

# 1 Tectonic interactions during rift linkage: Insights from analog and 2 numerical experiments

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4 *Timothy Chris Schmid<sup>1</sup>, Sascha Brune<sup>2,3</sup>, Anne Glerum<sup>2</sup>, and Guido Schreurs<sup>1</sup>*

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6 <sup>1</sup>*Institute of Geological Sciences, University of Bern*

7 <sup>2</sup>*Helmholtz Centre Potsdam – GFZ German Research Centre for Geosciences, Potsdam, Germany*

8 <sup>3</sup>*University of Potsdam, Potsdam-Golm, Germany*

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10 *Corresponding author Timothy Schmid: [timothy.schmid@geo.unibe.ch](mailto:timothy.schmid@geo.unibe.ch)*

11 *Institute of Geological Sciences, University of Bern, Baltzerstrasse 1+3, CH-3012 Bern, Switzerland*

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13 **Keywords:** *Numerical modelling, analog modelling, stress deflection, rift interaction, rift  
14 propagation*

## 16 **Abstract**

17 Continental rifts evolve by linkage and interaction of adjacent individual segments. As rift  
18 segments propagate, they can cause notable re-orientation of the local stress field so that  
19 stress orientations deviate from the regional trend. In return, this stress re-orientation can  
20 feed back on progressive deformation and may ultimately deflect propagating rift segments  
21 in an unexpected way. Here, we employ numerical and analog experiments of continental  
22 rifting to investigate the interaction between stress re-orientation and segment linkage. Both  
23 model types employ crustal-scale two-layer setups where pre-existing linear heterogeneities  
24 are introduced by mechanical weak seeds. We test various seed configurations to investigate  
25 the effect of i) two competing rift segments that propagate unilaterally, ii) linkage of two  
26 opposingly propagating rift segments, and iii) the combination of these configurations on  
27 stress re-orientation and rift linkage. Both the analog and numerical models show counter-

28 intuitive rift deflection of two sub-parallel propagating rift segments competing for linkage  
29 with an opposingly propagating segment. The deflection pattern can be explained by means  
30 of stress analysis in numerical experiments where stress re-orientation occurs locally and  
31 propagates across the model domain as rift segments propagate. Major stress re-orientations  
32 may occur locally, which means that faults and rift segment trends do not necessarily align  
33 perpendicularly to far-field extension directions. Our results show that strain localization and  
34 stress re-orientation are closely linked, mutually influence each other and may be an important  
35 factor for rift deflection among competing rift segments as observed in nature.

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## 37 **1. Introduction**

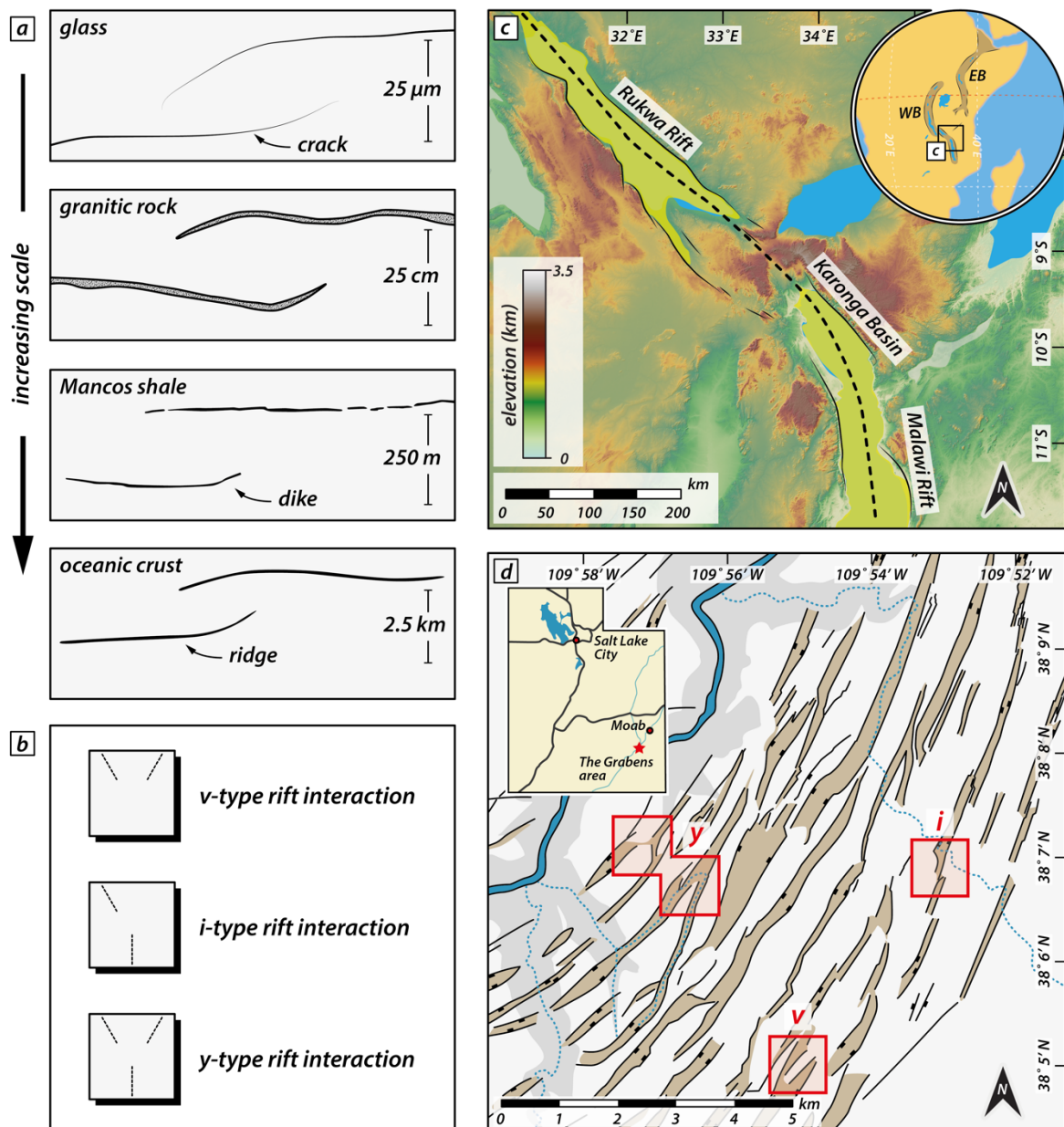
38 Continental rifting involves brittle faulting and the formation of subsiding rift basins. In places  
39 where individual rift segments are in proximity, they may interact and link when segments  
40 propagate and the rift system matures (Morley et al., 1990; Nelson et al., 1992; Rosendahl,  
41 1987). The propagation and linkage of formerly isolated rift segments resembles the  
42 propagation and interaction of extension fractures on a micro-scale (e.g., Childs et al., 1995;  
43 Willemse, 1997; Willemse et al., 1996; Fig. 1a). Indeed, analytical solutions and models have  
44 been used to describe crack growth and to predict its direction (e.g., Macdonald and Fox,  
45 1983; Mills, 1981). Such cracks occur in a variety of materials over a vast order of magnitude  
46 in length scale from micro-scale cracks in glass to km-scale ridge interaction structures in  
47 oceanic crust (Pollard and Aydin, 1984; Fig. 1a).

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49 Propagation and interaction of individual rift segments occur in continental rift systems at  
50 various styles and scales (Fig. 1b) and have been intensively studied over the years. The East  
51 African Rift System (EARS) constitutes a narrow rift with an Eastern and Western branch that  
52 propagate southward and northward, respectively (EARS; e.g., Ebinger et al., 2000; Morley  
53 et al., 1990; Nelson et al., 1992; Bonini et al., 2005; Bosworth, 1985; Brune et al., 2017; Corti

54 et al., 2019; Glerum et al., 2020; Heilman et al., 2019; Koehn et al., 2008; Kolawole et al.,  
 55 2018) comprising different sub-parallel deformed regions (inset Fig. 1c). On smaller scale,  
 56 interaction of segmented grabens has been studied for example in in the Canyonlands National  
 57 Park, Utah, a part of the Basin and Range wide rift (Allken et al., 2013; Trudgill, 2002; Schultz-  
 58 Ela and Walsh, 2002), where various styles of graben interaction are attributed to the  
 59 underlying strata (e.g., salt layer) or pre-existing weaknesses (Fig. 1d).

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63 **Figure 1:** Similar linkage structures occurring at a vast range of spatial scales. a) Propagation and linkage of segments at  
64 different scale from micro cracks in glass to linkage of oceanic ridge segments. Redrawn after Pollard and Aydin, (1984). b)  
65 Rift-interaction types investigated in this study. c) Rukwa Rift and Malawi rift along the Western Branch of the East African  
66 Rift System (EARS; inset). The two basins link obliquely via the Karonga Basin and form an i-type interaction zone. Rift axis  
67 redrawn after Kolawole et al., (2021). WB: Western Branch; EB: Eastern Branch of the EARS. d) Rift-related linked graben  
68 structures in the Canyonlands National Park, USA. Red rectangles mark areas with distinct interaction geometries (v-, i-, and  
69 y-geometries; see b) and text for detail). Redrawn after Allken et al., (2013).

70

71 Structural inheritance is thought to control nucleation and strain distribution along newly  
72 formed normal faults as weak fabrics can precondition and weaken a heterogenous upper  
73 crust (e.g., Collanega et al., 2018; Heilman et al., 2019; Kolawole et al., 2018; Morley, 2010;  
74 Morley, 1999; Kolawole et al., 2021; Morley et al., 2004). Pre-existing weak fabrics may appear  
75 as large shear zones (Daly et al., 1989), suture zones along adjacent basement terranes (Corti,  
76 2012; Corti et al., 2007) or upper crustal fabrics. Resulting rift structures may form as initially  
77 isolated segments that propagate along strike, interact and evolve into continuous zones of  
78 deformation with time as they link (Nelson et al., 1992). Rift segments link through previously  
79 un-rifted interaction zones resulting in a characteristic geometry that persists during later rift  
80 stages (Nelson et al., 1992).

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82 Recent strain accommodation in the Rukwa-North Malawi segment of the western branch of  
83 the EARS (Fig. 1c) shows dominant dip-slip faulting parallel to the border faults (Kolawole et  
84 al., 2018; Morley, 2010) driven by the reactivation of pre-existing basement fabrics (Heilman  
85 et al., 2019). There, the concentration of seismicity in the SE and NW of the Rukwa and  
86 Northern Malawi Rift, respectively suggest subsequent propagation and linkage of the rift  
87 segments with a flip in the boundary fault polarity near the interaction zone (Heilman et al.,  
88 2019 and references therein).

89

90 Pre-existing structures as well as fault interaction across multiple scales disturb the regionally  
91 inferred stress orientation (Morley, 2010; Oliva et al., 2022). In return, stress re-orientations  
92 within and adjacent to rift segments influence the style of progressive deformation. Ultimately,  
93 stress re-orientation may even favor pure dip-slip behavior even for extensional faults with an  
94 oblique orientation to the regional extension (e.g., Morley, 2010; Corti et al., 2013; Morley,  
95 2017; Philippon et al., 2015). This interplay between pre-existing structures and local re-  
96 orientation of the regional stress field affects how propagating rift segments interact. Under  
97 favorable conditions, it may even cause deflection of propagating rift segments (Nelson et al.,  
98 1992).

