

Response to David Rees Jones :

First of all, the authors would like to thank you for the constructive reviews, that helped us improve the quality of the manuscript,. In particular, we appreciated your remarks about the clarity of notations, the justification of numerical choices, and the clarity of modelling assumptions, and have made changes to the text accordingly. Please find below our responses to your remarks.

Best regards,
Fabien Souillé, on behalf of the authors.

- Code/data: I was not sure whether the code developed/used in this study was already available and if so where? The submission guidelines suggest: Authors are encouraged to deposit software, algorithms, and model code in FAIR-aligned repositories/archives whenever possible. These research outputs are then cited in the manuscript using the received DOI and included in the reference list. The manuscript must then include a section entitled "Code availability" or, in the case of data and code, "Code and data availability".
The code developed for this study is not directly available. However, most of it is implemented and available in the open-source TELEMAC-MASCARET software within the module KHIONE dedicated to ice modeling. This includes both single- and multiple-class models. The only difference is that the semi-implicit theta scheme with the linear system presented in this publication is not incorporated in KHIONE yet.
- L18: Could mention geophysical contexts like plumes of Ice Shelf Water under floating ice sheets.
Thank you for this suggestions. I have added a sentence and references on the topic of ice shelf water plumes.
- L48: Sentence probably needs splitting.
As recommended, I split the sentence.
- L50: Perhaps add or substitute a geophysical example (e.g. from seismology)
Thank you for pointing out seismology, because there is some really insightful literature in this field especially on how to deal with models uncertainties. I have changed the sentence to include references to both geophysical and environmental modeling.
- L56: A different kind of thing, but probabilistic methods are sometimes used in processing observational data (Frazer, E. K., Langhorne, P. J., Leonard, G. H., Robinson, N. J., & Schumayer, D. (2020). Observations of the size distribution of frazil ice in an Ice Shelf Water plume. Geophysical Research Letters, 47, <https://doi.org/10.1029/2020GL090498>)
Thank you for this very interesting reference, which I have added to the introduction.
- L74: The choice of e to denote thickness was a bit confusing as the letter is usually reserved for Euler's number. I'm not sure if there is a particular precedent or motivation. Perhaps e related to edge length? Similarly, the definition of a above equation (2) might be somewhat confusing (as it denotes only part of the surface

area, not the whole). If there is a source of both of these choices come from, it would be good to mention it.

The choice of notation for the thickness was difficult to make since « t » and « h » were already used. We chose e because it stands for « épaisseur » in French (literal translation of thickness), but this was probably not the best choice as you pointed out. So I've changed the notation to: λ .

For the definition of «a»: in previous works, there is a common assumption that frazil crystals grow from their peripheral area (which we denoted a), but Holland and Feltham (2005) proposed in their model that the crystals melt from their whole surface (denoted s). In the present study, we kept similar assumptions to be consistent with previous modeling works. L55 we expose this hypothesis : «We suppose that frazil crystals grow from their peripheral area a_i but melt from their surface s_i (Holland and Feltham, 2005)».

- L92: not very clear what ΔT refers to at this stage (I see you come back and discuss it later, so it would be sensible to add some cross references and/or consider consolidating the discussion)

ΔT was here used as a generic notation to refer to either the radius or the thickness, depending on the choice that is made for the scaling. It is actually similar to f/H in your paper (D. W. Rees Jones and A. J. Wells: Frazil-ice growth rate and dynamics). $f = f_2 = 1 \Leftrightarrow \Delta T = \text{thickness}$ and $f = f_3 = H/R \Leftrightarrow \Delta T = \text{radius}$. I have added a sentence to explain that more clearly.

- L108-109: It seems to be assumed that seeding and secondary nucleation produce crystals of the same size. This might be as good as any other assumption but could be spelled out more clearly

Absolutely, this was not clearly mentioned. I added a sentence to state this assumption.

- L118: while relying on this precedent is fine, I think the main issue is that the flocculation rate in this form is independent of the number of the number of crystals. The nucleation term written (2) in equation (1) will be proportional to the square of the number of crystals while the flocculation term written (3) will be linear in the number of crystals. [I would think that both these terms arise from the same type of processes and should both be quadratic.]

