sulfate record increases due
478 to anomalous sulfate peaks (Traversi et al. 2009) caused by impurity movement and marks the end of the sulfate
479 record being indicative of the climate history. In the upper $\sim 90\%$ of the ice sheet, the low effective diffusivity
480 suggests that sulfuric acid is not readily diffusing in liquid-like veins; however, in the deep ice a connected vein
481 network appears to allow the climate variations to be replaced by peaks generated after deposition. The low effective
482 diffusivity for the cold ice in the upper portion of the ice sheet suggests that sulfuric acid and other impurity records
483 might be preserved in deep, old ice if the ice temperature remains well below the pressure melting point.

484

485



486 **6 Appendix A**

487 We used the following age and depth ranges for the interglacial and glacial maximum periods.

Interglacials				
Name	Age Start (ka)	Age end (ka)	Depth Start (m)	Depth end (m)
Holocene	0	11.3	0	312
Last Interglacial (LIG)	118	132	1522	1705
Marine Isotope Stage 7	242	246	2250	2276
Marine Isotope Stage 9	327	339	2509	2558
Marine Isotope Stage 11	399	423	2676	2738
Glacial Maximums				
Name	Age Start (ka)	Age end (ka)		
Last Glacial Maximum (LGM)	18	30	438	576
Penultimate Glacial Maximum (PGM)	136	148	1737	1795
Marine Isotope Stage 8	253	263	2299	2353
Marine Isotope Stage 10	345	365	2571	2608
Marine Isotope Stage 12	435	450	2753	2764
Marine Isotope Stage 14	537	553	2875	2891
Marine Isotope Stage 16	634	677	3007	3033
Marine Isotope Stage 18	740	759	3106	3122

488

489 **7 Code Availability**

490 All code is available upon request to: tjfudge@uw.edu. The calculation of the mean scaled gradient is well described
491 in Barnes et al., 2003 and was replicated here.

492 **8 Data Availability**

493 All EDC sulfate data is available upon request to: mirko.severi@unifi.it

494 **9 Author contributions**

495 TJF designed the study and wrote the manuscript. RS, BHH, and TJF performed scaled mean gradient analysis. LV
496 and TJF performed the volcanic width analysis. MS provided the sulfate data. EDW inspired and guided the project.
497 All authors contributed to the manuscript.

498 **10 Competing Interests**

499 The authors declare no competing interests

500

501 **11 Special Issue Statement**

502 We would like this manuscript included in two special issues:

503 [Ice core science at the three poles](#)

504 [Oldest Ice: finding and interpreting climate proxies in ice older than 700 000 years](#)



505

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