99

100 Rift propagation and segment interaction has been investigated by analog modelling studies  
101 that examined linkage of two segments across a transfer zone (e.g., Zwaan et al., 2016;  
102 Zwaan and Schreurs, 2017; Corti, 2012; Acocella et al., 1999; Bellahsen and Daniel, 2005).  
103 Bellahsen and Daniel (2005) studied the control of existing faults on new fault growth under  
104 multiphase extension. They suggested that pre-existing faults may disturb the local stress field  
105 and impede linkage of newly forming faults which also occurs in natural examples of  
106 multiphase extension (Duffy et al., 2015). Such stress deflections have been reported and  
107 studied in natural settings such as the North Malay Basin, Thailand, due to the vicinity of pre-  
108 existing faults (Tingay et al., 2006; Tingay et al., 2010). While analog experiments are an  
109 effective tool to simulate mechanical (brittle and ductile) deformation processes occurring in  
110 continental rifts in 3D, accessing information about stresses is challenging. In contrast,  
111 numerical modelling experiments provide direct access to element-wise stress tensors that  
112 can be interpreted in terms of stress regimes and orientations under extension (Brune and  
113 Autin, 2013; Duclaux et al., 2020). Despite the impact of stress distribution on faulting and  
114 rift segment interaction, only recently numerical studies made use of it to gain further insights  
115 into rift evolution and continental break-up (e.g., Glerum et al., 2020; Mondy et al., 2018).

116 However, these studies mostly focus on larger-scale deformation to evaluate stresses over the  
117 entire time span of rifting up to continental break-up.

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119 Here we use crustal-scale analog and numerical models to investigate rift propagation and  
120 strain localization in early rifting stages when rift segments interact. Both types of models  
121 document enigmatic rift segment deflection when two sub-parallel rift segments propagate  
122 approximately in the same direction and compete for linkage with an opposingly propagating  
123 segment. To understand the reason for rift segment deflection, we analyze the stress  
124 distribution in early rifting stages and its interplay with strain localization that initiates above  
125 pre-existing structures. Our experiments show that relatively simple rift segment interactions  
126 can cause locally complex stress patterns that deviate from the regional stress field. Such  
127 stress re-orientations occur in transient stages and can change over time and with progressive  
128 deformation due to subsequent changes in material strengths.

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130 **2. Analog model**

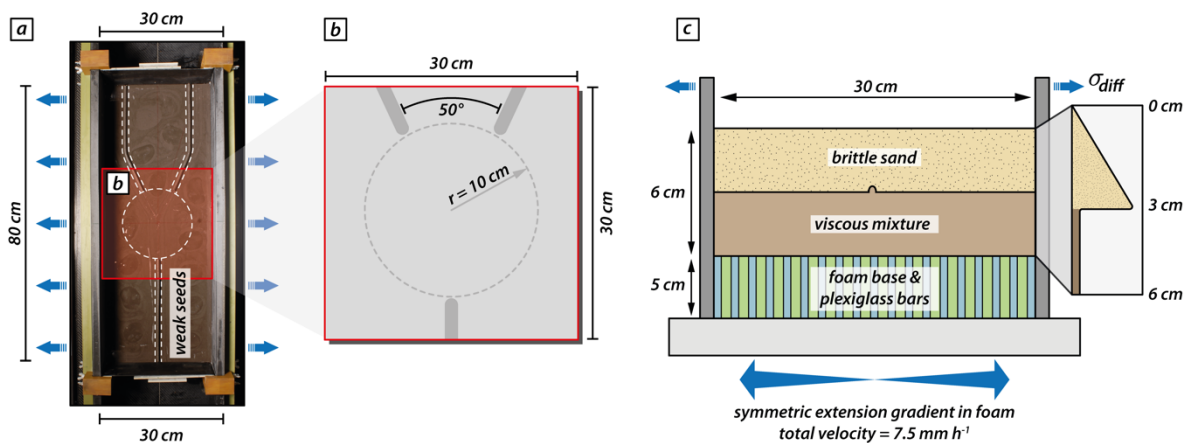
131 The presented analog modelling experiment shows unexpected features such as rift deflection.  
132 It motivates our numerical study, and we use the analog model as a reference for examining  
133 strain and stress distribution in numerical experiments.

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135 **2.1. Analog model setup**

136 For the analog reference model, we use a simplified two-layer crustal scale setup with a brittle  
137 and a viscous material to simulate upper crustal brittle faulting and lower crustal viscous  
138 deformation, respectively. The base of the model consists of a set of alternating plexiglass  
139 and foam bars which are compressed prior to the model preparation by two mobile sidewalls  
140 (Fig. 2a). During the experiment the computer-controlled sidewalls extend and provide a  
141 symmetric extension gradient as the model base expands and the model vertically thins. For  
142 monitoring the surface deformation evolution, we use a stereoscopic camera setup to take  
143 top view photos and stereo image pairs every 60 s for quantitative deformation analysis by  
144 means of 3D stereo Digital Image Correlation (Adam et al., 2005). The model was scanned  
145 every 20 min in a medical XRCT scanner for gaining insights on internal model evolution.

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149 **Figure 2:** Analog modelling setup. a) Top view of the experimental apparatus with two mobile side walls that extend  
150 orthogonally. The entire model comprises an area of 80 x 30 cm and three viscous seeds are placed on top of the viscous layer  
151 before sieving in the brittle sand layer. The central model part where propagating rift segments interact contains no seeds.

152 *b) Zoom in of the seed configuration into the analyzed model area (i.e., 30 x 30 cm). The two competing seed segments form*  
153 *an intermediate angle of 50°. The model center contains an area with a radius of 10 cm where weak seeds are absent. c)*  
154 *Sketch of the model cross section. The model setup consists of a brittle sand layer representing the upper brittle crust on top*  
155 *of a viscous mixture of PDMS and corundum sand imitating the lower ductile crust.*

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## 157 **2.2. Model geometry, rheological layering, and material properties**

158 For simulating upper crustal deformation, we use dry quartz sand with a bulk density of 1560  
159  $\text{kg m}^{-3}$  and an internal friction coefficient of 0.72 (Schmid et al., 2020a). For the lower viscous  
160 model part, we use a quasi-Newtonian PDMS/corundum sand mixture (weight ratio 1:1) with  
161 bulk density of  $1600 \text{ kg m}^{-3}$  and a viscosity of  $\sim 1 \times 10^5 \text{ Pa s}$  (Zwaan et al., 2018). Hence, the  
162 brittle-viscous setup has a density gradient that avoids density instabilities and spontaneous  
163 upwelling of the viscous layer. The model features viscous rods placed on top of the viscous  
164 model layer before sieving in the quartz sand (Fig. 2). These rods act as mechanically weak  
165 seeds and localize faulting in the upper brittle model domain. The used seed configuration  
166 includes three individual seed segments. The model includes a y-seed configuration with one  
167 seed segment perpendicular to the extension direction on one side (hereafter called frontal  
168 segment) whereas on the opposing side of the model center two obliquely placed seeds  
169 (hereafter called rear segments) form an intermediate angle of  $50^\circ$  (Fig. 2; see also Fig. 1b,d).  
170 The three seed segments hypothetically merge at the model center. However, we exclude  
171 weak seeds in an area with a radius  $r = 10 \text{ cm}$  around the model center to allow free  
172 interaction of the propagating rift structures (Fig. 2b). The analog model comprises an initial  
173 area of 80 cm by 30 cm and has a total thickness of 6 cm (each layer 3 cm) which represents  
174 a 30 km thick continental crust. In accordance with the numerical setup, the effectively  
175 analyzed model area is restricted to 30 x 30 cm. The mobile sidewalls move with an extension  
176 velocity of  $5 \text{ mm h}^{-1}$  each (totaling in  $10 \text{ mm h}^{-1}$ ), which results in a maximum extension of  
177 40 mm at the final model stage after 4h.

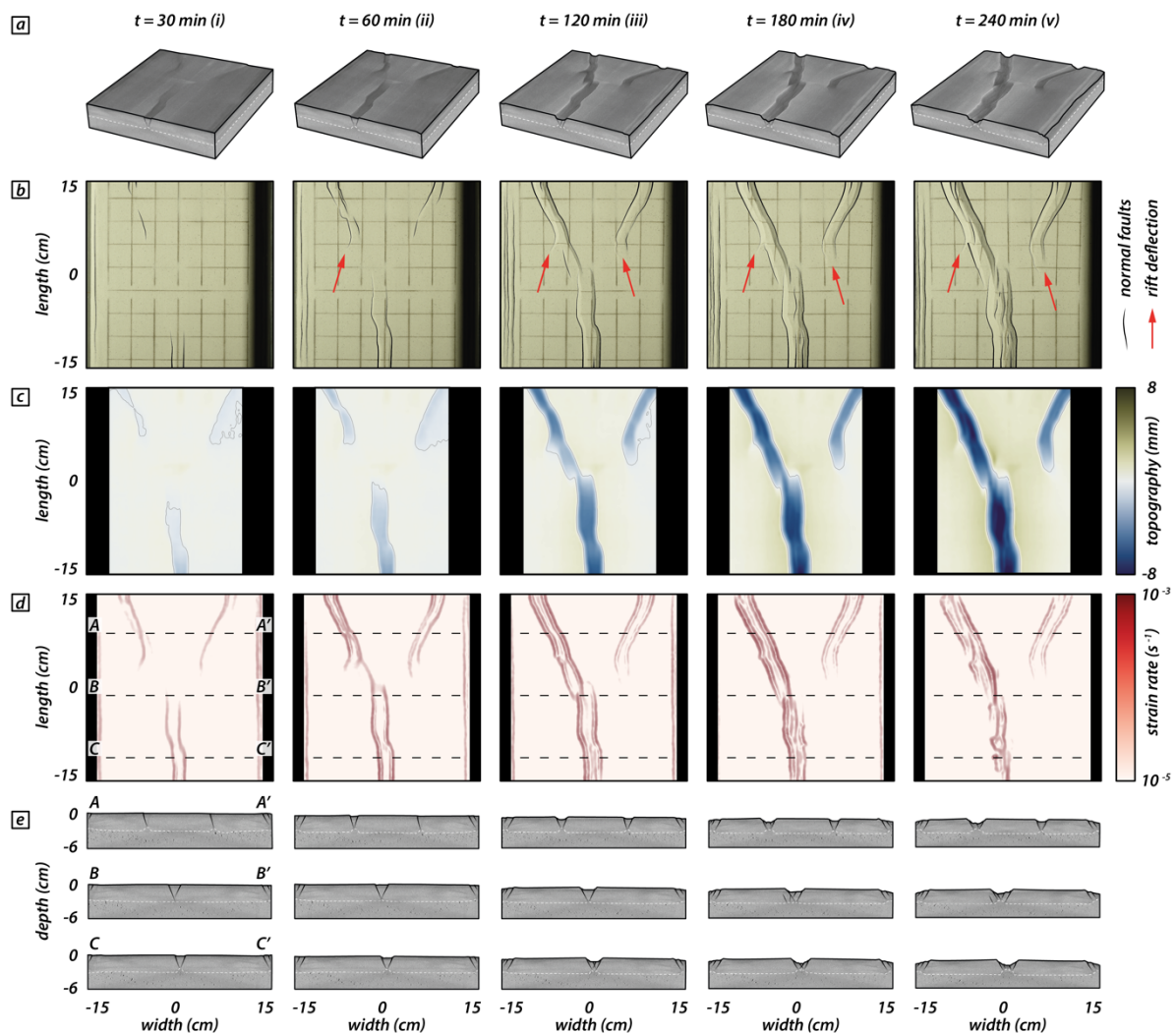
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### 179 **2.3. Analog model results**