Yes I agree with your reasoning on the Flocculation term. I've added more details on this in the text. I think that there is really a gap in terms of modeling on flocculation, If we compare for example with what is done in sedimentology. As a perspective for this work on multiple-size-class models, It would be interesting to study the sensitivity to different formulations for both nucleation and flocculation, and see what formulations fit better to the observed evolution of distributions.

But I think the difficulty here would be that distributions only gives us an idea of the global balance between all processes combined. Validating a single process independently from the others would be a challenge.

- L180-188: the reference to Appendix A could have been put nearer the start of the paragraph? The whole of this paragraph was relatively technical and could have been move the appendix, perhaps. It would have been good to have a brief explanation of how this scheme was chosen (especially because other studies didn't necessarily use the same method).

As suggested, I moved most of the details on the numerical scheme to appendix A.

Concerning the scheme choice, the stability condition imposes to set time steps that are relatively small. This condition is less restrictive with a semi-implicit method which allows larger time steps. This was the reason why we chose the semi-implicit method. Indeed, the used Monte Carlo sampling for the uncertainty investigation implies thousands of simulations to run, and the scheme choice was in this context crucial to lower the cost of the study. The semi-implicit approach was also used by Wang and Doering (2005). For other studies, the time scheme is not always described, so my guess is that Euler forward is more often used since it is easier to implement, despite the fact that the time step constraint can be very limiting.

- Figure 1: I found the graphs quite cluttered and hard to read. Could experiment with different colour schemes, larger figure panels (there was quite a bit of white space), perhaps one/two less data series. Perhaps the left panel would have been more useful plotted at some later time instead (e.g. $t=300s$).

To improve the visual clarity of this Figure, I removed $m=500$ and $m=1000$. I also added the number of crystals per class at $t=300s$ on left plot.

- Table 1: I spotted n_{max} here but missed where it was discussed in the main text (I saw it in L160 but not discussed in paragraph starting L105). It is quite a significant fudge factor so needs discussing somewhere in words when the notation is first introduced.

The parameter is introduced in section 2.1 when secondary nucleation is presented. $\langle n \rangle = \max(N, n_{max})$ is the average number of particles per unit volume that take part in the collisions, and n_{max} is a fitting parameter controlling the efficiency of the collision breeding». Perhaps this needed a reminder in L160, that I added.

- Section 3: I felt this could be moved to an appendix/supplement. It wasn't clear that anything was particular to this manuscript. Perhaps a one paragraph summary could go at the start of what is currently section 4.

The other referee, Dr Mark Loewen, had the same comment. I moved most technical details of section 3.3 in an appendix. However I do prefer to keep a section 3 to explain the methodology that was followed, I think it's important to understand part 4 and 5, especially for readers that are not familiar with statistical methods.

- L360: Presumably the minimum threshold is related to what you assume about nucleation. Secondary nucleation is about breaking off fragments of ice off, so presumably the minimum size might relate to this process and might not necessarily be the same as the scale that might be expected from classical nucleation theory.

Yes, that is true, thank you for pointing it out. Secondary nucleation may feed a specific radius, or a selected range of radii, which are not necessarily the smallest radius class. That way, secondary nucleation would be independent from the discretization of the radius space but at the cost of new parameters. In the current study, we assumed that secondary nucleation only feeds the smallest radius class in the model, which makes the minimum radius connected to secondary nucleation. This assumption needed to be stated out more clearly in the text, and we adapted the manuscript accordingly.

- L386-395: Of course, in geophysical contexts, this uncertainty is even worse.