180 In the analog model three different rift segments initiate above the weak seeds and propagate  
181 toward each other. Thereby, the two rear segments compete for linkage with the frontal  
182 segment. After 30 min (i.e., 5 mm extension; Fig. 3(i)), brittle deformation localizes along two  
183 rift boundary faults forming the frontal rift segment. Rifting in the rear segments localizes first  
184 along right-dipping rift boundary faults and after 60 min (i.e., 10 mm extension; Fig. 3(ii))  
185 both rear segments develop a set of two conjugate rift boundary faults (Fig. 3a,b (ii)).  
186 Interestingly, instead of advancing straight forward, the fault tips deflect and propagate away  
187 from each other (Fig. 3b,d (ii)). This is partially due to the rift propagation over the area where  
188 no seeds are present where rifting perpendicular to the extension direction is favored.  
189 However, after 120 min (i.e., 20 mm extension; Fig. 3 (iii)) rift tips deflect and turn away from  
190 one another. Rift tips deflect from an initially oblique orientation and rotate into an inverted  
191 oblique direction (with respect to the extension direction). The frontal and the rear left rift  
192 segment propagate further and, as they approach one another, form an en-echelon basin that  
193 convergently overlaps with the frontal rift segment (Morley et al., 1990; Fig. 3b,d (iii)). After  
194 180 min (i.e., 30 mm extension; Fig. 3(iv)), intra-rift faults develop in the frontal and left rear  
195 rift segments. Note that strain rate is successively localized in the two fully linked rift segments  
196 whereas the right rear segment experiences minor strain rate values (Fig. 3d (iv)). At the final  
197 model stage (i.e., after 240 min and 40 mm extension; Fig. 3 (v)), the right rear segment  
198 propagated minimally with a rift tip turned away from the linked segments (Fig. 3b,d (v)). The  
199 fully linked frontal and left rear segments continuously accommodated displacement resulting  
200 in deeper rift structures compared to the abandoned right rear segment (Fig. 3c,e (v)).

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**Figure 3:** Analog modelling results documenting deflection of the right rear segment and cessation of faulting activity. Distinct time steps (i.e., after 30 min and after every hour) show the model evolution. a) CT volumes of the investigated model domain at distinct time steps. White dashed lines indicate the brittle-viscous interface. b) Top views and line drawings indicating observable normal faults at the model surface. Red arrows indicate rift tips that deflect and turn away from one another. c) Topography from digital elevation models of the model surface. d) Strain rates obtained from 3D stereo DIC. Black dashed lines indicate positions of 3 transects through the CT volume. e) Rift transects A-A', B-B', and C-C'. White dashed lines indicate the brittle-viscous interface.

211

### 212 3. Numerical modelling

213 We perform a series of numerical models to investigate rift linkage interaction and to analyze  
214 occurring surface stresses. Similar to the analog experiment, the numerical model consists of  
215 a two-layer crustal setup with laterally homogenous material layers where boundary-  
216 orthogonal extension with constant velocity is applied.

### 217 **3.1. Numerical model setup**

218 We use the open source, finite-element code ASPECT to solve the extended Boussinesq  
219 equations of momentum, mass, and energy in combination with advection equations for each  
220 compositional field (Gassmüller et al., 2018; Glerum et al., 2018; Heister et al., 2017;  
221 Kronbichler et al., 2012; Rose et al., 2017; Glerum et al., 2020). Since the numerical models  
222 are motivated by the analog model, the two setups are designed in a similar way. To this aim,  
223 we employ a numerical setup where the rheologies of upper and lower crust are brittle and  
224 ductile, respectively, and independent of temperature just like in the analog model. However,  
225 the numerical models operate on the true scales of the continental crust over tens of  
226 kilometers and millions of years, while the analog model is a scaled, cm-sized representation  
227 that evolves on hour-scale. Additionally, the numerical setup applies maximum extension  
228 velocities at the side walls and extension velocities at the base that linearly increase from the  
229 center towards the model boundaries. In contrast, maximum extension velocities at the side  
230 walls in the analog model are achieved via compression of a basal foam plexiglass setup (prior  
231 to the model run) that extends homogeneously during the model run.

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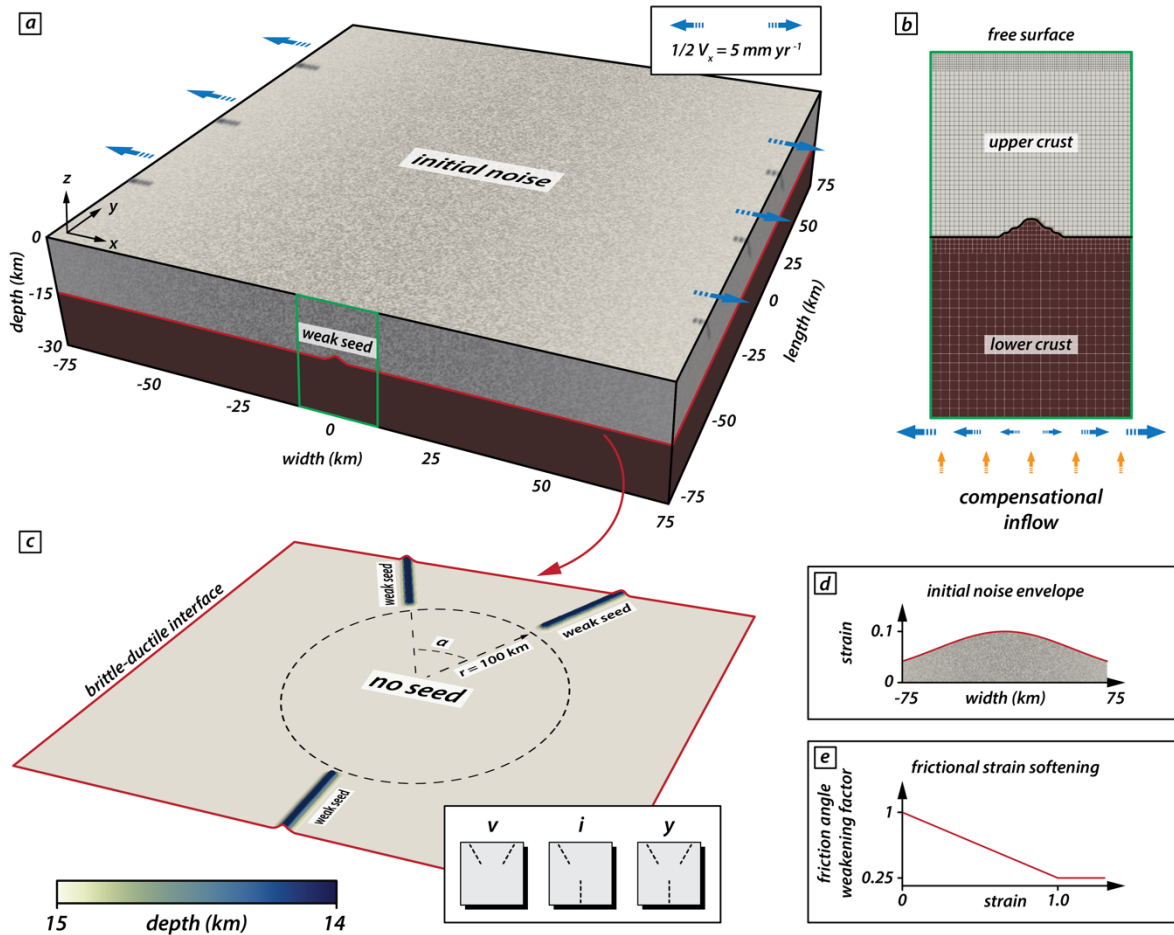
233 The presented numerical experiments cover a rectangular cuboid domain of 150 km width  
234 and length in the horizontal x- and y-direction, respectively, and 30 km in depth along the  
235 vertical z-axis (Fig. 4a). The entire model domain is divided into 1.53 million hexahedral,  
236 second-order elements. For the upper 15 km of the model, we use a cell resolution of 750 m,  
237 with an additional refinement at the uppermost km which yields near-surface elements with a  
238 resolution of 375 m. The grid resolution for the lower 15 km of the model is 1500 m. At the  
239 left and right model sides, we apply a symmetrically distributed outflow velocity of  $\frac{1}{2} V_x = 5$   
240  $\text{mm yr}^{-1}$ , resulting in a total extension velocity of  $10 \text{ mm yr}^{-1}$  (Fig. 4a,b). After a total model  
241 time of 4 My, the model has therefore experienced a total extension of 40 km. While  $V_x$  is  
242 prescribed at the left and right model sides,  $V_y$  and  $V_z$  are left free. We compensate material

243 loss through the side boundaries by compensational inflow at the model base and the  
244 horizontal  $V_x$  component increases linearly from the model center towards the lateral model  
245 boundaries (Fig. 4b). The front and back lateral boundaries allow for free slip and the top of  
246 the model features a free surface boundary condition (Rose et al., 2017).

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248 The model includes two rheological layers represented by compositional fields, namely a 15  
249 km thick visco-plastic upper crust with a density of  $2700 \text{ kg m}^{-3}$  and a 15 km thick iso-viscous  
250 lower crust with a density of  $2900 \text{ kg m}^{-3}$  and a constant viscosity of  $1 \cdot 10^{20} \text{ Pa s}$ . For the upper  
251 crust, the viscous viscosity is fixed to  $2 \cdot 10^{28} \text{ Pa s}$ , such that plastic deformation is always  
252 enabled. We introduce initial and dynamic mechanical weaknesses in the upper crust in two  
253 ways. (i) Mechanically weak seeds: At distinct positions near the brittle-ductile interface, the  
254 upper model layer is locally 10% thinned and the lower model layer elevates like the viscous  
255 weak seeds in the analog model setup. These mechanical seeds weaken the upper crustal  
256 strength and localize brittle faulting. Our experiments include three different seed  
257 configurations: v, i, and y (Fig. 4c; see also Fig. 1b-d), where seeds within a central model  
258 area (i.e.,  $r = 100 \text{ km}$ ) are absent. For each configuration, the rear seeds form an intermediate  
259 angle of  $10^\circ$ ,  $30^\circ$ , or  $50^\circ$ . (ii) Friction softening: For each element, an initial plastic strain value  
260 of 0 (resulting in strong material) to 0.1 (weaker) is randomly assigned and reduces the  
261 maximum friction angle of  $26.56^\circ$  by a maximum of 10%. This reflects the structural  
262 heterogeneity of natural settings and allows for more randomized strain patterns in the central  
263 model domain where the mechanical seeds are absent. The initial plastic strain noise is  
264 distributed over the entire model width with an amplitude following a Gaussian curve parallel  
265 to the extension direction that is repeated along the model length (y-direction, Fig. 4d). During  
266 continuous extension, the effective friction angle linearly reduces to 25% of the maximum  
267 friction angle (i.e., to  $6.64^\circ$ ) for plastic strain between 0 and 1 while it remains constant at  
268  $6.64^\circ$  for plastic strains  $> 1$  (Fig. 4e). This corresponds to a reduction of the effective friction

269 coefficient from 0.5 to 0.12. The cohesion of the upper crust remains constant at  $5 \cdot 10^6$  Pa for  
 270 all conducted experiments.  
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 274 **Figure 4:** Numerical model setup for iso-viscous models. a) The model domain comprises a volume of 150 x 150 x 30 km. Blue  
 275 arrows indicate the applied boundary-orthogonal extension. The green rectangle indicates the position of the zoom-in in b).  
 276 The red line indicates the initial depth of the brittle-ductile interface (as defined by the interface between the two rheological  
 277 layers) indicated in c). b) Initial conditions and mesh refinement (arrows not to scale). c) Position and configuration of the  
 278 mechanical weak seeds at the brittle-ductile interface. The setup comprises an area with radius  $r = 100$  km where no weak  
 279 seeds are present. Three different seed configurations refer to  $y$ -,  $i$ -, and  $v$ -models (see text for details). d) Initial amplitude of  
 280 strain along the  $x$ -axis. The Gaussian distribution is constant along the  $y$ -axis; also see grey shade in a). Note that while the  
 281 strain amplitude follows a Gaussian distribution, the location of the initial strain is random. e) Linear weakening with strain  
 282 applied to the friction angle.

283

### 284 **3.2. Model limitations**

285 Just like the analog model (Sec. 2), our crustal scale two-layer numerical setup does not  
286 comprise a lithospheric mantle layer and no asthenosphere. Further, the iso-viscous setup  
287 does not account for a temperature-dependent viscosity. However, we focus on an early rifting  
288 phase where the influence of the deforming mantle lithosphere can be neglected. The crustal-  
289 scale setup strongly limits the computational effort for calculating deformation in 3D (Allken  
290 et al., 2011, 2012; Katzman et al., 1995; Zwaan et al., 2016) and hence, our simplifications  
291 allow for a higher model resolution; a necessity to depict early stages of rifting and the  
292 coalescence of brittle deformation. Several alternative model runs have been performed  
293 including a temperature- and pressure-dependent viscosity. Those tests reproduced first-order  
294 features (i.e., strain rates, rift geometry and stress distribution) of the presented models in  
295 this study, which further justified the choice of a simplified iso-viscous setup. Note that we  
296 apply frictional softening as a function of strain within each cell. For simplicity, we do not  
297 include normalization accounting for cell size (Lavie et al., 2000) nor viscoplastic  
298 regularization techniques (Duretz et al., 2019; Jacquy and Cacace, 2020). Moreover, our  
299 model does not include the influence of melting or magma intrusions nor sedimentation and  
300 erosion.

301

### 302 **3.3. Post-processing**

303 Numerical models pose the advantage that they grant direct access to stress tensors of each  
304 individual cell. We exploit this opportunity by investigating surface stresses to deduct the  
305 stress regime and the effect of different seed configurations on stress distribution. ASPECT  
306 provides post processors that calculate the magnitude and orientation of the maximum  
307 horizontal stresses and the Regime Stress Ratio (RSR) (Glerum et al., 2020). This stress  
308 regime characterization is calculated according to the scheme of the World Stress Map  
309 (Zoback, 1992). The RSR value maps possible stress regimes to an interval between 0 and 3.

310 For isotropic and homogenous materials, the standard rules of Andersonian faulting are  
311 applied (Anderson, 1905). For RSR values  $< 1$ , faulting occurs in an extensional stress regime  
312 whereas for RSR values  $> 2$  compressive stress regimes generate thrust faults. Strike-slip  
313 faults occur for values  $1 \geq \text{RSR} \leq 2$ . We extract data of maximum horizontal compressive  
314 stress together with the stress regime and investigate them in areas where the strain rate  
315 exceeds a threshold of  $10^{-16} \text{ s}^{-1}$  and deformation occurs. For visualization, surface stresses  
316 from an originally unstructured grid are resampled on an equidistant grid.

317

### 318 **3.4. General model evolution of the reference model**

319 In this section we describe the numerical modelling results focusing particularly on the general  
320 evolution of our reference model with a y-seed configuration and an intermediate seed angle  
321 of  $50^\circ$  (Figure 5). At the early stage (i.e., after 0.5 million years), three distinct rift segments  
322 develop above the initial seed positions bounded by a pair of conjugate rift boundary faults  
323 (Fig. 5a (i)). This early stage is characterized by a symmetric evolution of the two competing  
324 rear segments, which results in a symmetric subsidence inside of the graben structures (Fig.  
325 5b (i)). For each rift segment, faulting activity is localized along the rift boundary faults. In  
326 the central model domain, however, strain rates depict a more distributed deformation pattern  
327 with multiple minor faults (Fig. 5c (i)). Note that the two rear segments propagate and show  
328 curved fault segments that initially deflect and turn away from each other resulting in rift  
329 segments with a curved geometry expressed in the topography (Fig. 5b (i)), similar to the rift  
330 evolution in the analog model. Once they overlap with the propagating frontal segment, faults  
331 symmetrically curve inwards and towards the frontal segment. The change from localized  
332 strain rates above the seeds to distributed strain rate patterns in the central model domain is  
333 best seen in transects (Fig. 5d (i)).

334

335 After the first million years, deformation has prominently localized along the left of the two  
336 rear segments and along the frontal segment (Fig. 5a,c, (ii)). While deformation in the frontal  
337 segment is localized along the rift boundary faults, inward migration occurred in the left rear  
338 segment with developing intra-rift faults and only the left-dipping rift boundary fault active.  
339 Similarly, the right rear segment shows faulting along the right-dipping rift boundary fault but  
340 activity along intra-rift faults is lacking. In the central model domain, formerly distributed  
341 deformation localized between the frontal and left rear rift segment (Fig. 5d (ii)). While strain  
342 rates indicate a shift from a symmetric to an asymmetric deformation phase, topography is  
343 still symmetric which implies that the shift is imminent and has not affected the topography  
344 after the first million years (Fig. 5b (ii)).

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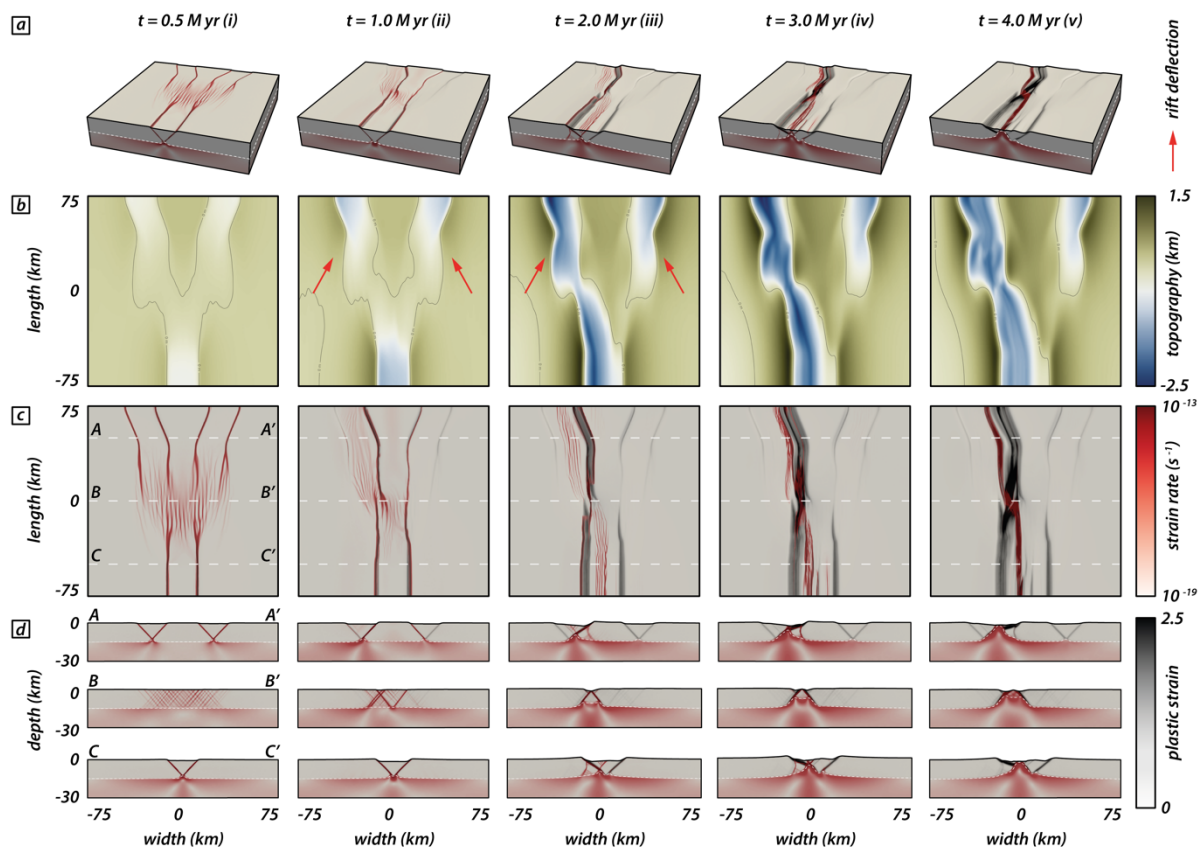
346 After two million years, deformation is entirely localized along the frontal and left rear  
347 segment. Only the right-dipping rift boundary fault of the frontal segment is active and inward  
348 migration led to a set of pervasive intra-rift faults (Fig. 5a,c (iii)). The left rear segment depicts  
349 a similar deformation pattern as in the previous step, but strain mainly accumulates along the  
350 left-dipping rift boundary fault causing an asymmetric graben geometry (Fig. 5d (iii)). Note  
351 that, after two million years, fault activity along the right rear segment completely ceased with  
352 no further strain accumulation visible (Fig. 5a,c,d (iii)). The topography reflects this completed  
353 switch from a symmetric to an asymmetric deformation stage with enhanced subsidence along  
354 the frontal and left rear segments and their linkage throughout the central model domain (Fig.  
355 5b (iii)).

356

357 With ongoing extension, deformation subsequently localizes along the axial rift zone that links  
358 the frontal and left rear segments (Fig. 5a,c,d (iv,v)) and faulting activity along rift boundary  
359 faults ceases. The linked structure reaches maximum depth inside of the rift after three million  
360 years. After four million years, however, the basin experiences minor uplift due to increase



361 upward motion of the underlying viscous material (Fig. 5d (iv,v)). Note that the basin depth  
 362 of the right rear rift segment remains stable after two million years and does not experience  
 363 further subsidence nor uplift.  
 364



365  
 366 **Figure 5:** Modelling results of the reference model documenting cessation of fault activity along the right rear segment while  
 367 the left rear and frontal segments link. Distinct time steps show the model evolution. a) Model box showing logarithmic strain  
 368 rates (red) and plastic strain (black) in the brittle and viscous model domain. White dashed lines indicate the brittle-viscous  
 369 interface. b) Top views showing the model topography. Red arrows indicate rift tips that deflect and turn away from one  
 370 another. Black lines refer to the zero-elevation height. c) Top views of the model showing strain rates (red) and corresponding  
 371 plastic strain (black) at distinct model run times. White dashed lines correspond to the three rift transects A-A', B-B', and C-C'  
 372 in subfigure d). d) Rift-axis perpendicular transects A-A', B-B', and C-C' parallel to the extension direction.