- L415: e and r should be italicized.
This was corrected.
- Around equation (25), perhaps link back to equations (3) and (4) somewhere. It's not immediately clear how great a range of uncertainty there is in Nu as Nu is a rather complex function of the parameters. Also the logarithm symbol should be in Roman font.
As recommended, I added reference to equations (3) and (4) and used Roman font for the logarithm symbol.
During our investigations, we also tested Nu as an uncertain parameter initially. To estimate the uncertainty bounds and PDF of Nu we propagated the PDF of turbulent parameters and radius through equations (3) and (4) and obtained interesting results for the PDF of Nu: it was multi-modal, and the values ranged from approximately 1 to 25. However, we decided not to use this approach since Nu can't be considered independent from the radius. Considering it as uncertain may need more sophisticated modeling of probabilistic dependency and can be a good perspective (this is actually linked to your following remark). Another challenge may concern the discontinuities in functions (3) and (4).
- L445: A more general issue is that the choices of parameters are not independent, but may trade off against each other. I think this could have been discussed more strongly at various points.
Yes, we tried to make choices of parameters to avoid modeling the probabilistic dependency for this first investigation. For example, we decided not to take Nu as a parameter, but selected turbulent parameters instead. But I agree that underlying dependencies remain. To clarify this, I have added more elements on dependency in section 4 and put more emphasis on that in the «limits and perspectives» section.
- L447 paragraph: the meaning of these symbols is defined quite a long way removed from this section, perhaps a bit more of a reminder of what they mean might help follow the paragraph.
I have added a reminder for the description of these parameters.
- Figure 4: there seems to be a slope variation from w proportional to r to w proportional to r^n where $n \approx 1/2$ which will reflect different dynamical regimes. The simplified approach (a_d) constant doesn't have this feature.
Yes, that's true. If we look at the data on which most of the laws were fitted, there is a significant scattering for small radius. The slope can therefore vary a lot for small radius. We kept the law of w that represented the median behavior, so that adding a constant envelope actually contains the majority of the scattered data.
To take the change of slope into account, an option would be to model the dependency of w to r. Another idea would be to take a simple but generic enough law, like $w = a \cdot r^n$ and model the uncertainty of both parameters a and n. But defining proper distribution for a and n is not straightforward.
- Table 2: consider adding a column of parameter names/descriptors
I added a column of description like in Table 1.
- Table 3: Perhaps add the simplifications ($t_s, a_d=0$) in cases 1-2 more clearly (or in the caption).
I added $t_s=0, a_d=0$ in the caption.

- L517/L550: the sensitivity to initial conditions that you find is quite worrying for users of such models, as the initial conditions will be hard to know/control. I wondered if you had thought about what controls how long the ICs matter for?

It seems that the IC is very influential prior to maximum supercooling. So I don't think the model is actually good at predicting the maximum supercooling point, unless it has been calibrated on very similar setups to the one modeled. With that said, what surprised me the most when I started working on these models is the steady state that does not depend on thermal growth parameters but only on the heat sink rate (in the absence of buoyancy removal). So IC doesn't really matter to estimate the amount of frazil past maximum supercooling. Of course when adding buoyancy, the steady state depends on the equilibrium between the heat sink rate and buoyancy removal, in which case I don't think the well mixed models are precise enough to be predictive. However, this limitation doesn't come from IC but rather from buoyancy velocity.
- L572: the logic seems somewhat back-to-front here. It isn't the initial distribution that's key here, but rather what you assume about nucleation processes.

I agree on the fact that the assumption of nucleation feeding only the first class has an impact on the results. However, I think that there may be other reasons why we observe different behaviors by changing minimum radius (2b) and initial distribution (2c). In fact the volume growth rate is bigger for small classes so, the more initial concentration is attributed to small classes, the faster the maximum supercooling is reached. This is what we observe in case (2c). In this section I wanted to stress out that it is a combination of both nucleation and thermal growth processes that makes the choice of r_{\min} and initial distribution important. I rephrased part of this section to reflect that idea.
- L634: this might be true in the lab but not clear in the field

I modified the text to precise that.
- L639: this is convincingly shown in this context (and is an important outcome), but might not be true in more complex situations where there is more complex evolution of the mean crystal size.

I agree that the SSC model might not be representative of the complex situations observed in reality. The fact that the mean radius is the most influential here is only a property of the SSC model, since using it requires choosing a mean radius, and does not necessarily reflect reality.
- L741: Rees Jones, D. W. (not Jones, D. W. R)

It was corrected, I apologize for the mistake.