374

### 375 **3.5. Early localization patterns for v-, i-, and y-seeds**

376 To investigate the influence of different seed configurations, we compare v- (Fig. 6a-c), i-  
 377 (Fig. 6d-f), and y-seed (Fig. 6g-i) configurations for different intermediate angles (i.e., 10°,  
 378 30°, and 50°) at an early stage after 0.5 million years. y- and i-seed configurations provide a

379 setup where rift structures opposingly propagate towards the model center where rift linkage  
380 eventually occurs. In contrast, rift structures in the v-seed configuration propagate  
381 approximately in the same direction, which has a consequence on the overall strain rate  
382 distribution.

383

384 The early stage in v-seed experiments (Fig. 6a-c) is characterized by a zone of localized and  
385 distributed deformation in the rear and frontal part of the experiments, respectively. The  
386 transition from localized to distributed deformation occurs where the two competing rift  
387 segments deflect and rotate away from one another. Note that the fault deflection successively  
388 decreases towards the left and right model sides, where faults strike perpendicular to the  
389 extension direction. This is consistent with observations for experiments with a y-seed  
390 configuration. However, there the two competing rear segments rotate back and eventually  
391 bend towards the propagating frontal segment (Fig. 6g-i).

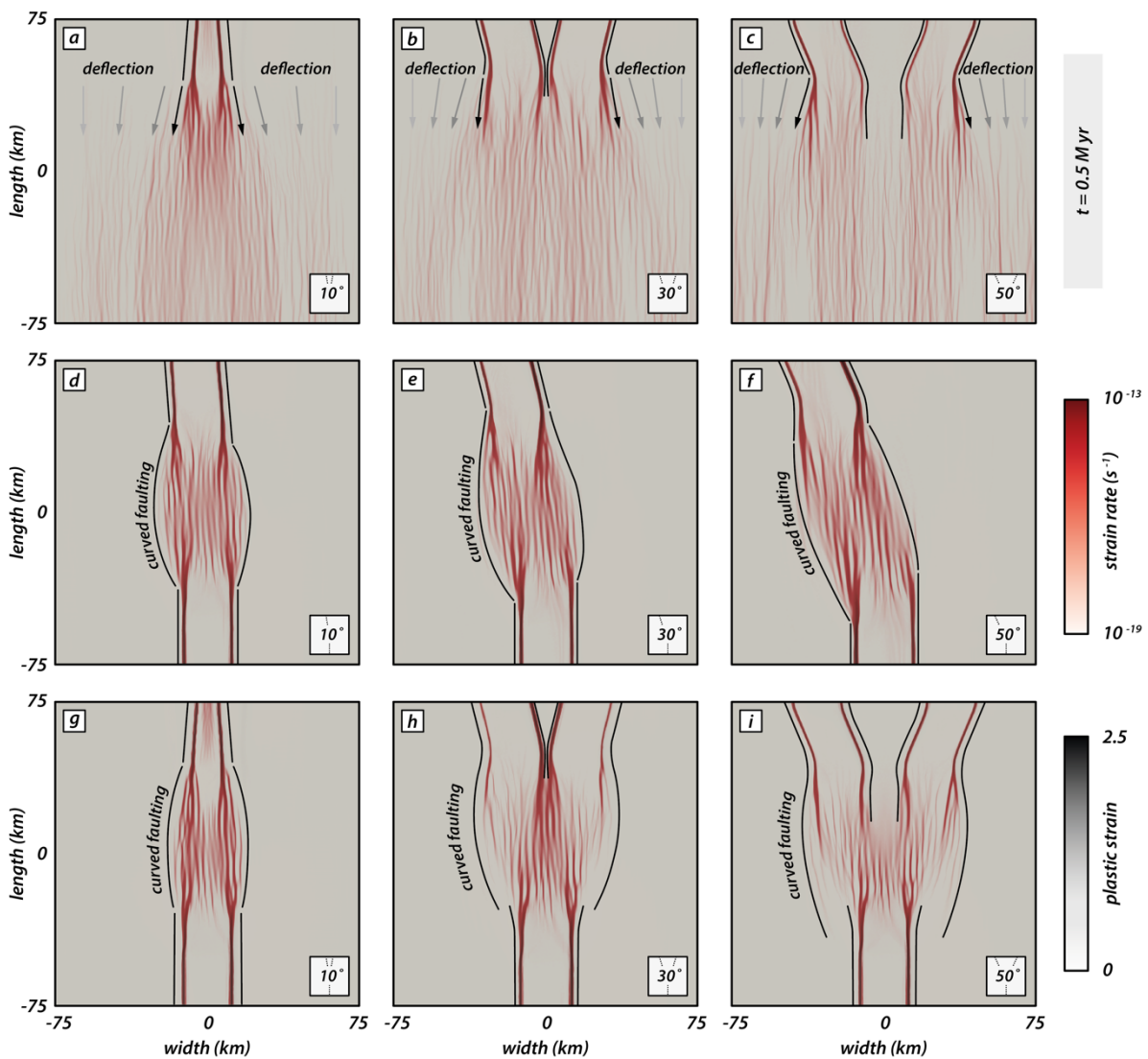
392

393 For experiment with a i-seed configuration (Fig. 6d-f) two opposingly propagating rift branches  
394 form. Since the right rear segment is absent, both opposingly propagating rift segments link  
395 in the model center where deformation is distributed onto intra-rift faults. The overall strain  
396 rate field is localized, and no strain rate deflection occurs.

397

398 Models with a y-seed configuration (Fig. 6g-i) depict a strain rate pattern where deformation  
399 is localized along rift boundary faults at the model margins where seeds are present and a  
400 distributed en-echelon strain rate pattern in the model center. Note that for the model with  
401 an intermediate angle of  $10^\circ$  the two competing rear segments are close enough resulting in  
402 a zone where strain is localized along only one rift boundary fault per rift segment (i.e.,  
403 outward-dipping faults with respect to the model box) that overlap and form a central graben  
404 with minor intra-rift faults. For larger intermediate angles, two individual rift segments

405 (bounded by two rift boundary faults) form that propagate towards the model center. While  
 406 the strain rate pattern due to the competing rear segments is identical for experiments with  
 407 a y- and v-seed configuration, the additional frontal segment in experiments with a y-seed  
 408 configuration causes localization of strain rates in a single rift branch bounded by two rift  
 409 boundary faults. This contrasts with the v-seed configuration where strain rates in the frontal  
 410 model domain occur distributed over the entire model domain (Fig. 6a-c).  
 411



412  
 413

414 **Figure 6:** Types of rift segment linkages depending on the seed configuration at an early phase after 0.5 million years. Model  
 415 top views show strain rates (logarithmic) and plastic strain in red and black colors, respectively. a-c) v-seed configuration for  
 416 intermediate angles of 10°, 30°, and 50°. d-f) i-seed configuration for intermediate angles of 10°, 30°, and 50°. g-i) y-seed  
 417 configuration for intermediate angles of 10°, 30°, and 50° (reference model). Black lines confine deformed areas. For models

418 *with a v-seed configuration (a-c), competing rift segments deflect away from each other resulting in a fan-shaped geometry.*  
419 *Note that fault strike successively re-orientates into an orientation perpendicular to the extension direction towards the left and*  
420 *right model sides. Curved faulting occurs in models with an i- and y-seed configuration (d-j) where rift segments interact.*

421

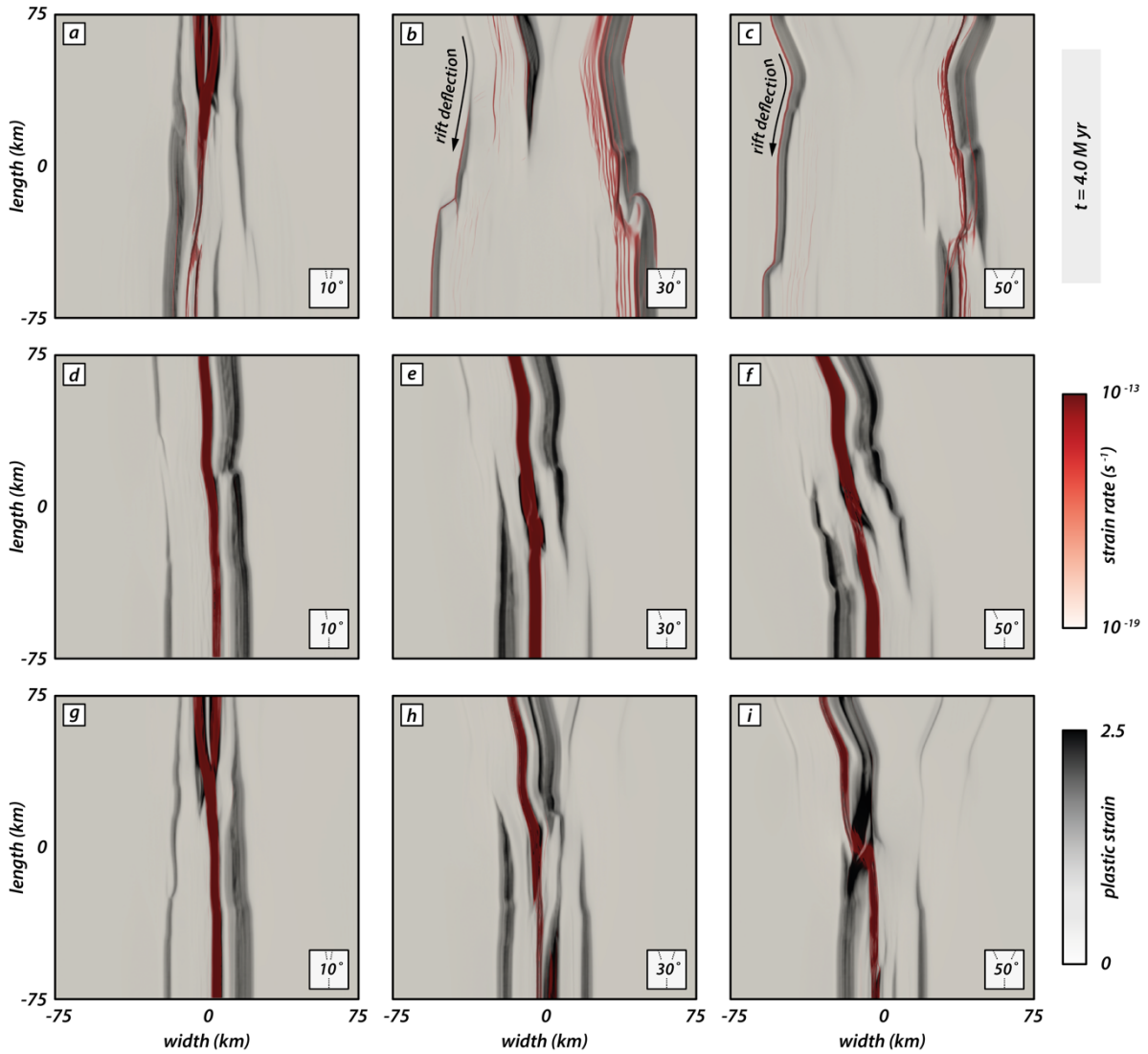
### 422 **3.6. Final rift geometry and localization patterns for v-, i-, and y-seeds**

423 The final model stage after four million years best illustrates differences in rift geometry  
424 between the models with different seed geometry and an intermediate angle (Fig. 7). Rift  
425 deflection is well visible in v-seed models (Fig. 7 a-c) and most prominent in experiments with  
426 a larger intermediate angle (Fig. 7b,c). Above the seeds, two short individual rift segments  
427 form bounded by a pair of conjugate rift boundary faults. However, as the rifts propagate  
428 towards the model center, strain is mainly accommodated along the boundary faults that dip  
429 towards the model center. Hence, the larger part of the model subsides uniformly and builds  
430 a broad rift zone confined by two large boundary faults. When the two rift segments  
431 propagate, they deflect and turn away from one another resulting in a gradually wider rift.  
432 For intermediate angles of 30° and 50°, both competing rift segments show active faulting  
433 along intra-rift faults in the rear model part, but a zone of continuous faulting activity has  
434 developed along the right side of the rift.

435

436 Models with an i-seed configuration show a continuous and straight rift geometry for all  
437 intermediate angles (fig. 7d-f). For an intermediate angle of 10°, the rift structure is nearly  
438 orthogonal with respect to the extension direction. Note that most plastic strain is  
439 accommodated along the left-dipping rift boundary fault (Fig. 7d). For larger intermediate  
440 angles, the rift subsequently experiences more segmentation with small left stepping  
441 segments towards the rear model part (Fig. 7e,f). Strain accommodation occurs mainly on the  
442 right-dipping rift boundary fault for the frontal model part and switches to the left-dipping  
443 boundary fault in the rear model part.

444



445

446 **Figure 7:** Influence of seed configuration on the final rift geometry after 4 million years. Strain rates (logarithmic) and plastic  
 447 strain are indicated by red and black colors, respectively. a-c) v-seed configuration for intermediate angles of 10°, 30°, and  
 448 50°. d-f) i-seed configuration for intermediate angles of 10°, 30°, and 50°. g-i) y-seed configuration for intermediate angles  
 449 of 10°, 30°, and 50° (reference model).

450

451 The most prominent difference occurs in models with a y-seed configuration and various  
 452 intermediate angles. For an intermediate angle of 10°, the final rift geometry resembles that  
 453 of a continuous straight rift segment (Fig. 7g). Both competing rear seeds are close enough  
 454 such that they build one rift system rather than two distinct branches. For y-seed models with  
 455 a larger intermediate angle (Fig. 7h,i), two individual rear rift segments form and compete for  
 456 linkage with the frontal rift segment. Plastic strain well illustrates the asymmetric strain

457 accommodation focused along the left-dipping rift boundary fault of the left rear segment,  
458 whereas the right rear segment only experienced minor strain accommodation (Fig. 7h,i). In  
459 both cases, high strain rates are localized in the axial rift zone and witness activity along the  
460 linked frontal and left rear segments.

461

462 Note that all experiments with an intermediate angle of  $10^\circ$  (Fig. 7a,d,g) form continuous  
463 straight rift segments, regardless of the seed configuration. Additionally, the final rift geometry  
464 for y- and v-seed configurations for an intermediate angle of  $10^\circ$  is similar with a gently wider  
465 rift in the frontal model part (Fig. 7a,g). In contrast, for i-seed configurations the rift width is  
466 similar along the entire length with a minor lateral offset (Fig. 7d). Strain rates are localized  
467 in the axial rift zone throughout the entire model length forking into two close zones in the  
468 rear end where the competing seeds are located.

469

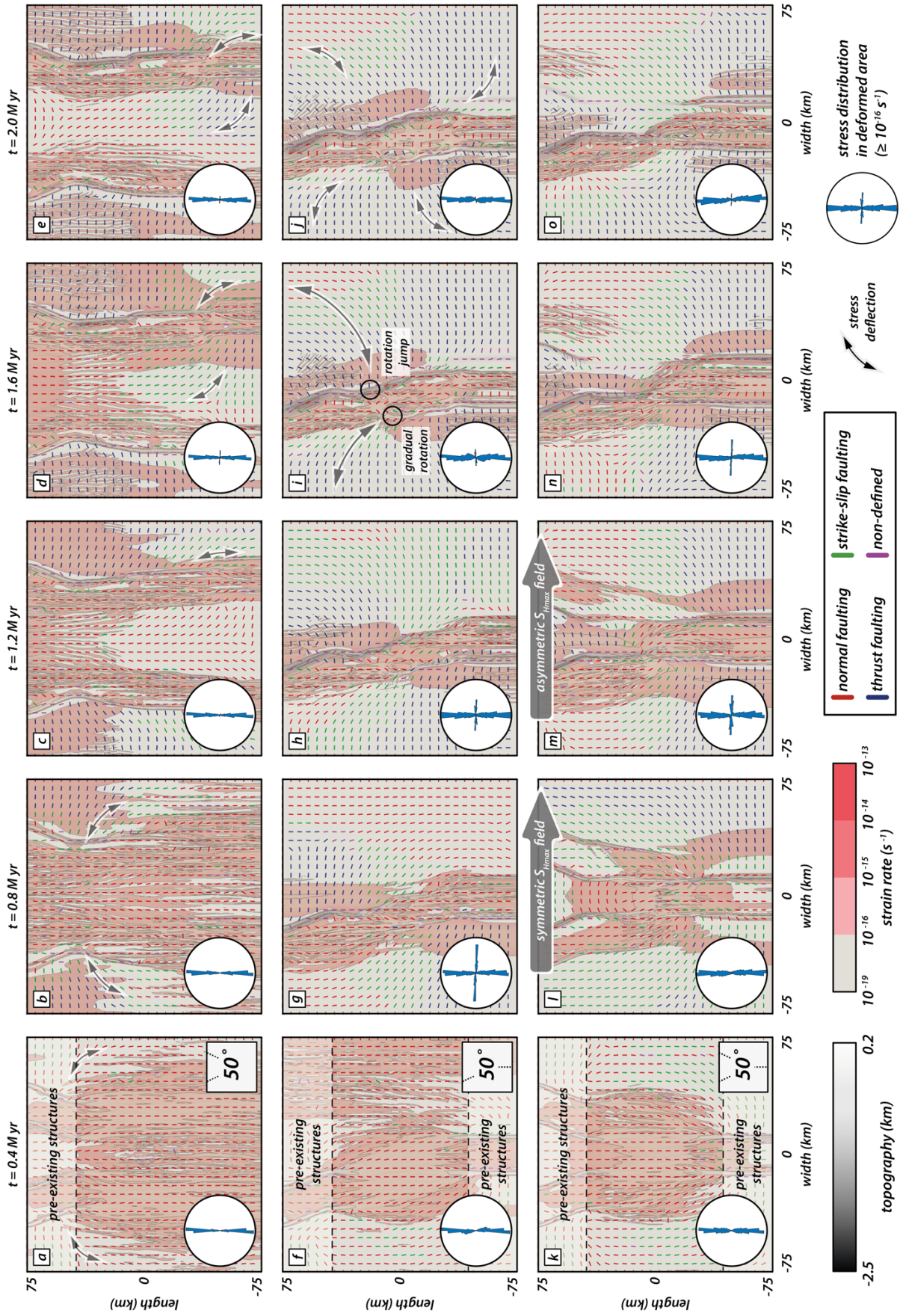
### 470 **3.7. $S_{Hmax}$ evolution with progressive deformation**

471 In this section we present the distribution and orientation of the maximal horizontal  
472 compressive stress component  $S_{Hmax}$  with progressive rift evolution and segment linkage. We  
473 focus on models with v-, i-, and y-seed configurations and an intermediate angle of  $50^\circ$  (Fig.  
474 8; see also supplementary Figures S1-S3) distinguishing between model zones with pre-  
475 existing structures (i.e., weak seeds) and a central zone where material strength is isotropic.

476

477 Our models depict two distinct phases within the first two million years: early strain  
478 accommodation over a wider model domain followed by strain localization and linkage of  
479 propagating rift segments (see also supplementary Figures S4-S6). Consequently, we focus  
480 on  $S_{Hmax}$  in the first two million years of deformation and its effect on rift propagation. Figure  
481 8 shows the orientation of  $S_{Hmax}$  and the stress regime based on the common color scheme of  
482 the World Stress Map (Heidbach et al., 2018). Note that  $S_{Hmax}$  orientation and the stress regime

483 alone do not suffice to discriminate between locations where stresses exceed crustal strength  
484 and faulting occurs. Strain rate values provide further necessary information, and we use a  
485 threshold of  $10^{-16} \text{ s}^{-1}$  that splits the model into locations of active deformation (i.e.,  $\geq 10^{-16} \text{ s}^{-1}$ )  
486 and tectonically inactive domains (i.e.,  $< 10^{-16} \text{ s}^{-1}$ ).  
487



488

489



490 **Figure 8:** Interplay of rift localization and surface stresses. Top views show the distribution of the maximum horizontal  
491 compressive stress component  $S_{Hmax}$  (not scaled to the magnitude) in models with an intermediate angle of  $50^\circ$  at early  
492 deformation stages (i.e., until 2 million years). a-e) v-seed configuration. f-j) i-seed configuration. k-o) y-seed configuration.  
493 Black colors refer to topographic elevation and red colors mark zones where strain rates exceed a threshold of  $10^{-16} s^{-1}$ . Color  
494 coding for the stress regime marks normal, strike-slip, and thrust faulting in red, green, and blue, respectively, using the  
495 common color scheme of the World Stress Map (Heidbach et al., 2018). Elements where the stress regime is non-defined are  
496 marked purple. Black arrows highlight stress deflection of  $S_{max}$ . Rose diagrams show the distribution of  $S_{Hmax}$  orientation in  
497 zones where active faulting occurs (i.e., strain rate  $\geq 10^{-16} s^{-1}$ ). Large grey arrows for the y-seed configuration mark the change  
498 from a symmetric to an asymmetric  $S_{Hmax}$  distribution.

499

### 500 **3.7.1. Effect of $S_{Hmax}$ re-orientation on rift propagation of competing rift** 501 **segments (v-seed models)**

502 Early stages in our numerical experiments are characterized by rift deflection and curved fault  
503 traces in the model center where rift segments interact (see Fig. 6). Hereafter we refer to that  
504 phenomenon as arcuate faulting. Arcuate faulting mainly occurs in experiments with larger  
505 intermediate angles ( $>10^\circ$ ) in early stages (Fig. 6), especially if two competing rift segments  
506 are present (v-, and y-seed configurations). Moreover, we have shown that deflection of  
507 propagating rifts occurs when deformation is symmetrically distributed along both competing  
508 rift branches. This is clearly visible for the v-seed configuration (Fig. 8a-e). Assuming  
509 orthogonal extension and isotropic material properties,  $S_{Hmax}$  is expected to align perpendicular  
510 to the extension direction producing pure dip-slip normal faults (Anderson, 1905). However,  
511 the model shows an immediate  $S_{Hmax}$  re-orientation at early deformation stages (i.e., after 0.4  
512 million years; Fig. 8a) from a N-S to a E-W orientation in the vicinity of the underlying weak  
513 seeds such that dip slip faults are favored over oblique-slip faults with a strike-slip component.  
514 With progressive extension (Fig. 8b-e),  $S_{Hmax}$  re-orientations successively propagate into the  
515 isotropic zone without pre-existing structures, concomitant with the rift propagation.  
516 Consequently, the position of the front where stress rotation occurs propagates over time  
517 resulting in the deflection of the propagating rift arms away from each other.

518

519 There is a distinct difference between stress deflection along weak structures and E-W  
520 deflections of  $S_{Hmax}$  in zones where strain rates are below the set threshold of  $10^{-16} s^{-1}$ . The v-  
521 seed configuration shows localized strain accumulation along one rift boundary fault per  
522 segment (i.e., the outer one) resulting in a rift zone with a broad graben system that subsides  
523 (Fig. 8e).  $S_{Hmax}$  re-orientation inside of the graben is in parts identical to the E-W orientation  
524 of  $S_{Hmax}$  outside of the graben. While local  $S_{Hmax}$  rotations may be explained by small  
525 differences in the maximum and intermediate principal stress components, such E-W stress  
526 re-orientation in our model occurs systematically and suggest that this feature reflects the  
527 influence of the strength anisotropy (Morley, 2010). The initial  $S_{Hmax}$  deflection near weak  
528 structures locally favors dip-slip faulting but also has regional influence on the overall stress  
529 regime.

530

### 531 **3.7.2. $S_{Hmax}$ evolution in sub-parallel rift segments (i-seed models)**

532 During the early stage (i.e., after 0.4 million years, Fig. 8f), the distribution of  $S_{Hmax}$  resembles  
533 the distribution from the v-seed configuration described in the previous section. Stress  
534 deflection mainly occurs in zones where a weak fabric is present.  $S_{Hmax}$  values in the central  
535 zone rotate by a small amount and reflect arcuate faulting (see Fig. 6). Since the two rift  
536 segments propagate in opposing directions, linkage is efficient and localizes in a short time  
537 (Fig. 8f-j).  $S_{Hmax}$  values deflect accordingly along propagating faults, which affects the entire  
538 model domain. This deflection does not occur symmetrically on both sides of each rift segment.  
539 Rather, it shows two distinct zones: 1) E-W orientations outside the rift deflect into a parallel  
540 orientation near the rift boarder or 2) N-S orientations outside of the rift deflect into E-W  
541 orientations near faults (Fig. 8j).

542

543 We find that  $S_{Hmax}$  orientations deflect gradually from E-W to N-S along abandoned rift  
544 boundary faults where activity ceased (Fig. 8h-j; upper left and lower right model domain). In

545 contrast,  $S_{Hmax}$  re-orientations from N-S to nearly E-W towards active rift boundary faults are  
546 followed by a rapid flip back to N-S along the faults (Fig. 8h-j; lower left and upper right model  
547 domain). The two types of re-orientation seem to correspond with two types of deformed  
548 zones. Where deformation is strongly localized along rift boundary faults, jumps in the  $S_{Hmax}$   
549 orientation occur. In contrast, zones where inward migration of fault activity activates intra-  
550 rift faults,  $S_{Hmax}$  re-orientation occurs gradually.

551

### 552 **3.7.3. Rift arm competition and deflection (y-seed models)**

553 A prominent feature in our models with two competing rift segments is the deflection of rift  
554 branches and arcuate strain rate patterns (Fig. 8a-e) in the model with a v-seed configuration.  
555 Moreover, the i-seed configuration demonstrates a gradual  $S_{Hmax}$  re-orientation over a broader  
556 pre-weakened zone due to formerly active boundary faults. One could therefore expect that  
557 both features should occur in the model with y-seed configuration (Fig. 8k-o).

558

559 Indeed, early stages (i.e., after 0.4 million years; Fig. 8k) are characterized by a symmetric  
560 stress field with re-oriented  $S_{Hmax}$  values near the two rear rift segments. However, in contrast  
561 to the v-seed configuration,  $S_{Hmax}$  re-orientation also occurs near the frontal pre-existing weak  
562 fabric along developing rift boundary faults. In the isotropic zone,  $S_{Hmax}$  values dominantly  
563 show a N-S direction. The general N-S orientation reflects the regional stress field due to an  
564 E-W extension as predicted by Anderson (1905) in isotropic areas, into which rift segments  
565 have yet to propagate. With ongoing extension, all three rift segments propagate into the  
566 isotropic zone and cause a re-orientation of  $S_{Hmax}$  (Fig. 8l). Note that after 0.8 million years  
567 the stress re-orientation occurs symmetrically. This contrasts with the i-seed configuration  
568 where  $S_{Hmax}$  values deflect into either an E-W orientation along active rift boundary faults or  
569 gradually turn into a fault parallel direction over a broader weakened zone (see subsection  
570 3.7.2.). The early symmetric stress distribution in the y-seed configuration model is

571 unarguably due to the symmetric seed configuration (see also Fig. 8a-e). At this stage, dip-  
572 slip faulting along the competing sub-parallel rift segments is favored over oblique slip faults  
573 as in models with a v-seed configuration. It is only after 1.2 million years, when fault activity  
574 along the right rear segment ceases that deformation localizes along the left rear and frontal  
575 segments and linkage intensifies (Fig. 8m). Successively, localization and linkage occur  
576 coevally with a switch from a symmetric to an asymmetric stress distribution and resembles  
577 more the stress distribution in the i-seed configuration model (Fig. 8f-j). The model state after  
578 1.2 million years (Fig. 8m) also marks the switch from a symmetric to an asymmetric stress  
579 distribution that was formerly dominated by the competing rear rift segments with dip-slip  
580 faulting favored along the two competing rift segments (see also v-seed configuration; Fig. 8  
581 a-e). After 1.2 million years the system is dominated by the linkage of two obliquely oriented  
582 segments (i.e., i-seed configuration). Note that after 1.2 million years dip-slip faulting mostly  
583 occurs along the competing rift segment that links with the opposingly propagating segment  
584 whereas dominantly oblique slip faults occur along the abandoned rift segment where activity  
585 ceases.

586

587 The symmetry switch is also visible in rose diagrams of stress orientations within the active  
588 faulting zone (i.e., strain rate  $\geq 10^{-16} \text{ s}^{-1}$ ). A dominantly N-S oriented  $S_{Hmax}$  distribution changes  
589 to a bimodal distribution with a second E-W orientation (Fig. 8l-n). Similarly, bimodal  $S_{Hmax}$   
590 distribution is also visible in the experiment with an i-seed configuration but occurs earlier.  
591 Since the experiment with an i-seed configuration is never in the state of an early symmetric  
592 stress distribution linkage is facilitated and occurs earlier (Fig.8g-i).

593

#### 594 **4. Discussion**

595 Despite the relatively simple setup of our experiments, the interaction of individual weak seeds  
596 generates a complex evolution of linkage patterns. In the following we discuss the effect of

597 pre-existing structures on  $S_{Hmax}$  re-orientations and how, in return, stress re-orientation  
598 influences rift propagation and rift segment linkage.

599

#### 600 **4.1. Effect of pre-existing structures on rift segment propagation,** 601 **interaction, and $S_{Hmax}$**

602 Previous modelling studies demonstrated that pre-existing weaknesses may cause local re-  
603 orientations of  $S_{Hmax}$  resulting in extensional faults with an oblique orientation to the regional  
604 extension direction which exhibit pure dip-slip behavior (e.g., Morley, 2010; Corti et al., 2013;  
605 Morley, 2017; Philippon et al., 2015). This contrasts the expected (assuming Andersonian  
606 faulting theory) occurrence of faults with an oblique slip component above pre-existing  
607 structures that are obliquely oriented with respect to the extension direction (Tron and Brun,  
608 1991; Withjack and Jamison, 1986). Our  $S_{Hmax}$  analysis documents two types of stress re-  
609 orientation, either gradually or by a jump along faults (Fig. 8i). A potential explanation for the  
610 two types of stress deflection is that cessation of boundary fault activity (and subsequent  
611 faulting activity along intra-rift faults) creates a broad zone of reduced crustal strength. Hence,  
612  $S_{Hmax}$  orientations successively re-orient along those formerly active faults and eventually  
613 rotate into a N-S orientation along active intra-rift faults. In contrast, where faulting activity  
614 is strongly localized along rift boundary faults, re-orientation occurs rapidly by a jump from E-  
615 W to a N-S orientation. This suggests that formerly active faults act as a wider zone of pre-  
616 weakened material, where stresses deflect sequentially rather than with a rapid jump. Similar  
617 observations have been made in previous studies of numerical models (Gudmundsson et al.,  
618 2010; Kattenhorn et al., 2000). These experiments suggest that earlier fractures lead to  
619 subzones (within a broader damage zone), where stresses subsequently rotate away from the  
620 regional stress field. Although our analog and numerical models do not feature elastic  
621 deformation, they indicate that stress deflection is an ongoing process, even after elastic  
622 material failure. Such a stress deflection further implies that stress orientations in rocks with

623 pre-existing weaknesses can substantially deviate from predicted orientations in isotropic  
624 media (Anderson, 1905).

625

626 It has been proposed that early faulting and propagation in the Rukwa and North Malawi Rifts  
627 (Fig. 1c) were guided by pre-existing basement fabrics (Heilman et al., 2019). This region is  
628 further shaped by a flip in the boundary fault polarity in the present-day geometry within the  
629 interaction zone between Rukwa Rift and North Malawi Rift (Bosworth, 1985). Our i-seed  
630 models show identical geometries for increasing intermediate angles (Figs. 7h,i and S5), where  
631 plastic strain near pre-existing structures is mostly accommodated along prominent rift  
632 boundary faults that flip fault polarity from the frontal to the rear rift segment. This flip in  
633 fault polarity occurs prominently in models with an intermediate angle  $\geq 10^\circ$ . We speculate  
634 that the increasing obliquity of the southward propagating rift segment favors asymmetric  
635 graben evolution with one dominant boundary fault accommodating a larger amount of strain.  
636 In contrast to small intermediate angles (i.e.,  $10^\circ$ ), seed configurations with a higher obliquity  
637 provoke local rotation of  $S_{Hmax}$  within the interaction zone into a strike-slip regime near the  
638 subordinate boundary fault (Fig. S5). Hence, strain accommodation along incipient faults  
639 within the dip-slip regime is favored. This facilitates propagation along those dominant rift  
640 boundary faults and eventually defines the final rift geometry.

641

642 Kolawole et al. (2018) further propose two different types of strain accommodation at early  
643 rift phases. Prominent strain accommodation localized onto a discrete and narrow zone along  
644 large rift boundary faults (Style-1; sensu Kolawole et al., 2018) and faulting distributed over  
645 a broader zone, where fault clusters may reflect pre-conditioning of the material (Style-2;  
646 sensu Kolawole et al., 2018). With this perspective, jumps and gradual rotation of  $S_{Hmax}$   
647 orientations are comparable to Style-1 and Style-2 strain localization, respectively, as  
648 proposed by Kolawole et al. (2018). Hence, the type of weakness (narrow discrete zone or

649 distributed cluster zone) should also be reflected by the stress re-orientation distribution  
650 (Morley, 2010).

651

#### 652 **4.2. Local $S_{Hmax}$ re-orientation and its influence on rift segment interaction** 653 **and rift deflection**

654 A particular observation in our experiments with a v-, and y-seed configuration is that two  
655 sub-parallel rift segments, which propagate approximately in the same direction deflect away  
656 from each other at early stages. This is somewhat surprising as one would expect the two rift  
657 segments to cut towards each other by minimizing fault length. The occurrence of rift  
658 deflection in both analog and numerical experiments validates that the results are robust and  
659 require discussing the role of  $S_{Hmax}$  re-orientation and how it influences rift segment  
660 interaction.

661

662 We speculate that, while both rear rift segments in our y-seed models equivalently  
663 accommodate strain in the early stages (i.e., when the overall stress distribution is symmetric;  
664 Fig. 8),  $S_{Hmax}$  orientations are dominated by the influence of the two competing rear rift  
665 segments that accommodate strain in equal parts. It is only after fault activity along one rear  
666 segment ceases that deformation localizes along the active rear and frontal segments and  
667 linkage intensifies. Strain localization and linkage occur coevally with a switch from a  
668 symmetric to an asymmetric stress distribution resembling the stress distributions in v-, and  
669 i-seed configuration models, respectively. The switch from a symmetric to an asymmetric  
670 stress distribution in y-seed models also marks the switch from a system that was formerly  
671 dominated by the competing rear rift segments (i.e., v-seed configuration) to a system that  
672 is dominated by the linkage of two obliquely oriented segments (i.e., i-seed configuration).

673

674 In models with a v-seed configuration, however, the symmetric phase prevails and causes  
675 coeval  $S_{Hmax}$  re-orientations and rift deflection that cause divergence of the two propagating  
676 rift segments. A similar process of extensional segment interaction via stress rotation is known  
677 from mid-ocean ridge settings: Pollard and Aydin (1984) argue that paths of two opposingly  
678 propagating oceanic ridges weakly diverge due to shear stresses that divert propagating ridges  
679 as they approach each other. Once the two ridges overlap, the stress field changes causing  
680 convergence and intersection. Similarly, Nelson et al. (1992) describe interference of  
681 compressional zones of propagating cracks diverting their tips before they overlap and turn  
682 back toward another. In this respect, our models with a v-seed configuration suggest that  
683 stresses also cause divergence of two rift segments that propagate approximately in the same  
684 direction. However, overlap never occurs (as they propagate approximately in the same  
685 direction) and hence, the two segments remain in a stress field that further diverts their paths.  
686

687 Only in models with a y-seed configuration, compressional zones and rift deflection can be  
688 overcome once the opposingly propagating rift segment links with one of the competing rift  
689 segments. Linkage occurs after about the first million years, concurrently with rift deflection  
690 and abandonment of the right rear segment (Figs. 9a and S6). Moreover, remaining activity  
691 in the right rear segment depicts low strain rates along numerous arcuate intra-rift faults (Figs.  
692 9b and S6). This suggests that linkage and rift abandonment are closely coupled and faulting  
693 along the linked segments intensifies when the activity along the remaining rift segment  
694 ceases. In addition, the left rear segment displays a rather asymmetric half graben geometry  
695 (Figs. 5c,d, 7i and S4) with one prominent rift boundary fault accommodating a larger part of  
696 plastic strain similar to our models with a i-seed configuration (see also Figs. 7e,f and S5).  
697 Dominant strain accommodation occurs along the west-dipping rift boundary fault of the left  
698 rear segment coinciding with jumps in the  $S_{Hmax}$  orientation (Fig. 8m-o). Our modelling results  
699 show that stress deflection along rift segment tips is a mechanical consequence of the



700 interaction between weak zones and far-field stresses offering a potential explanation for  
701 naturally occurring rift deflection. However, we must emphasize that complexities in natural  
702 rift settings pose additional difficulties that require further investigations of stress orientations.

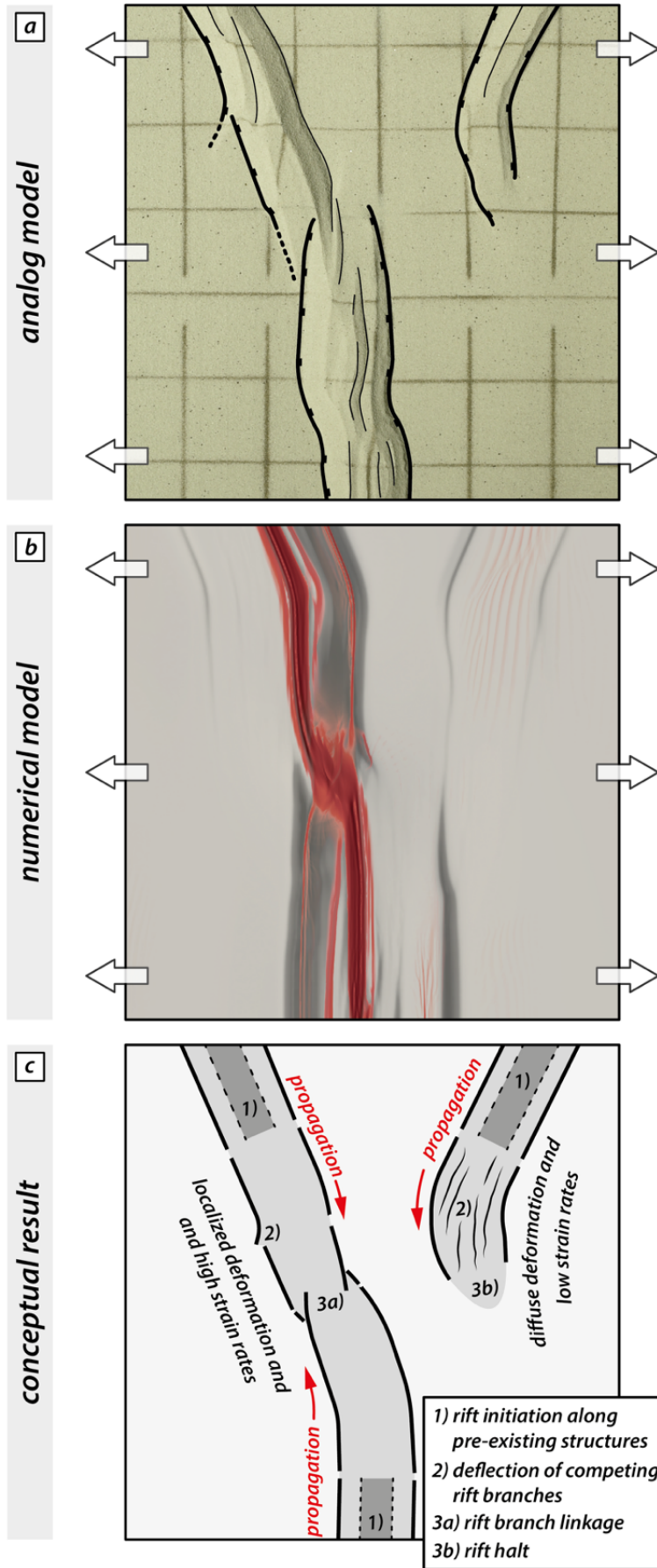
703

704 An example of rift deflection in nature has been described in the Main Ethiopian Rift.  
705 Geophysical and geologic studies evidence that pre-existing structures controlled the  
706 approximately 11 Ma southward propagation of the Northern Main Ethiopian Rift and its  
707 contemporaneous westward deflection along the Yerer-Tullu Wellel Volcanotectonic  
708 Lineament (YTVL; Abebe et al., 1998; Keranen and Klemperer, 2008; Muhabaw et al., 2022).  
709 Only after the rotation of the principal stress direction at about 5-6 Ma (Bonini et al., 2005),  
710 extension along the YTVL ceased and deformation localization along the Central Main  
711 Ethiopian Rift became more favorable. Our models document similar rift deflections and  
712 moreover indicate that, even in the absence of changing plate motions, rift segments deflect,  
713 and may cease while competing rift segments prevail and strain further localizes.

714

715 For the Canyonlands National Park, it has been proposed that it is mainly the lateral offset  
716 between pre-existing structures that explains the diversity of structures (Allken et al., 2013;  
717 Fig. 1d). With larger offsets, interaction between adjacent rift segments is limited and  
718 competing grabens persist and endure ongoing propagation coevally. We find that stresses,  
719 in combination with the geometry of pre-existing structures, play an important role and that  
720 they have a mutual effect on one another. Hence, stress distribution must be considered as  
721 an important factor especially in early rifting stages when segments link and predetermine  
722 strain localization during subsequent progressive rifting.

723



**Figure 9:** Summary plot showing the geometric similarity of rift segment linkage, deflection of competing branches and abandonment in analog and numerical models. a) Observed key features at the final stage of the analog model. b) Final strain and strain rate pattern in the numerical reference model. c) Conceptual interpretation of rift deflection and linkage based on our analog and numerical results (for details see text).

## 725 **5. Conclusions**

726 We present a series of analog and numerical rifting experiments. Our results suggest that,  
727 even in a relatively simple iso-viscous two-layer crustal setup, pre-existing weaknesses  
728 substantially disturb the regional stress pattern, which impacts rift propagation and the overall  
729 rift evolution. The complex stress re-orientation is distinct for different seed configurations  
730 (i.e., v-seed, i-seed, and y-seed) and closely interacts with the final rift geometry. The most  
731 important findings of our study can be summarized as follows:

732

- 733 • Our numerical experiments reproduce rift segment deflection seen in our analog  
734 models. This highlights the robustness of our results and their applicability to  
735 interpreting rift segment propagation, interaction, and linkage in natural settings of  
736 continental rifting.
- 737 • Pre-existing structures may control localization of rift segments that successively  
738 propagate into previously undeformed areas. Consequently, stress re-orientation  
739 initially occurs along pre-conditioned zones and propagates, coevally with rift segment  
740 propagation and strain accrual, into formerly undeformed areas.
- 741 • Interacting stresses between two competing rift segments may cause outward  
742 deflection of the propagating rift tips resulting in a successively broader rift geometry  
743 along-strike.
- 744 • Outward deflection of competing rift segments is less prominent if an opposingly  
745 propagating rift segment is present. With progressive extensional deformation, strain  
746 accrual along one of the competing rift segments prevails whilst faulting activity along  
747 the other segment ceases. Coevally, the general stress orientation changes from a  
748 symmetric to an asymmetric distribution indicating the onset of rift linkage.
- 749 • Our modelling results reproduce first-order structures of natural examples from the  
750 East African Rift System and, on smaller scale, graben structures in the Canyonlands

751 National Park. The combined investigation of surface stresses and strain localization  
752 provides an explanation for distinct rift deflection among competing rift segments and  
753 rift linkage structures where ongoing deformation and stresses mutually affect each  
754 other.

755

756 While changes in rift orientation are often used to infer regional palaeo-movements, we  
757 demonstrate that local stress field re-orientations can occur under constant plate motions.  
758 Albeit on a smaller scale, implications from our observations corroborate findings from  
759 previous studies (Brune; 2014; Duclaux et al., 2020; Gapais et al., 2000). Locally, stress and  
760 strain can largely deviate from a regional, far-field pattern and instead represent local  
761 deformation interference. In addition, the observed stress re-orientations change over time  
762 indicating that stresses measured in natural examples may depict transient stages that change  
763 with progressive deformation due to subsequent changes in material strengths (Morley et al.,  
764 2004). This implication must be considered in processing local fault-slip data when interpreting  
765 the evolution of rifts at any scale.

## 766 **Data availability**

767 Rheological measurements of the used analog materials are available in the form of open  
768 access data publications provided by the GFZ Data Service (brittle materials: Schmid et al.,  
769 2020a; Schmid et al., 2020b; viscous materials: Zwaan et al., 2018).

770

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## 782 **CRedit authorship contribution statement**

783 Timothy C. Schmid: Conceptualization, Methodology, Investigation, Formal Analysis, Writing  
784 – original draft, Visualization, Data curation. Sascha Brune: Conceptualization, Methodology,  
785 HPC funding acquisition, Supervision, Project administration, Writing – review & editing. Anne  
786 Glerum: Methodology, Software, HPC funding acquisition, Writing – review & editing. Guido  
787 Schreurs: Writing – review & editing, Supervision, Project administration, Funding acquisitions,  
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789

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